

## Printed subthreshold organic transistors operating at high gain and ultralow power

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**Abstract (<125 words):** Overcoming the trade-offs between power consumption, fabrication-cost and signal-amplification has been a long-standing question for wearable electronics. We report a high-gain, fully inkjet-printed Schottky-barrier organic thin-film transistor amplifier circuit. The transistor signal amplification efficiency is 38.2 siemens per ampere, which is near the theoretical thermionic limit and an ultralow power consumption of <1 nanowatt. The use of a Schottky-barrier for the source gave the transistor geometry-independent electrical characteristics and accommodated the large dimensional-variation in inkjet-printed features. These transistors demonstrated good reliability with negligible threshold-voltage shift. We demonstrate this capability with an ultralow-power high-gain amplifier for the detection of electrophysiological signals and showed a signal-to-noise ratio >60 decibels and noise voltage <0.3 microvolt per hertz<sup>1/2</sup> at 100 hertz.

25 **One Sentence Summary:** A low cost, low power, stable and flexible transistor fabrication technology is used to make a wearable micro-volt resolution electrophysiological signal detection amplifier.

30 **Main Text (<2500 words):** Organic thin-film transistors (OTFTs) have driven the development in low-cost, large-area electronics, including emerging application areas (1-9) such as wearable technologies. These applications require devices that are bendable and stretchable without affecting their electrical behavior (10, 11). Organic semiconductors have been widely investigated for this application, but circuits usually require a large operating voltage, leading to high power consumption and unsuitability for battery powered operation (12-14). The most challenging part of wearable electronics is the sensor interface, which is an analog application requiring low-voltage, low-power circuits with high gain, very high input impedance, low noise (15) and simple, low-cost fabrication (16, 17).

40 To meet these requirements, we use an inkjet-printed circuit technology (18) based on a subthreshold Schottky-barrier (SB) OTFT that operates near the OFF state. This approach has three

main advantages (19). First, these transistors exhibit a steep subthreshold slope, which allows the use of a low operating voltage and leads to a high transconductance efficiency. Second, the current-voltage ( $I$ - $V$ ) characteristics are independent of the channel length for a broad range of device geometries. These characteristics are ideal for printed electronics, because the variation in the typical inkjet-printed feature size of  $\sim 40$   $\mu\text{m}$  can be as much as 10  $\mu\text{m}$  (fig. S12). Third, the intrinsic gain of the SB-OTFT transistor is large (e.g.,  $> 1000$ ) and independent of channel length and electrical bias, with a  $V$ - $I$  signal amplification efficiency approaching the theoretical limit of  $q/k_{\text{B}}T$ , where  $q$  is the elementary charge,  $k_{\text{B}}$  is Boltzmann's constant, and  $T$  is temperature.

Defect density must be minimized within the printed structure to ensure a good Schottky-barrier contact at the source-semiconductor interface (19). The Schottky-contact energy barrier for hole injection into the organic semiconductor is established by the difference between the work function of the metal and the highest occupied molecular orbit level (HOMO) in the organic semiconductor (19-21). We used 2,7-dioctyl[1]benzothieno[3,2-b][1]benzothiophene (C8-BTBT) as the semiconductor (Fig. 1A), which exhibits fast growth ( $< 1$  min) of large crystals ( $> 50$   $\mu\text{m}$ ), and has a lower HOMO level compared to pentacene or other derivatives (22) to yield a good Schottky-barrier ( $> 0.2$  eV) (19, 21). Polyvinyl cinnamate (PVC) was used as the dielectric layer to provide a smooth interface between the semiconductor and dielectric, thus minimizing carrier trapping and scattering (23). A fluoropolymer encapsulation layer (CYTOP) protected the device from environmental effects. Silver was used for the metal parts. All of these materials were formulated as inks with good jetting properties (Fig. S1), and all of the fabrication steps for the individual SB-OTFT transistors and amplifier circuits reported here were carried out using a single inkjet printer tool.

The SB-OTFT demonstrated a near-zero threshold voltage ( $V_{\text{T}} = -0.01$  V; Fig. 1B) along with an ultrastep subthreshold slope of  $SS = 60.2$  mV/decade (Fig. 1C) that approached the theoretical thermionic limit (20):

$$SS_{\text{theoretical}} = \ln(10)v_{\text{th}} = 59.6 \text{ mV/decade (at } T = 300\text{K)}, \quad (1)$$

where  $v_{\text{th}} = k_{\text{B}}T/q$  is thermal voltage. In addition, this steep  $SS$  is repeatable (Fig. S3). The small  $V_{\text{T}}$  and steep  $SS$  were resulted from the low trap density (20):

$$V_{\text{T}} = V_{\text{T,theoretical}} + \frac{Q_{\text{t}}}{C_{\text{i}}}, \quad (2)$$

and

$$SS = SS_{\text{theoretical}} \left( 1 + \frac{q^2 D_{\text{t}}}{C_{\text{i}}} \right), \quad (3)$$

where  $Q_{\text{t}}$  is the trap carrier density, coulomb per  $\text{cm}^2$ ,  $D_{\text{t}}$  is the defect trap density, per eV and  $\text{cm}^2$ , and  $C_{\text{i}}$  is the gate insulator capacitance, farad per  $\text{cm}^2$ .  $Q_{\text{t}}$  and  $D_{\text{t}}$  can be affected by defects in the semiconductor bulk (e.g., grain boundaries and stacking faults) and at the semiconductor/dielectric interface (e.g., interface roughness and atomic species/vacancies on dangling bonds). The relatively large semiconductor crystals in the TFT channel ( $> 50$   $\mu\text{m}$ , providing good coverage over the channel, Fig. S2E) significantly reduce grain boundaries and stacking faults, as compared to the amorphous or micro-polycrystalline phases. The printed polymer dielectric layer was free of dangling bonds and provided a smooth semiconductor/dielectric interface (with roughness of 2.1

Å, Fig. S2C). This was comparable to the roughness of the silicon/silicon dioxide in state-of-the-art CMOS technologies. Thus, reducing the defect density to a very low level gives the best values for  $V_T$  and  $SS$ ; furthermore, the variation in these values between devices was much less than for other vacuum-deposition-based TFT technologies (Fig. 1, D and E, Table S1). Note that TFTs with a large  $C_i$  are effective in reducing  $V_T$  and  $SS$  (24), but lead to higher operating current. While this boosts the switching speed in logic circuits, it does not benefit the low-power, low frequency operation of analog sensor interfaces.

We investigated the nature of the defect density and of the Schottky barrier through the density of states (DOS), see Fig. 2A and the effective Schottky-barrier height ( $\Phi_{\text{eff}}$ ). These results suggest that the DOS comprises a small and constant background of deep states ( $g_{\text{deep}} = 6.59 \times 10^{14} \text{ cm}^{-3} \text{ eV}^{-1}$ ), a broad spectrum of delocalized states with a characteristic energy of 24.8 meV near  $\nu_{\text{th}}$ , and a steeply rising number of localized tail states with a characteristic energy of 6.7 meV. In addition, because the DOS was dominated by extended states (following  $\sqrt{E - E_{\text{HOMO}}}$ ), there was a clear mobility edge for energies above the HOMO level (i.e.,  $E > E_{\text{HOMO}}$ ) characteristic of a small overall DOS. Because the semiconductor source potential ( $\phi_s$ ) cannot be neglected in low-voltage TFTs, this term was included in the DOS calculation (Eqs. S5 to S23).

The source-side Schottky-barrier height ( $\Phi_{\text{eff}}$ ) decreased with increasing  $-V_{\text{GS}}$ , so that the drain current ( $I_{\text{DS}}$ ) was modulated by the gate bias.  $\Phi_{\text{eff}}$  could be extracted from temperature-dependent  $I$ - $V$  measurements (Fig. S6). In the subthreshold regime,  $\Phi_{\text{eff}} = \zeta_0 V_{\text{GS}} + \Phi_{\text{eff},0}$ , where  $\zeta_0$  is a coefficient that describes the modulation of Schottky-barrier height by  $V_{\text{GS}}$  (19).  $\Phi_{\text{eff}}$  showed a good initial Schottky-barrier of  $\sim 0.51$  eV and a high barrier-lowering factor of  $\zeta_0 = 1.24$ . This result suggests that charge-carrier injection was mainly by thermionic emission with smaller contributions from thermionic field emission and tunneling (see inset of Fig. 2B). Note that above a certain  $V_{\text{GS}}$  level (in the case shown at  $-0.34$  V), barrier lowering saturated, and the transistor behaved ohmically in the above-threshold regime. This change occurred when the source-side depletion width reached just a few nanometers and allowed charge carriers to tunnel through the Schottky barrier (Fig. S5C) (20, 21). The small defect density and the presence of a good Schottky barrier in the subthreshold regime were prerequisites for a high  $g_m$  and  $r_o$ .

The near-zero  $V_T$  was important for low-power operation, whereas the ultrastep  $SS$  was important for high transconductance ( $g_m = \partial I_{\text{DS}} / \partial V_{\text{GS}}$ ) and transconductance efficiency ( $g_m / I_{\text{DS}}$ ) (Eqs. S2 and S3). In addition, the SB-OTFT operation was channel-length independent with a large output resistance ( $r_o = \partial V_{\text{DS}} / \partial I_{\text{DS}}$ , Fig. 1F), which was provided by the Schottky barrier at the source-semiconductor contact. Thus, the SB-OTFT could provide a high intrinsic gain (defined as  $A_i = g_m r_o$ ) (25), resulting from the high transconductance and output resistance.

Both the transconductance and output resistance had an exponential dependence with an inverse proportionality on  $-V_{\text{GS}}$ , because of the response of SB-TFTs in the subthreshold regime (Fig. 2C) as was also the case previously with an inorganic SB-TFT (19). In comparison with other TFT technologies, the SB-OTFT transconductance and output resistance is about ten times higher at similar currents, i.e.,  $g_m = 3.8 \times 10^{-8}$  siemens and  $r_o = 3.2 \times 10^{10}$  ohm at  $I_{\text{DS}} = 1$  nA (1-6, 24, 26-28). The intrinsic gain  $A_i$  was determined from the theoretical expression (19):  $A_i = \frac{SS_{\text{theoretical}}}{SS} n \exp\left(\frac{v_{\text{sat}}}{n v_t}\right)$ , where  $n$  is the ideality factor (here,  $n=1.6$ ). These devices showed a high and constant value for  $A_i$  of  $\sim 1100$  in the subthreshold regime (Fig. 2D), which is much larger than that of the inorganic SB-TFT and Si-MOSFET because of the ultrastep  $SS$ . More importantly,  $g_m / I_{\text{DS}}$  for the SB-OTFT was  $\sim 38.2$  S/A, approaching the theoretical limit for TFT technologies of

$q/k_B T$  (i.e., 38.7 S/A at  $T = 300$  K). The high  $g_m/I_{DS}$  (indicating a large  $g_m$  at low  $I_{DS}$ ) was essential for an amplifier circuit to achieve high gain at low power. The SB-OTFT reported here exhibited more efficient  $V$ - $I$  signal amplification compared to the other reported devices (Fig. 2E).

The usability of inkjet-printed OTFTs is commonly limited by their short shelf life and operational instabilities (29, 30). However, when the transfer ( $I_D$ - $V_{GS}$ ) characteristics of representative SB-OTFT were tested over a period of 3 months under ambient conditions, no appreciable changes were observed (Fig. 3A). The threshold voltage shift was  $<1$  mV and the transconductance efficiency changed by  $<1\%$ , thus far superior in ambient environment operation and storage than typical OTFTs where these changes are typically  $>100$  mV and  $>20\%$ , respectively (30, 31).

Similarly, the effect of electrical and illumination stress was very small (29-31). Electrical stress was applied under an ON-state condition (i.e.,  $V_{GS} = V_{DS} = -3$  V), in which a conducting channel was formed and charge carriers were more likely to be trapped compared to the nearly OFF-state condition. The transfer characteristics of the device before and after stress were almost identical (Fig. 3C). The threshold voltage shifted by  $<30$  mV with the characteristic decay time of  $\sim 10^3$  s and the transconductance efficiency changed by  $<2\%$  (Fig. 3D). Because of the wide band-gap of C8-BTBT (Fig. S7), the device demonstrated excellent light stability (Fig. 3E) under visible light illumination stress ( $10$  mW/cm<sup>2</sup>), with a photocurrent  $<10$  attoampere/ $\mu$ m and a threshold voltage shift within 1 mV (Fig. 3F).

Noise ultimately limits the minimum detectable signal in any circuit, especially at the low frequencies of many electrophysiological signals ( $<100$  Hz). The low-frequency noise response of SB-OTFTs showed both  $1/f$  and white noise (Fig. 3G). As expected, these noise components were proportional to the current as  $I^2$  and  $I$ , respectively (Fig. S8, C and D, Eqs. S29 and 30). Thus, by operating in the subthreshold regime, the noise was reduced, giving rise to a signal-to-noise ratio (SNR) of 63 dB over the cut-off frequency of the TFT (Fig. 3H), which is sufficient for most low-frequency analog applications. The flicker noise coefficient is fabrication process dependent, and the value in our devices was  $1.8 \times 10^{-22}$  V<sup>2</sup> farad, which is one order of magnitude lower than that found in typical amorphous Si and metal oxide based TFTs and two orders of magnitude lower than conventional OTFTs (Table S1) (32). The root-mean-square noise voltage referred to the gate  $\sqrt{\langle v_{gn}^2 \rangle}$  for all noise sources is  $<0.3$   $\mu$ V/Hz<sup>1/2</sup> at 100 Hz (Fig. 3H), which is a few orders of magnitude lower than that of other TFT technologies for the same operating current.

We integrated amplifier circuits from pairs of SB-OTFTs in a common-source configuration, i.e., a drive transistor  $T_D$  and a bias transistor  $T_B$  (Fig. 4A). Because of the very high  $A_i$  of the SB-TFT, the amplifier demonstrated steep  $V_{out}$  characteristics and a voltage gain ( $A_V = \partial V_{out} / \partial V_{in}$ ) of 260 V/V at the peak (Fig. 4B). Because transistor  $T_B$  operated in the subthreshold regime with a bias current  $I_B = 342$  pA in the saturation regime, the power consumption was  $<1$  nW (Fig. 4C). Compared to other TFT amplifiers, this high-gain amplifier enabled high-resolution ( $<4$   $\mu$ V) of electrophysiological signal detection (Fig. 4D). In addition, the gain-bandwidth product was scalable by gate bias. Given a maximum electrophysiological signal frequency of 50 Hz (33), the SB-OTFT had a relatively large allowed bias window for analog circuit design (0.13 V) compared to the variation of  $V_T$ .

Such an amplifier can be used to monitor human electro-oculogram (EOG) signals, which are essentially the corneo-retinal potentials ( $V_{EOG}$ ) that exist across the front (positive) and back (negative) of the human eye (Fig. 4F), typically in the range from 0.2 to  $\sim 1.0$  mV (34). This

technique is useful for eye-movement tracking, particularly in improving existing technologies that are bulky, costly and require high power (35). With a biasing electrode over the eyebrow and another electrode below the lower eyelid connecting to the amplifier input (Fig. 4, A and F), the  $V_{in}$  relation for the amplifier becomes

$$V_{in} = V_{bias} + \gamma V_{EOG}. \quad (4)$$

Here,  $\gamma$  is a coefficient that describes the direction of eye movement. In the configuration used,  $\gamma < 0$  corresponds to an upward movement of the eyeball, whereas  $\gamma > 0$  indicates the corresponding downward movement. Therefore, the amplifier output gives an amplitude of up to  $\sim 0.3$  V (Fig. 4G, Movie S1). The amplifier is also able to track horizontal eye movement (Fig. S11). The amplified EOG signal with amplitudes  $> 0.2$  V and SNR  $> 60$  dB has the potential to detect subtle eye movements for a better depiction of the virtual environment, e.g., depth-of-field rendering. Tracking eye movement is important in virtual and augmented reality (35). The ultralow power consumption of SB-OTFT based circuits can potentially operate from energy acquired from micro-harvesters (in the order of  $\mu\text{J}/\text{cycle}$ ) (8), although from a complete system standpoint this would require low-power versions of signal conditioning and transmission circuit stages.

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**Competing interests:** The authors declare no competing interests. **Data and materials availability:** All data are available in the main text or the supplementary materials.

**Supplementary Materials:**

Materials and Methods

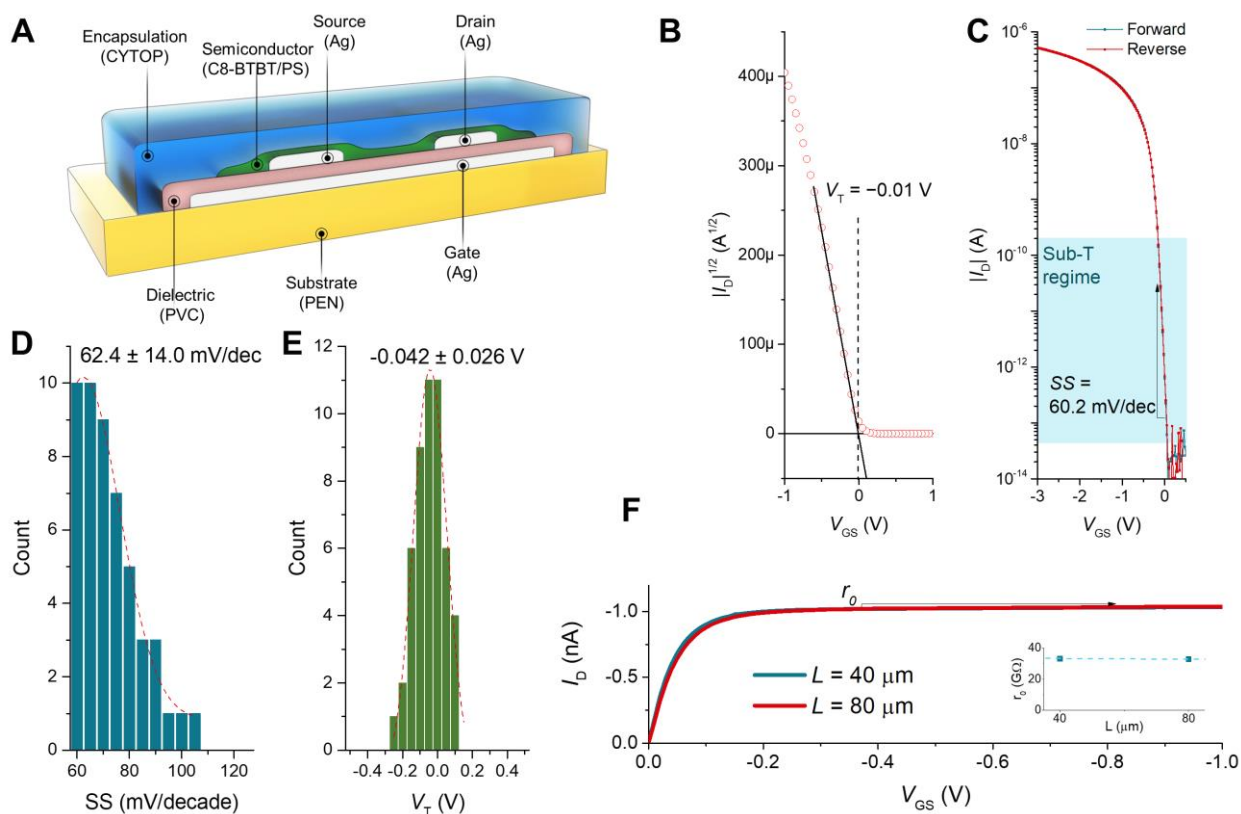
5 Supplementary Text

Figures S1-S12

Tables S1-S2

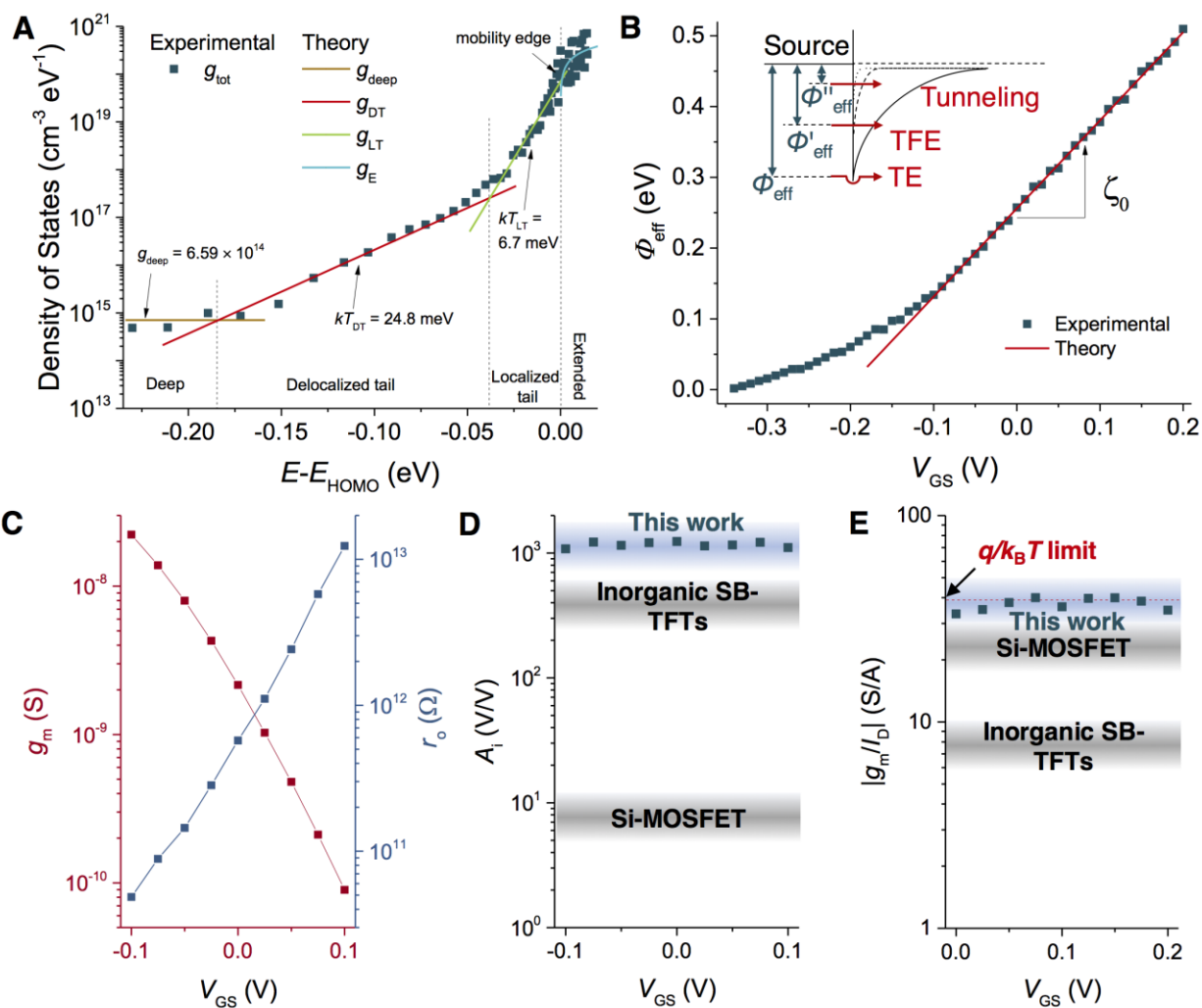
Movie S1

References (36-61)



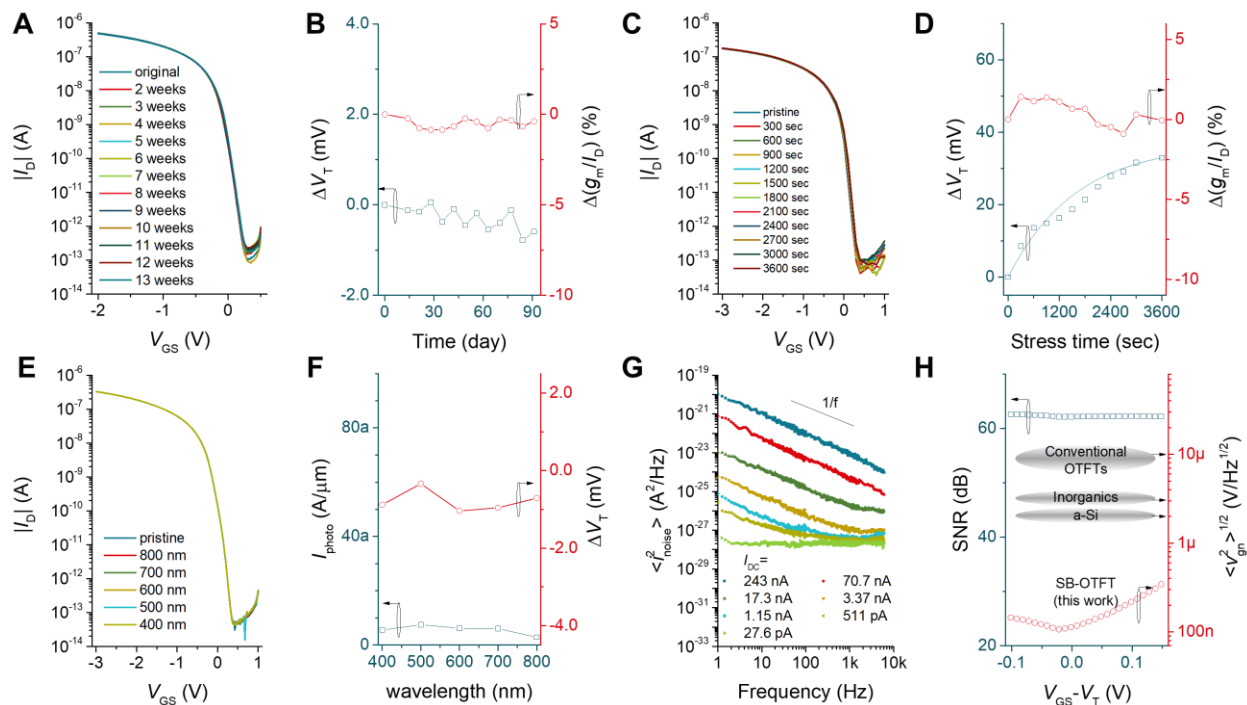
**Fig. 1. Device structure and electrical characteristics.** (A) Schematic cross section of the SB-OTFT. Measured transfer characteristics ( $I_D$  vs.  $V_{GS}$ ) of a typical device (B) on a linear-scale, indicating the threshold voltage ( $V_T$ ), and (C) on a log-scale, indicating the subthreshold slope ( $SS$ ). Statistical distributions of (D)  $SS$  and (E)  $V_T$  for 50 devices. The dashed lines indicate normal distributions. (F) Measured output characteristics ( $I_D$  vs.  $V_{DS}$ ) indicating the output

resistance ( $r_0$ ) of devices with different channel length ( $L$ ) showing a full overlap of the characteristics. Inset shows  $r_0$  vs.  $L$ .



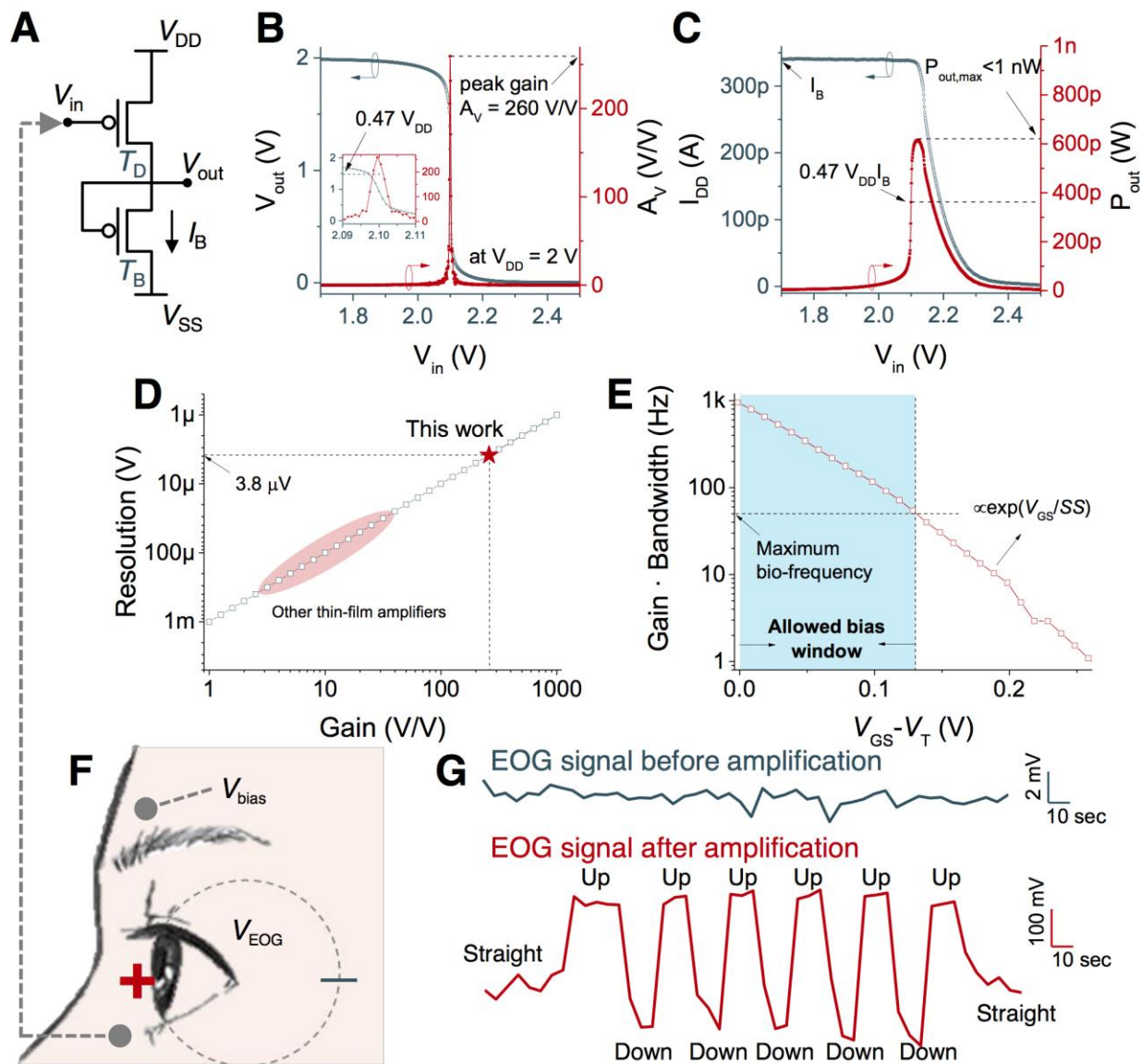
**Fig. 2. Static parameters.** (A) Density of states (DOS) for a typical device, indicating four different regimes: deep states, delocalized tail (DT) states, localized tail (LT) states, and extended states. The slopes in the DT and LT regimes indicate the characteristic energies ( $k_B T_{\text{DT}}$  and  $k_B T_{\text{LT}}$ , respectively). (B) Effective Schottky-barrier heights ( $\Phi_{\text{eff}}$ ) as a function of  $V_{\text{GS}}$ , indicating the gate-modulation factor ( $\zeta_0$ ) for the  $\Phi_{\text{eff}}$  lowering (inset: schematic energy band diagram showing variation in effective  $\Phi_{\text{eff}}$  and different charge carrier injection processes). (C) Experimental values for  $g_m$  and  $r_0$  as a function of  $V_{\text{GS}}$ . (D) Measured intrinsic gain ( $A_i$ ) as a

function of  $V_{GS}$ . (E) Experimental values of transconductance efficiency ( $g_m/I_D$ ) as a function of  $V_{GS}$  reaching the theoretical thermionic limit of 38.7 S/A.



**Fig. 3. Stability and reliability.** (A) Measured transfer characteristics for storage under ambient conditions for the times indicated and (B) change in absolute threshold voltage ( $\Delta V_T$ ) and change in relative transconductance efficiency [ $\Delta(g_m/I_D)$ ] as a function of time. (C) Measured transfer characteristics under negative bias stress ( $V_{GS} = V_{DS} = -3$  V) for the stress time indicated, and (D)  $\Delta V_T$  and  $\Delta(g_m/I_D)$  as a function of stress time. (E) Measured transfer characteristics under light exposure, (F) photocurrent ( $I_{photo}$  in  $A/\mu m$ ) and  $\Delta V_T$  for different wavelengths (400~800 nm). (G) Measured SB-OTFT current noise under different direct current biases ( $I_{DC}$ ). (H)

Signal-to-noise ratio (SNR) in the near-threshold and subthreshold regimes, and input-referred voltage noise density at 100 Hz.



**Fig. 4. Amplifier characteristics and demonstration of electro-oculography detection.** (A) Schematic circuit diagram of a common-source amplifier. (B) Measured output voltage ( $V_{out}$ ) and gain ( $A_V$ ) as a function of input voltage ( $V_{in}$ ). (C) Measured operating current ( $I_{DD}$ ) and power ( $P_{out}$ ) as functions of  $V_{in}$ . (D) Resolution of electrophysiological signal detection as a function of gain. (E) Gain-bandwidth product as a function of  $V_{GS}$  in the subthreshold regime. (F) Operating principle and circuit configuration for electro-oculography (EOG) amplification with the amplifier. (G) EOG signal obtained before and after amplification.