

# Strain Analysis and Engineering in Halide Perovskite Photovoltaics

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## Abstract

Halide perovskites are a compelling candidate for the next generation of clean energy harvesting technologies thanks to their low cost, facile fabrication and outstanding semiconductor properties. However, photovoltaic device efficiencies are still below practical limits and long-term stability challenges hinder their practical application. Current evidence suggests that strain in halide perovskites is a key factor in dictating device efficiency and stability. Here, we outline the fundamentals of strain within halide perovskites relevant to photovoltaic applications and rationalize approaches to characterize the phenomenon. We examine recent breakthroughs in eliminating the adverse impacts of strain, enhancing both device efficiencies and operational stabilities. Finally, we discuss further challenges and outline future research directions for placing stress and strain studies at the forefront of halide perovskite research. An extensive understanding of strain in halide perovskite is needed, which would allow effective strain management and drive further enhancements in efficiencies and stabilities of perovskite photovoltaics.

28 Halide perovskites have been thrust into the limelight in the last decade as a promising next-  
 29 generation photovoltaic technology, with power conversion efficiencies (PCEs) already  
 30 comparable with those of crystalline silicon devices. Reductions in unwanted energy losses  
 31 both in the absorber and at interfaces with contact layers have led to substantial efficiency and  
 32 stability improvements. Recently, strain control has been at the forefront of device efficiency  
 33 enhancement in halide perovskites by minimizing undesirable defect formation<sup>1,2</sup>, and  
 34 subsequent nonradiative recombination<sup>3</sup>. A certified PCE of ~24.4% has recently been reported  
 35 for FAPbI<sub>3</sub>-based devices (FA=formamidinium) through managing strain, with the  
 36 encapsulated devices retaining 90% of their efficiencies after 400 hours at maximum power  
 37 point<sup>1</sup>. Such encouraging achievements suggest that strain engineering is a powerful tool to  
 38 enhance efficiencies and stabilities for a broad range of practical applications. Despite such  
 39 successes, there is an apparent paradox whereby similarly strained halide perovskites exhibit  
 40 either beneficial or detrimental effects even when they are nominally similar materials.  
 41 Furthermore, the magnitude of strain is relatively high in halide perovskites compared to other  
 42 photovoltaic systems: for instance, residual strain values of up to ~2.4% for  $\alpha$ -FAPbI<sub>3</sub><sup>4</sup> and  
 43 ~1.6% for CsPbBr<sub>3</sub> have been reported<sup>5</sup>, whereas the performance of Si and Cu(In, Ga)Se<sub>2</sub>  
 44 devices drops substantially when residual strains exceed ~1%<sup>6,7</sup>. The high magnitude of strain  
 45 in perovskite materials is particularly striking since perovskites are one of the most  
 46 mechanically fragile of any photovoltaic materials<sup>8</sup>. It is therefore essential to thoroughly  
 47 understand the fundamentals of strain on different length scales, and to devise effective  
 48 methods to characterise it. This will allow us to move towards rational engineering as opposed  
 49 to empirical optimisation.

## 50 1. Definition and measurement of strain in halide perovskites

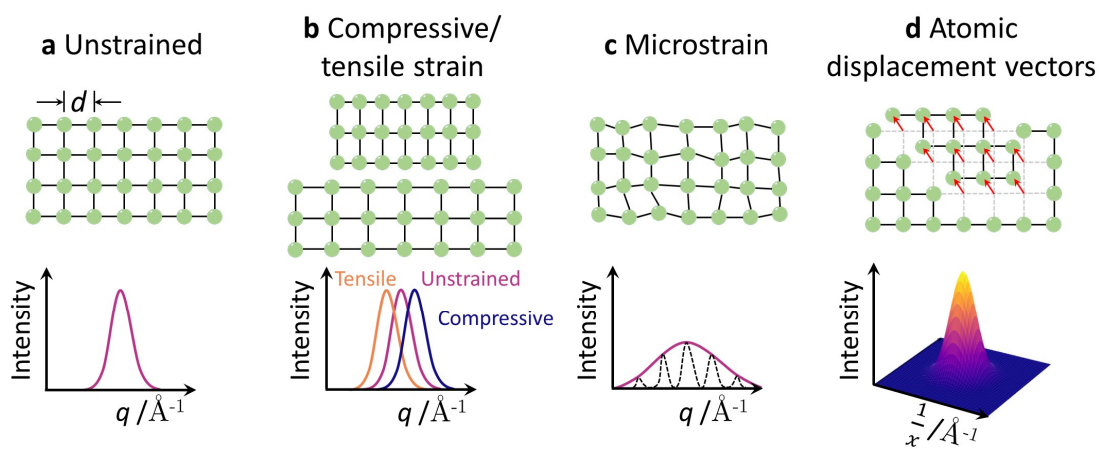
51 Strain ( $\epsilon$ ) is defined as a structural deformation of a material in response to applied stress( $\sigma$ )<sup>9</sup>.  
 52 This stress can be from external forces on the material, or internal structural defects such as  
 53 atomic interstitials and vacancies. Mathematically, strain is defined by a (second rank) tensor  
 54 with elements given by

$$55 \quad \epsilon_{kl} = \frac{1}{2}(e_{kl} + e_{lk}) .$$

56 Here,  $e_{kl} = \frac{\partial u_k}{\partial x_l}$ , where  $u_k$  are the components of the atomic displacement vector  $\mathbf{u}(x_1, x_2, x_3)$   
 57 along the Cartesian directions, and  $x_l$  are the Cartesian coordinates.  $\mathbf{u}(x_1, x_2, x_3)$  describes the  
 58 displacement of atoms from their expected positions in an unstrained material, and the above  
 59 definition is valid for small displacements. This particular linear combination of  $\frac{1}{2}(e_{kl} + e_{lk})$   
 60 removes the effects of rigid rotation from the strain tensor<sup>10</sup>. Determination of this tensor  
 61 constitutes a full description of strain in materials as the atomic displacement vectors can be  
 62 reconstructed from the tensor elements. Such a description not only encapsulates the length  
 63 changes along the Cartesian axes given by  $\epsilon_{x_1x_1}$ ,  $\epsilon_{x_2x_2}$ , and  $\epsilon_{x_3x_3}$ , but also the shear  
 64 components of the deformation given by the off-diagonal elements. In the elastic regime, the  
 65 magnitude of the deformation resulting from the applied stress (also a tensor,  $\sigma$ ) can be

66 calculated from Hooke's law,  $\sigma_{ij} = c_{ijkl}\epsilon_{kl}$ , where  $c_{ijkl}$  are the components of a (fourth rank)  
 67 stiffness tensor.

68 The aspects of strain in a material that can be determined with typical laboratory techniques  
 69 include tensile and compressive strain where the interplanar spacing,  $d$ , of the material is  
 70 compared to a reference "unstrained" value. Typically, the variation in  $d$  spacing is observed  
 71 through the Bragg peak positions in an X-ray diffraction (XRD) pattern. If a material is  
 72 compressively strained, the  $d$  spacing is decreased relative to the reference, and the peak  
 73 associated with that plane will shift to higher scattering vector,  $\mathbf{q}$  (Fig. 1a,b), with the converse  
 74 being true for tensile strain. The values of compressive/tensile strain are usually expressed as  
 75 a ratio of the  $d$  spacings and written as percentages. The nature of strain (compressive/tensile)  
 76 is indicated either explicitly or using a negative/positive sign indicating compressive/tensile  
 77 strain, respectively. However, we note that such measurements only probe the diagonal ( $\epsilon_{kk}$ )  
 78 elements of the strain tensor. Furthermore, applying a load (stress) to a material results in  
 79 deformations (strain) not only in the directions parallel to, but also perpendicular to the load  
 80 (Poisson's effect). For example, a biaxial in-plane tensile strain in a halide perovskite thin film  
 81 will also result in a compressive strain in the direction of the film normal<sup>11</sup>. This is especially  
 82 pertinent because many lab-based X-ray diffractometers are in Bragg-Brentano (reflection)  
 83 geometry with the scattering vector in the direction of the film normal. Consequently, it is the  
 84 out-of-plane compressive strain that will be probed.



85

86 **Fig. 1| Definition and measurement of strain in halide perovskites. a-d,** The top row gives  
 87 representations of (a) unstrained material, (b) compressive/tensile strain, (c) microstrain, and  
 88 (d) atomic displacement vectors. The bottom row indicates the signatures of each of these kinds  
 89 of strain in diffraction patterns.

90 An additional feature of strain reported is "microstrain" which is determined by Bragg peak  
 91 width analyses. If the diffraction measurement simultaneously samples a number of different  
 92  $d$  spacings within the sample, then contributions from all these  $d$  spacings will combine to  
 93 form a broadened Bragg peak (Fig. 1c). Microstrain is a measure of local deviations of  $d$   
 94 spacings from their average value and includes contributions from both locally compressively  
 95 and locally tensile strained regions. A commonly employed method for calculating microstrain

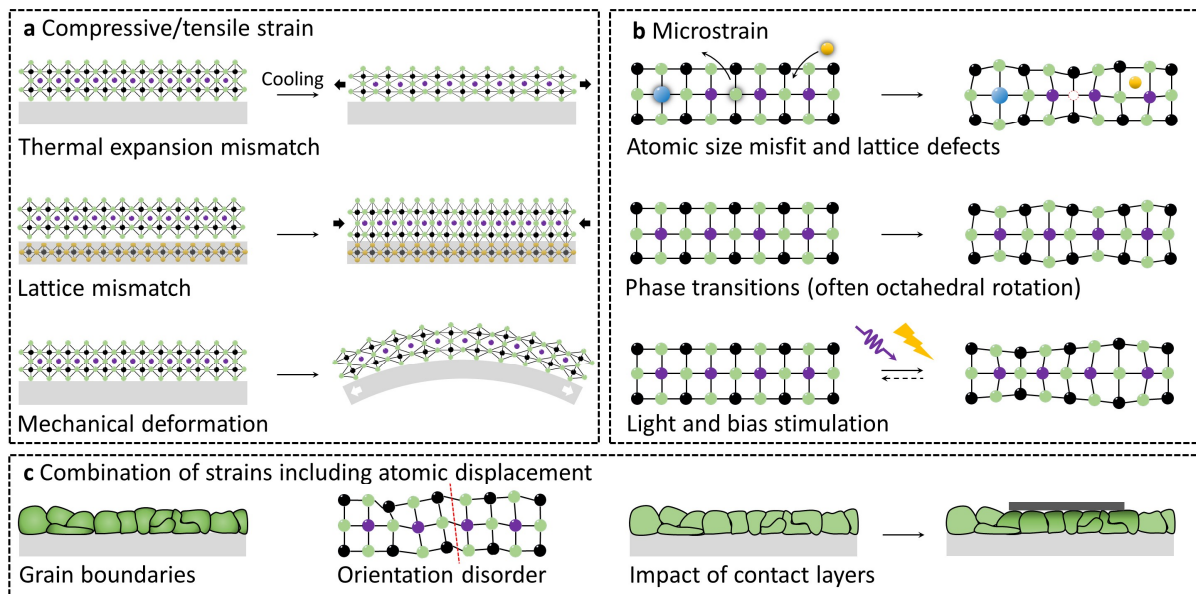
96 is Williamson-Hall analysis<sup>12</sup>. The method relies on the different  $q$ -dependence of Bragg peak  
97 broadening resulting from microstrain and crystallite size effects and, despite the method's  
98 ubiquity, there are many potential pitfalls. One must first remove any instrument broadening,  
99 which is typically accomplished using an external standard, but these standards may not be in  
100 the same form as the sample (e.g. thin film) and thus do not exactly reproduce the conditions  
101 of the diffractometer for the sample. Further, many nuances influence the crystallite size  
102 calculation including the distribution of sizes, shape, and crystallite length in each  
103 crystallographic direction. Errors in calculation of particle size will have a knock-on effect for  
104 determination of microstrain. Depending on the structural features (dislocations, variable  
105 lattice distortion, grain surface relaxation) there may or may not be a strong  $hkl$  dependence  
106 on microstrain peak broadening, which complicates matters further. Therefore, the researcher  
107 must be especially diligent when carrying out peak broadening analyses from lab-based bulk  
108 XRD measurements<sup>12</sup>. We also stress that microstrain values are not directly comparable to  
109 compressive/tensile strain values given that these two quantities probe different aspects of  
110 material structure. The community must avoid interchangeable use of these terminologies to  
111 prevent errors.

112 Another important consideration is the relevant length scale of the measurement. All of the  
113 lattice planes that satisfy the Bragg condition within the sample interaction volume (defined by  
114 the footprint and penetration depth) of the X-rays contribute to a given diffraction peak. In  
115 most lab-based X-ray diffractometers, the beam footprint is on the scale of millimetres squared  
116 and so the interaction volume is large and variations in  $d$  spacings in this region result in a  
117 broadened Bragg peak. However, by reducing the beam footprint and scanning the illumination  
118 position across the sample, the local shifts in  $d$  spacing can be resolved in the form of a tensile  
119 and compressive strain map. This forms the basis for spatially resolved XRD techniques  
120 (micro- or nano-XRD) that are able to measure local peak shifts or phase changes<sup>3</sup>. Overall,  
121 one must consider the relevant length scales to probe in halide perovskites given that  
122 heterogeneities are present over nanometers up to many micrometers<sup>13</sup>.

123 Other advanced techniques are required to probe beyond the diagonal elements of the strain  
124 tensor corresponding to structure expansion/contraction. By using the three-dimensional  
125 intensity distribution around Bragg peaks, one can reconstruct the electron density of a crystal  
126 to find the atomic displacement vectors (Fig. 1d) from which the off-diagonal components of  
127 the strain tensor can be extracted. Techniques that can provide such rich information even on  
128 the local scale include Bragg coherent diffraction imaging and Bragg ptychographic  
129 approaches<sup>14, 15</sup>.

## 130 **2. Origins of strain in halide perovskites**

131



132

133 **Fig. 2| Origins of strain in halide perovskites. a-c,** Representations of the origins of (a)  
 134 compressive/tensile strain, (b) microstrain, and (c) more complex strain (including  
 135 contributions from shear and rotational components of the strain tensor) in halide perovskites.  
 136 The blue and purple spheres are the A-site cations (e.g., Cs, MA and FA), and the black and  
 137 green spheres represent metal cations (e.g., Pb, Sn) and halide anions (e.g., Cl, Br and I),  
 138 respectively. In a, the orange spheres are substrate atoms. In (b), the red dashed circle is a  
 139 vacancy, and the yellow sphere represents the extrinsic dopant that causes a defect. The red  
 140 dashed line in (c) indicates a twin boundary.

141 Strain can result from a broad range of factors. During film formation, the precursor solutions  
 142 are deposited on a substrate and the film is annealed (often above 100 °C). If the film is formed  
 143 on a substrate with a lower coefficient of thermal expansion (CTE) than the perovskite ( $3.3\sim 8.4$   
 144  $\times 10^{-5} \text{ K}^{-1}$ )<sup>16</sup>, the substrate restricts the contraction of the perovskite along the in-plane  
 145 directions during cool down<sup>17</sup>, resulting in an in-plane tensile strain (Fig. 2a, top). Analogously,  
 146 compressive strain can occur in a film grown on a substrate with a higher CTE than that of the  
 147 perovskite, though this is less common in halide perovskite devices. This process of cooling  
 148 after annealing leads to strain (generally tensile) in the plane of the film.

149 Lattice mismatch between the substrate and the as-formed perovskite film can also generate a  
 150 tensile or compressive strain (Fig. 2a, middle), depending on whether the *d* spacing of the  
 151 substrate is larger or smaller than that of the perovskite overlayer. However, when the lattice  
 152 mismatch is large, the perovskite crystal randomly stacks on the bottom crystal, resulting in a  
 153 non-uniform interfacial distortion<sup>18, 19</sup>. Compressive and tensile strains can also be introduced  
 154 to films using mechanical film bending<sup>20</sup> (Fig. 2a, bottom) and hydrostatic pressure<sup>21</sup>.

155 While the above factors generally influence strains across the whole film, other factors  
 156 including atomic misfits, phase transitions, light/bias stimulation and grain boundaries are  
 157 linked to local strain (namely microstrain) in halide perovskites (Fig. 2b). For example, in  
 158 FAPbI<sub>3</sub>, the large FA cation within the 3D structure promotes local PbI<sub>6</sub> octahedral tilting (that

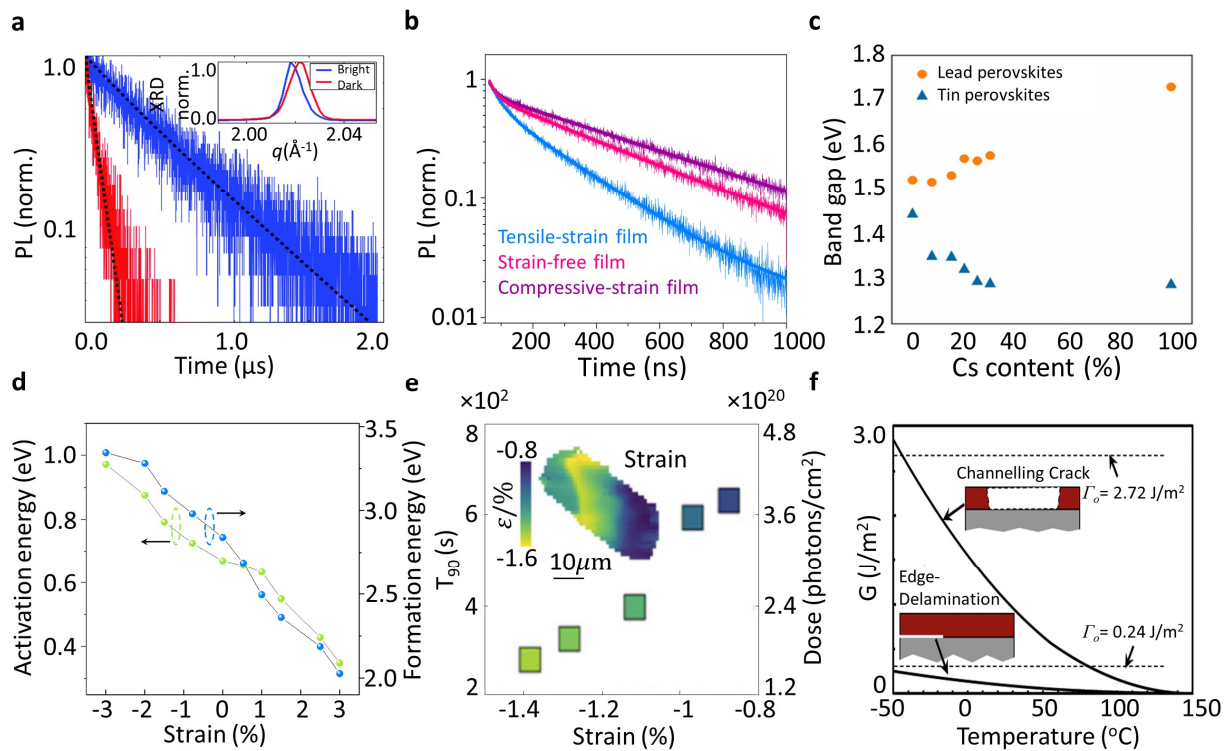
159 is, the Pb-X-Pb angle deviates from the ideal  $180^\circ$ ), resulting in a local lattice distortion<sup>2, 22</sup>.  
160 Further, local lattice distortion and strain are greatly elevated during phase transitions, where  
161 the symmetry of the crystal structure is generally lowered compared to the cubic perovskite  
162 structure. Group theoretical analysis<sup>23</sup> of octahedral tilting schemes yields group-subgroup  
163 relationships between perovskite phases (e.g.  $I4/mcm$  is a subgroup of  $Pm\bar{3}m$ ) which can  
164 transform in a continuous (i.e., second-order) fashion, leading to spontaneous strains due to  
165 continuous lattice parameter variation between phases upon octahedral tilting<sup>24-26</sup>. Large strains  
166 and complex microstructure, such as martensitic nano/microstructures<sup>26</sup> or crystal twins, can  
167 result from discontinuous (i.e., first-order) phase transitions with abrupt changes in lattice  
168 parameter, which can occur between non-group-subgroup related space groups (though first-  
169 order transitions between group-subgroup phases are also possible). Such a transition is found  
170 in MAPbI<sub>3</sub> (MA = methylammonium) between  $I4/mcm$  to  $Im\bar{3}$ , at a hydrostatic pressure of  
171  $\sim 0.3$  GPa<sup>21</sup>. Martensitic transformations are diffusionless and often involve a lowering of  
172 crystal symmetry (e.g., from cubic to tetragonal<sup>27</sup>) and when the resulting phases are mutually  
173 spatially constrained, strain minimisation can result in a change in shape and curvature of their  
174 interface<sup>28</sup>. Therefore, one can use strain to vary the nano/microstructure resulting from  
175 martensitic transformations, as has been shown, for example, in SrSnO<sub>3</sub><sup>29</sup>. Given the  
176 prevalence of symmetry lowering phase transitions in halide perovskites, a more detailed  
177 understanding of the strains resulting from phase transitions will be beneficial.

178 Photovoltaic devices are exposed to light and bias during operation, and these have been  
179 reported to affect microstrain in halide perovskites. Recently, it was suggested that microstrain  
180 is enhanced under illumination owing to light-assisted ion migration and/or phase segregation<sup>30</sup>.  
181<sup>31</sup>. Here, it was proposed that photoexcited carriers within perovskites interact with the soft,  
182 ionic lattice<sup>32, 33</sup>, resulting in greater microstrain. Further, some works attempt to probe the  
183 links between photo-induced lattice expansion and microstrain (excluding the effect of ion  
184 migration), but debate is ongoing<sup>34, 35</sup>. Under illumination, the polarity (positive or negative)  
185 and magnitude of bias voltages have also been reported to be critical in inducing microstrain  
186 by enhancing ion migration<sup>30</sup>.

187 Grain boundaries can be another source of strain in halide perovskites (Fig. 2c, left) with the  
188 resulting strain dependent on the specific interfacial conditions, e.g., orientation mismatch at  
189 boundaries<sup>36</sup>, of which there is striking variation within thin films. It has been proposed that  
190 defect formation at certain grain boundaries in perovskite thin films stems from asymmetric  
191 strain fields at junctions between grains<sup>37</sup>. The high degree of lattice distortion and non-  
192 stoichiometric chemical distribution at grain boundaries also cause large strain fields<sup>3, 38</sup>. Apart  
193 from boundaries between grains, sub-grain lattice orientation disorder<sup>36</sup> (Fig. 2c, middle) and  
194 twin boundaries may also be sources of strain in halide perovskite films<sup>39</sup>. Contact layers are  
195 also known to exert forces on the adjacent perovskite films in devices, however, we note that  
196 this is highly dependent on the thickness and stiffness of contact layers (Fig. 2c, right)<sup>20, 40</sup>.  
197 Given the substrate's comparatively larger thickness, the substrate will dominate in imparting  
198 strain in perovskite devices<sup>20</sup>.

### 199 3. The effects of strain on halide perovskites

200 Strain affects the structure of halide perovskites and consequently influences the optoelectronic  
 201 properties and stability of the material, which both need to be optimised to push perovskite  
 202 devices towards commercialization.



203

204 **Fig. 3| Effects of strain on halide perovskites.** **a**, Time-resolved photoluminescence (PL)  
 205 decays for differently strained regions of a MAPbI<sub>3</sub> film where the PL lifetime in the more  
 206 compressively strained region (red; dark PL) is shorter than that in the more tensile strained  
 207 region (blue; bright PL). Inset: 220 Bragg peaks show the difference in lattice spacing for the  
 208 two regions. **b**, Time-resolved PL decays for three (FAPbI<sub>3</sub>)<sub>0.85</sub>(MAPbBr<sub>3</sub>)<sub>0.15</sub> films with  
 209 different strains. **c**, The variation of bandgap with Cs content (scales with lattice parameter) for  
 210 FA<sub>1-x</sub>Cs<sub>x</sub>PbI<sub>3</sub> (orange circles) and FA<sub>1-x</sub>Cs<sub>x</sub>SnI<sub>3</sub> (blue triangles). **d**, Calculated strain-dependent  
 211 activation energies of ion migration and formation energies of halide vacancies in CsPbI<sub>2</sub>Br  
 212 where the minus/plus values in the *x*-axis represent compressive/tensile strain, respectively. **e**,  
 213 Time for the XRD pattern correlation coefficient to decay to 90 % of its original value (*T*<sub>90</sub>)  
 214 against strain. The right *y*-axis shows the sample's radiation dose. Inset: strain map of the  
 215 CsPbBr<sub>3</sub> microcrystal from which these data are taken. **f**, Calculated strain energy release rate,  
 216 *G*, for channelling cracks and edge delamination.  $\Gamma_0=0.24$  J/m<sup>2</sup> (polycrystalline) and  $\Gamma_0=2.72$   
 217 J/m<sup>2</sup> (single-crystalline) are minimum and maximum cohesive fracture energies for MAPbI<sub>3</sub>,  
 218 respectively. Figures adapted with permission from: **a**, ref.<sup>3</sup> under a Creative Commons licence  
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 222 American Chemical Society; **f**, ref.<sup>11</sup>, Elsevier.

### 223 3.1 Impact on device performance

224 The impact of strain on device performance is complicated by the apparently contrasting  
225 responses of different perovskite systems to strain and the problem of defining an ‘unstrained’  
226 reference state. To illustrate this, Fig. 3a shows time-resolved photoluminescence (PL) for a  
227 MAPbI<sub>3</sub> film<sup>3</sup>, showing regions with greater local compressive strain (relative to the average  
228 *d* spacing of the film) associated with the shorter PL lifetime, albeit the net strain on the film  
229 is unknown. By contrast, in (FAPbI<sub>3</sub>)<sub>0.85</sub>(MAPbBr<sub>3</sub>)<sub>0.15</sub> films (Fig. 3b) the PL lifetime is  
230 reported to increase upon compressive strain<sup>41</sup>. The reason for these apparent differences may  
231 be complex, and one must consider how to best benchmark absolute strain and how the different  
232 compositions or growth methods will impact both strain and recombination, for example, due  
233 to photodoping in the mixed-component systems<sup>42</sup>. Further, Fig. 3c shows that the bandgap  
234 increases for FA<sub>1-x</sub>Cs<sub>x</sub>PbI<sub>3</sub> but decreases for FA<sub>1-x</sub>Cs<sub>x</sub>SnI<sub>3</sub> with increasing Cs content, despite  
235 a reduction in lattice parameter in both cases<sup>43</sup>. These conflicting trends are attributed to  
236 octahedral tilting dominating over isotropic structural contraction in FA<sub>1-x</sub>Cs<sub>x</sub>PbI<sub>3</sub> and highlight  
237 the importance of probing the various structural degrees of freedom in the perovskite.

238 To maximise device efficiency, any nonradiative recombination must be minimised as this  
239 leads to power losses. While most defects in halide perovskites are benign, there are still deep  
240 trap states, which are detrimental to performance. For example, large variations of intragrain  
241 crystallographic orientations and concomitant local strains are linked to higher trap densities<sup>36</sup>.  
242 Furthermore, larger grains exhibit more sub-grain misorientation leading to higher trap  
243 densities. This suggests that managing local strain is even more critical than tuning  
244 morphological grain size for minimising nonradiative recombination. At the sub-grain level,  
245 twinning has been shown to be ubiquitous in FA-based systems due to the low twin formation  
246 energies<sup>44</sup>; however, these are expected to be benign from a carrier recombination perspective<sup>39</sup>.

247 To understand the effects of strain on the band structure of halide perovskites we must first  
248 recognise that the valence band is governed by the anti-bonding overlap between Pb *s* orbitals  
249 and halide *p* orbitals, while the conduction band is dominated by the non-bonding overlap  
250 between Pb *p* orbitals with minor contributions from halide *p* orbitals. Strain-induced changes  
251 in Pb-X bond lengths and bond angles (Pb-X-Pb) and associated changes in orbital overlaps  
252 lead to changes in band dispersion<sup>45</sup>. In turn, the bandgap increases with increasing tensile  
253 strain and decreases with increasing compressive strain<sup>41</sup>. Carrier mobility has been shown to  
254 increase with increasing compressive strain due to a reduction of carrier effective mass, though  
255 at higher magnitudes of compressive strain the mobility decreases due to generation of  
256 dislocations<sup>4</sup>. In 2D A<sub>2</sub>PbBr<sub>4</sub> [(A= *n*-butylammonium, *n*-octylammonium, or  
257 phenethylammonium (PEA)]<sup>46</sup>, mobility was increased by tuning the cation stacking pattern,  
258 alleviating microstrain, and reducing electron-phonon interactions. Few-layer 2D perovskites  
259 are also observed to exhibit an expanded structure accompanied by a blue-shifted PL emission  
260 peak in comparison to bulk 2D perovskites<sup>47</sup>. These reports show that strain is a key  
261 optimisation parameter and an essential structural lever for modulating device performance.

### 262 **3.2 Impact on device stability**

263 Intrinsic device stability is arguably the most important topic for halide perovskites. Ion  
264 migration has been associated with device instability, owing to degradation caused by atomic



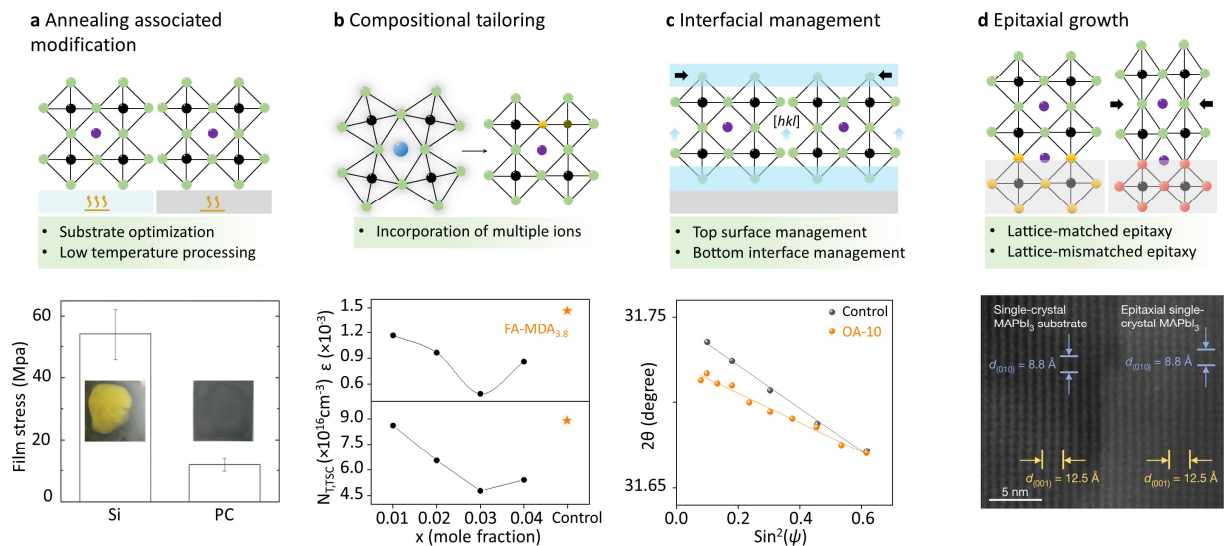
265 rearrangement or redox chemistry between migrated ions and other species. Evidence suggests  
266 that the halides are the most mobile ions in halide perovskites<sup>48, 49</sup>. The activation energy of  
267 ion migration has been reported to increase with compressive strain and decrease with tensile  
268 strain (Fig. 3d)<sup>16</sup>. This result is consistent with other reports of temperature-dependent  
269 conductivity measurements, where increasing compressive strain (relieving tensile strain)  
270 results in more stable films<sup>17</sup>. However, here substrates were generally bent along only one  
271 axis and so the stress applied was uniaxial. Due to Poisson's effect, there will be strains of  
272 opposite signs in the directions perpendicular to the direction of the uniaxial strain. As such,  
273 both tensile and compressive strains are expected to exist along different directions in the plane  
274 of the film. Furthermore, consideration of the impact of cracks and wrinkles in the film will be  
275 needed, as these features may be induced in such bending studies and will affect ion migration.  
276 The increasing compressive strain has also been shown to increase the radiation hardness of  
277 CsPbBr<sub>3</sub> microcrystals (Fig. 3e)<sup>5</sup>. Further, it has been reported that increasing hydrostatic  
278 pressure increases the range of miscible I:Br ratios in MAPb(I<sub>1-x</sub>Br<sub>x</sub>)<sub>3</sub>, and inhibits ion-  
279 migration-driven halide segregation<sup>50</sup>. There is still much to be understood about the  
280 fundamental mechanisms of halide segregation. Future work will need to explore, for example,  
281 how the material's elastic energy changes to maintain the coherency of the underlying lattice  
282 in mixed component systems with different ionic sizes<sup>51-53</sup>. The variation of the free energy as  
283 a function of composition in alloyed perovskite systems due to coherency strains may, in turn,  
284 dictate the miscibility gap of the mixed ions, and elucidate the specific segregation  
285 mechanisms<sup>54</sup>.

286 Strain also affects the stability of halide perovskites by regulating the thermodynamics of defect  
287 generation. Density functional theory (DFT) calculations show that the defect formation energy  
288 of the halide vacancies in halide perovskites decreases with increasing tensile strain<sup>16</sup>. The  
289 defect concentration in tensile strained (FAPbI<sub>3</sub>)<sub>0.85</sub>(MAPbBr<sub>3</sub>)<sub>0.15</sub> is reported to be higher than  
290 that in its less strained counterpart ( $2.6 \times 10^{16}$  versus  $2.1 \times 10^{16}$  cm<sup>-3</sup>)<sup>55</sup>. Moreover, calculations  
291 show that shallow defect states in MAPbI<sub>3</sub> close to the valence band become even shallower  
292 under low hydrostatic pressure (0.3 GPa)<sup>56</sup>. However, in FAPbI<sub>3</sub> and FA<sub>0.75</sub>CS<sub>0.25</sub>PbI<sub>3</sub>, a higher  
293 pressure (2 GPa)<sup>45</sup> has been reported to transform shallow iodide vacancy defects into deep  
294 level states. More critical comparisons on identical perovskite systems will be needed to  
295 generalize such observations. Other computational work suggests that under biaxial strain  
296 (compressive or tensile), the lowest energy phase of MAPbI<sub>3</sub> is orthorhombic, indicating the  
297 presence of a thermodynamic driving force for a phase transition from the desired cubic phase<sup>57</sup>.  
298 This contrasts with work on the FA-based analogue showing that compressive strain suppresses  
299 the transition from the cubic to tetragonal phase<sup>4</sup>. This is another example of nominally similar  
300 halide perovskite systems responding to strain in opposing ways. Similar phase retention  
301 behaviour is observed in MAPbI<sub>3</sub> films where the tetragonal to orthorhombic phase transition  
302 is suppressed for smaller grains<sup>58, 59</sup> which are speculated to be more strained than larger grains,  
303 but such strain fields are yet to be directly characterised.

304 Halide perovskites must also be mechanically robust for photovoltaic applications. Grain  
305 boundaries and film edges are sites of abrupt changes in film morphology and structure, and  
306 therefore, are regions where stress may be concentrated<sup>60</sup>. Upon loading, a mechanical fracture

307 can propagate along these boundaries and edges when stress overwhelms the energetic  
 308 threshold of film fracture. For a given stress value and film thickness, the propensity to crack  
 309 and delaminate can be calculated through their respective strain energy release rates,  $G$ , which  
 310 scales according to  $G \propto \frac{\sigma^2 h}{E}$ , where  $\sigma$  is the stress,  $E$  is the Young's modulus, and  $h$  is the  
 311 layer thickness<sup>60</sup>. If  $G$  exceeds the cohesive fracture energy,  $\Gamma_0$ , (an index reflecting the  
 312 resistance to mechanical fracture), film cracking and/or delamination will occur.  $\Gamma_0$  for halide  
 313 perovskites falls below  $1.5 \text{ J/m}^2$ , making these materials mechanically vulnerable to fracture<sup>61</sup>.  
 314 Furthermore, film cracking is more energetically favourable than film delamination (Fig. 3f)<sup>11</sup>.  
 315 An aim of the field should be to obtain absorber layers with higher  $\Gamma_0$  ( $>10 \text{ J/m}^2$ )<sup>61</sup> to improve  
 316 mechanical robustness, with this being even more pertinent for flexible devices<sup>62</sup>. Any residual  
 317 stresses in perovskites, or adjacent device layers, require careful consideration to minimise the  
 318 potential of fracturing.

#### 319 4. Strain engineering to mitigate adverse effects



320

321 **Fig. 4| Strain engineering in halide perovskite photovoltaics.** **a**, Annealing associated  
 322 modification: Top, diagram of how to prevent annealing associated strain generation. Bottom,  
 323 measured stress for MAPbI<sub>3</sub> formed on silicon (Si) and polycarbonate (PC) at 100 °C. Insets  
 324 show photographs of perovskites deposited on both substrates after 45 h of dry heat aging  
 325 (85 °C and 25% relative humidity). The perovskite degrades to the yellow phase only on the Si  
 326 substrate where there is greater thermal expansion mismatch. **b**, Compositional tailoring: Top,  
 327 illustration of the incorporation of multiple ions. Bottom, microstrain calculated in samples  
 328 consisting of fluorine-doped tin oxide (FTO)/mesoporous TiO<sub>2</sub>/(FAPbI<sub>3</sub>)<sub>1-x</sub>(MC)<sub>x</sub> (upper plot)  
 329 and trap density (lower plot) where MC is a mixture of MDA<sup>2+</sup> and Cs<sup>+</sup> with the equimolar  
 330 amount and the orange star indicates control perovskite [(FAPbI<sub>3</sub>)<sub>0.972</sub>(MDACl<sub>2</sub>)<sub>0.028</sub>],  $N_{T, TSC}$   
 331 is the trap density. **c**, Interfacial management: Top, illustration of perovskite interface  
 332 management. Bottom, linear fit of  $2\theta$  vs.  $\sin^2 \psi$  for n-octylammonium modified (OA-10) and  
 333 control perovskite films. A reduced slope represents a reduction of tensile stress in perovskite

334 films. **d**, Epitaxial growth: Top, illustration of epitaxial growth. Bottom, a high-resolution  
335 transmission electron microscope image showing the interfacial area of the epitaxial single-  
336 crystal MAPbI<sub>3</sub> on a MAPbI<sub>3</sub> substrate. Figures adapted with permission from: **a**, ref.<sup>20</sup>, Wiley-  
337 VCH (bottom); **b**, ref.<sup>1</sup>, AAAS (bottom); **c**, ref.<sup>55</sup>, Wiley-VCH (bottom); **d**, ref.<sup>76</sup>, Springer  
338 Nature Ltd (bottom).

#### 339 **4.1 Annealing associated modification**

340 Thermal annealing is a routine step in the preparation of polycrystalline perovskites. However,  
341 any thermal expansion mismatch between the perovskite thin film and adjacent layers must be  
342 considered for strain generation in perovskites and ideally minimised. The bottom panel in Fig.  
343 4a shows aged MAPbI<sub>3</sub> films, one on a Si substrate that has degraded to a yellow colour (PbI<sub>2</sub>),  
344 and the other on a polycarbonate substrate that remains as a dark absorbing film. The difference  
345 is attributed to the larger in-plane tensile strain in the film on Si due to the greater CTE  
346 mismatch between the perovskite and Si substrate<sup>20</sup>. One way to address this issue is through  
347 modification of the annealing process. A variation on standard film fabrication procedures is  
348 where the perovskite film is fully formed on the substrate at room temperature (with or without  
349 post-annealing) thereby lowering the tensile strain in the film<sup>20</sup>. Another low-temperature  
350 approach has enabled photovoltaic devices to deliver an impressive PCE of 23.1% by utilising  
351 amine assisted crystallization at room temperature<sup>63</sup>. While these reports show the potential for  
352 annealing associated modification of the deposition process, further innovation is required to  
353 avoid large residual strains in films. Comparing solution-processed films with thermally  
354 evaporated ones (in which the photoactive black perovskite phase can form even without  
355 annealing<sup>64</sup>) may also provide insights into any hidden role played by solvent molecules<sup>65</sup>.  
356 Further, a vertical temperature gradient through the device stack is present during annealing  
357 which can set up chemical and strain gradients along this direction<sup>41</sup>. However, such features  
358 are difficult to characterise without depth-dependent and multi-dimensional techniques.

#### 359 **4.2 Compositional tailoring**

360 Compositional tailoring (Fig. 4b) involves substituting different ions into lattice sites and has  
361 been reported to be effective in reducing microstrain in halide perovskites. For example, the  
362 incorporation of ions with radii smaller than FA (253 pm)<sup>66</sup>, such as MA (217 pm)<sup>66</sup> and Cs  
363 (167 pm)<sup>66</sup> at the A-site<sup>2, 22</sup>, and/or Br (196 pm)<sup>2</sup> and Cl (181 pm)<sup>2</sup> at the X-site<sup>2, 22</sup>, has been  
364 proposed to reduce the microstrain in FAPbI<sub>3</sub>-based systems, with a potential strain relaxation  
365 mechanism where these additives could decrease the B-X bond lengths and the B-X-B bond  
366 angles<sup>2</sup>, a conclusion ascertained from DFT calculations<sup>2</sup>. Another approach is to incorporate  
367 larger cations (such as PEA) to form 2D or 2D/3D perovskites. This has been shown to promote  
368 preferred crystal orientation, reduce Bragg peak width, and increase Bragg peak area (generally  
369 indicative of reduced microstrain), yielding improved device efficiency over the 3D-only  
370 controls<sup>67</sup>. Recently, simultaneous incorporation of smaller Cs and larger MDA  
371 (methylenediammonium; 262 pm) A-site cations into FAPbI<sub>3</sub> has been shown to suppress  
372 microstrain without increasing the bandgap<sup>1</sup>. The reported microstrain relaxation was  
373 associated with a decrease in defect concentration (Fig. 4b, bottom) and nonradiative  
374 recombination losses, leading to a PCE of ~24.4% and promising stability<sup>1</sup>. In this work, the

375 authors note that the simultaneous introduction of larger and smaller ions could modify the  
376 local strain in the perovskite lattice, yet the octahedral tilting they proposed, as a mechanism  
377 for such microstrain reduction still requires direct experimental evidence. Determination of  
378 local atomic structure, for example with pair distribution function analysis, high-resolution  
379 electron diffraction, or X-ray fine structure measurements, will be of great value to confirm or  
380 refute proposed strain relaxation mechanisms.

381 Compositional tailoring has achieved success in enhancing device performance but commonly  
382 used compositions employing multiple A-site cations (FA/MA or Cs/FA/MA) can be  
383 detrimental to the mechanical stability of films<sup>68</sup>. We note that current ubiquitous spin coating  
384 procedures including antisolvent quenching can result in smaller grain sizes, which is  
385 correlated with a greater susceptibility to film fracture<sup>68</sup>, accompanied by reduced cohesive  
386 fracture energy,  $\Gamma_0$ . The common practice of using a slight excess of PbI<sub>2</sub> in precursor solutions  
387 to achieve higher performances<sup>69</sup> can also reduce film stability<sup>68</sup>, in addition to any photo-  
388 instabilities it introduces<sup>69</sup>. By contrast, additive species such as 1,2-diaminopropane and  
389 butylphosphonic acid 4-ammonium chloride can be added to precursor solutions to crosslink  
390 between perovskite grains<sup>61,70</sup>, thereby reducing the film's propensity to mechanically fracture.

### 391 **4.3 Interfacial management**

392 It is of interest to examine if strain can be modified by inserting an interlayer between the  
393 perovskite and the substrate. To date, some compliant interlayers have been reported to relax  
394 strain in perovskites caused by CTE mismatch between substrate and perovskite<sup>40,71</sup>. However,  
395 given the thickness of typical interlayers (roughly a few tens of nanometers), the underlying  
396 mechanisms are likely to be more complex. For example, the use of a compliant  
397 polytriarylamine interlayer on top of a Si substrate, with a thickness of ~20 nm, will not have  
398 a significant effect on alleviating the CTE-mismatch-induced strain experienced by the top  
399 perovskite layers<sup>20</sup>. Furthermore, despite claims in recent reports<sup>18,72</sup>, interlayers are unlikely  
400 to facilitate strain control by minimising the lattice mismatch at the interface because the  
401 perovskite growth processes are unlikely to be epitaxial<sup>73</sup> (owing to the typical solution  
402 processing methods of these films) and thus lattice matching with the substrate may, in general,  
403 not be relevant. In other work, it has been proposed that the top contact layer can be modified  
404 for use as a strain compensation layer. For example, poly[5,5-bis(2-butyloctyl)-(2,2-  
405 bithiophene)-4,4'-dicarboxylate-alt-5,5'-2,2'-bithiophene], which has a larger CTE than  
406 CsPbI<sub>2</sub>Br, could lead to a compressive strain on the CsPbI<sub>2</sub>Br layer below, offsetting the  
407 residual tensile strain in the perovskite<sup>16</sup>. Constructing a 2D perovskite layer on top of the 3D  
408 perovskite has also been shown to reduce the residual stress in perovskites (Fig. 4c, bottom)<sup>55</sup>,  
409 based on 2D perovskites exhibiting a lower Young's modulus. However, care must be taken to  
410 determine whether the corresponding stress and strain reductions in perovskites are truly due  
411 to the mechanical action of these contact layers. Furthermore, efforts should be made to assess  
412 the stress and strain evolution at the 2D/3D interface, which is likely to be influenced by the  
413 orientation of the 2D perovskites relative to their 3D counterpart<sup>74</sup>. Besides, not all 2D/3D  
414 interfaces will be sufficiently stable without further management: for instance, it has been

415 proposed that invasive ions from the underlying 3D perovskites interrupt the Van der Waals  
416 interactions in the 2D layered perovskites<sup>75</sup>, which could drive layer slippage and even fracture.

417 Epitaxy is a well-established material growth technique for introducing controlled levels of  
418 strain into materials by judiciously choosing lattice-matched/mismatched substrates. Recently,  
419 successful epitaxial growth of  $\alpha$ -FAPbI<sub>3</sub> was achieved by tailoring the composition of a  
420 MAPbCl<sub>x</sub>Br<sub>3-x</sub> single-crystal substrate<sup>4</sup>. Further compositional tuning also allowed various  
421 levels of strain to be imparted onto the  $\alpha$ -FAPbI<sub>3</sub> epilayer, with low levels of compressive strain  
422 found to notably enhance the phase stability. In other work, solution-phase epitaxial growth of  
423 MAPbI<sub>3</sub> was enabled by using lattice-matched MAPbI<sub>3</sub> substrates (Fig. 4d, bottom), with an  
424 impressively large area (>25 cm<sup>2</sup>) and the absence of dislocations<sup>76</sup>. Heteroepitaxial growth of  
425 halide perovskites directly on top of charge transport layers could be useful, but this will further  
426 necessitate expanding the current library of contact layers for halide perovskite devices.  
427 Another interesting aspect of epitaxial growth is the tendency of phase impurities to epitaxially  
428 grow alongside perovskite phases<sup>77</sup>. Given the critical need to minimise phase impurities,  
429 especially those harmful to performance and stability, further research is needed to understand  
430 the role of epitaxial growth in stabilising halide perovskites. Nonetheless, epitaxial techniques  
431 are comparatively expensive and their use risks negating this key advantage of halide  
432 perovskites (that of inexpensive and scalable processability) in the face of more efficient and  
433 stable III-V or Si-based equivalents.

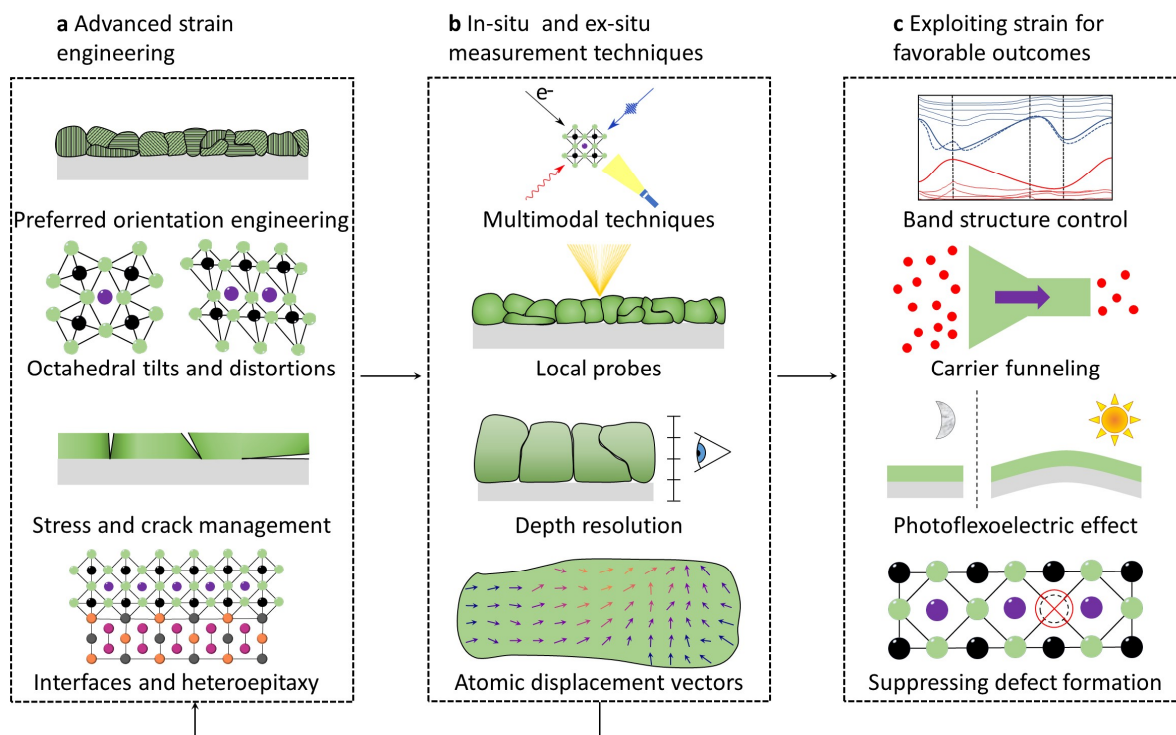
434 In addition, it has been reported that some contact layers in perovskite photovoltaics suffer  
435 intrinsic fragility<sup>78</sup> or inadequate adhesion to the perovskites<sup>79-81</sup>, and perovskite photovoltaics  
436 with these undesirable characteristics will have reduced stress tolerance and mechanical  
437 integrity. Approaches to address this issue have been proposed, including the design of  
438 materials in which strong covalent bonds can form after deposition from the solution state<sup>78</sup>,  
439 the use of self-assembled interlayers<sup>79</sup>, low-dimensional perovskite capping layers on top of  
440 3D perovskites<sup>80</sup>, and graded interfaces<sup>81</sup>. Such approaches will be critical for increasing the  
441 film's cohesive fracture energy,  $\Gamma_0$ , and ensuring mechanical integrity of the operating  
442 devices<sup>79</sup>.

## 443 **5. Challenges and Outlook**

444 To tackle the many outstanding questions about the links between strain and material properties  
445 in perovskite photovoltaics, we propose an iterative workflow whereby strain is tuned with  
446 advanced strain engineering (Fig. 5a), and more fully characterised using powerful in situ and  
447 ex situ measurement techniques (Fig. 5b) with the knowledge fed back to improve strain  
448 engineering approaches. Finally, the result (Fig. 5c) will be the ability to not only suppress the  
449 adverse effects of strain but also to exploit strain for favourable outcomes in devices.

450

451



452

453 **Fig. 5| Workflow for unlocking the potential of strain in perovskite devices.** **a**, Strain can  
 454 be applied systematically to halide perovskites with advanced strain engineering. **b**, Cutting-  
 455 edge in- and ex-situ characterisation techniques are employed to fully characterise strain in  
 456 halide perovskite samples. **c**, Strain is exploited to realise better devices.

## 457 5. 1 Advanced strain engineering

458 The association of nonradiative losses with orientation mismatch<sup>36</sup> suggests that engineering  
 459 films with greater preferred orientation should be targeted for improved optoelectronic  
 460 properties. Further, one may be able to channel strain into crystallographic directions where  
 461 the presence of strain is less detrimental. For example, microstrain in FAPbI<sub>3</sub> in the  $\langle 111 \rangle$   
 462 direction, in particular, is associated with reduced film stability<sup>22</sup>. Meanwhile, Young's  
 463 modulus exhibits a directional anisotropy, with  $E_{110} < E_{100}$ , in cubic MAPbX<sub>3</sub> (X=Br, Cl),  
 464 making it more facile for deformations to propagate in the  $\langle 110 \rangle$  direction than in the  $\langle 100 \rangle$   
 465 direction<sup>82</sup>. However, care must be taken when extracting microstrain values for particular  
 466 crystallographic directions from individual Bragg peaks, especially if assuming negligible  
 467 crystallite size broadening based on the morphological grain size given the presence of sub-  
 468 grain heterogeneities<sup>13</sup>.

469 Tilting of the octahedral BX<sub>6</sub><sup>4-</sup> units has been well studied in other systems based on the  
 470 perovskite structure<sup>83</sup> and octahedral distortions including cation off-centring are also known  
 471 in halide perovskites. However, study of octahedral tilting and distortion in halide perovskites  
 472 are often limited to Goldschmidt tolerance factor considerations, which leads to difficulties  
 473 when considering molecular cations that can form hydrogen bonds<sup>83</sup>. Although some  
 474 researchers have noted the link between octahedral tilting/distortion and strain<sup>84</sup>, explicit

475 characterisations of such structural features will be critical given the relationship between  
476 important properties such as bandgap and the largest metal-halide-metal bond angle<sup>85</sup>.

477 The determination of stress and its impact on halide perovskites are under-explored areas and  
478 merit more dedicated studies. Uncompensated stress can cause fracturing of the perovskite  
479 leading to the loss of Ohmic contacts with the charge-selective layers and allowing the ingress  
480 of detrimental moisture and oxygen<sup>86</sup>. Stress analyses could explain film wrinkling, a common  
481 mechanism by which a film can relieve stress during formation<sup>87</sup>. Additionally, from a  
482 thermodynamic perspective, spatially heterogeneous distributions of stress can result in the  
483 local accumulation of elastic energy in films which is relevant given ions will rearrange until  
484 their diffusion potentials are uniform across the film<sup>51</sup>. Stresses will occur on both short and  
485 long length scales, with delamination<sup>60</sup> and cracking<sup>88</sup> being problems at the scale of tens of  
486 micrometres, and emergent microstresses at the grain boundaries and other operational stresses  
487 being concerns on the nano to atomic scale<sup>11, 89</sup>. Defect pooling at grain interfaces can reduce  
488 the material's mechanical integrity by causing these interfaces to serve as fracture points<sup>61</sup>.  
489 Given the presence of sub-grain features and misorientations, solving the above issues will  
490 likely require strategies beyond just tuning processing methods. Stress can either be determined  
491 directly through curvature measurements<sup>20, 35, 87</sup>, or calculated indirectly, via interrogation of  
492 strain ascertained through material structure studies with use of appropriate proportionality  
493 constants. However, such indirect calculations of stress come with many complexities<sup>88, 90</sup>. A  
494 common method for the measurement of the in-plane residual stresses in films is the  $\sin^2 \psi$   
495 method<sup>88, 90</sup>.

496 Beyond the perovskite bulk, regulating harmful strain at the substrate/perovskite interface is  
497 also crucial for improving device performance. Substrates with aligned nanopores have been  
498 employed to confine crystallization and impose compressive strain on perovskites, with the  
499 strain level dependent on the pore size<sup>91</sup>. Similar results are achieved through the manufacture  
500 of textured perovskite films<sup>92</sup>. Furthermore, perovskites that are infiltrated into a porous  
501  $\text{TiO}_2/\text{ZrO}_2/\text{Carbon}$  layer have been shown to exhibit improved mechanical integrity<sup>68</sup>. Another  
502 emerging approach to reduce strain is to apply 2D materials (such as transition metal  
503 dichalcogenides) as interlayers where the Van der Waals interaction enables perovskites to  
504 expand/contract freely during heating/cooling processes<sup>93</sup>. We note, however, the key  
505 challenge of synthesising such 2D materials without dangling bonds. Regulating strain may  
506 also involve management of lattice matching at semiconductor interfaces, but this approach  
507 has received limited attention to date, probably for two reasons: first, there is a debate about  
508 whether typical solution processing of perovskites involves epitaxial growth processes, and  
509 second, the perovskite structure is sufficiently accommodating that device efficiencies have  
510 not suffered appreciably in the face of poorly lattice-matched interfaces. However, it is possible  
511 that some more sinister ramifications of poor lattice matching, such as delamination of layers  
512 in the device stack, only become evident after stringent, long-term operational testing<sup>61</sup>.  
513 Looking towards conventional semiconductors, epitaxial growth of perovskites on hetero-  
514 substrates can allow control of strain at the interface by tuning lattice matching. A halide  
515 perovskite epilayer grown on functional layers, without secondary transfer of the epitaxial  
516 perovskite, could be desirable for high-performance photovoltaic applications, and there is a

517 wealth of literature on other materials with domain matching epitaxy<sup>94</sup>. Additionally, the  
518 application of intervening layers allows high-quality epitaxial growth of semiconductors with  
519 strain applied through carefully chosen lattice mismatches<sup>60</sup>.

520 The strain management methods proposed in this work will apply to a wide range of perovskite-  
521 based photovoltaics including Sn- and Sn-Pb based devices<sup>95, 96</sup> and tandem cells<sup>97</sup>. Effective  
522 strain management will also be critical when employing perovskites on flexible substrates or  
523 upscaling to full-sized modules<sup>68</sup>. Additionally, daily or seasonal temperature, spectral, and  
524 light intensity changes present important challenges for both stress and strain management in  
525 operating devices.

## 526 **5.2 In and ex situ measurement techniques**

527 Multimodal correlative microscopy can provide deep insights into the underlying physics in  
528 halide perovskites. For example, correlation of strain fields with photoluminescence,  
529 compositional, and time-resolved data sets will uncover new links between strain and material  
530 performance<sup>3, 36</sup>. Integrating tight environmental control (of light, heat, humidity, UV exposure,  
531 etc.) and sample encapsulation with these measurement techniques will allow assessment of  
532 the evolution of strain in operating devices.

533 The presence of spatially heterogeneous strain in halide perovskite samples necessitates the use  
534 of local probes for studying strain. Current state-of-the-art film compositions contain many  
535 components which are known to segregate<sup>98, 99</sup>. To adequately assess the impacts of this  
536 segregation, a local approach to strain characterisation is required, further employing  
537 techniques such as scanning probe nano-XRD and X-ray fluorescence<sup>100</sup>. There is also  
538 evidence for rich chemical heterogeneity at the grain and sub-grain length scales<sup>13</sup>.

539  
540 In alloyed systems<sup>66</sup>, including the Pb/Sn and mixed cation and halide systems, it is still unclear  
541 how the alloyed components are arranged – whether they vary from unit cell to unit cell, or if  
542 they tend to form clusters across the sample<sup>101</sup>. In the case of mixed halide samples, given the  
543 evidence for inhomogeneous halide distributions<sup>42</sup>, the observed compressive/tensile “strain”  
544 may be due to changes in composition rather than stress. The community, therefore, needs to  
545 expand its understanding of local chemical distributions through the use of nanoscale chemical  
546 mapping techniques<sup>102, 103</sup>. Further, ion re-distribution in perovskites can gradually occur in  
547 response to strain, which, in turn, varies the strain distribution in perovskites<sup>16</sup>. To overcome  
548 these challenges, in-situ strain detection is required to view strain propagation in real-time.

549 Strain measurements to date have predominantly focused on the in-plane direction, and further  
550 work to simultaneously monitor out-of-plane strain is of interest. Furthermore, the relative  
551 importance of strain in the perovskite bulk and strain at the perovskite surface is currently  
552 unknown. A greater focus on depth-resolved techniques such as grazing incidence X-ray  
553 scattering will be required.

554 Most reports on strain to date have been on neat film samples as opposed to device stacks, and  
555 so these works miss the complex interactions of the perovskite with adjacent layers in the



556 device. Local heterogeneities in the perovskite material may also exist due to the influence of  
557 other device layers. Under electrical bias, there is evidence for photostrictive effects along with  
558 ion migration<sup>48</sup>, the activation energy of which can be altered by strain and is a vitally important  
559 phenomenon to understand in the field. Thus, developing techniques for in operando  
560 characterisations of microstrain in full and half devices for a complete understanding of the  
561 impact of strain will be fruitful.

562 Advanced coherent diffraction imaging techniques such as Bragg coherent diffraction imaging  
563 and Bragg ptychography can potentially allow reconstruction of full three-dimensional maps  
564 of atomic displacement vectors<sup>104</sup>. Resolving atomic displacement vectors would allow one to  
565 uncover strain fields in halide perovskites that may be characteristic of certain crystal  
566 aberrations, facilitating “defect diagnosis”<sup>14</sup>. Correlation of these strain fields with other  
567 optoelectronic data would yield rich insights and should be a future research focus. Atomic  
568 displacement vectors can also be resolved with cryogenic and high-resolution transmission  
569 electron microscopy measurements<sup>77, 105</sup>. However, attention must be paid to the beam dose in  
570 these measurements, which might introduce undesirable defects and/or lead to material  
571 decomposition if not performed carefully<sup>106, 107</sup>.

### 572 **5.3 Exploiting strain for favourable outcomes**

573 There is a wealth of literature on using strain in other material systems to enhance material  
574 characteristics by modulating band structures<sup>108</sup>, electrical conductivities<sup>109</sup>, electron-phonon  
575 coupling and ion migration<sup>110</sup>, among other properties<sup>111</sup>. Modulating strain could see several  
576 exotic effects realised in halide perovskites leading to new device use-cases. First-principles  
577 computational work suggests polar phases in lead-iodide perovskites may be realisable with  
578 large Rashba splitting at room temperature via tensile strain<sup>112</sup>. Being able to control Rashba  
579 splitting provides an additional lever for tuning carrier dynamics<sup>113</sup>. Distinct chemical domains  
580 with distinct bandgaps within mixed-component films may allow for carrier funnelling and  
581 devices based on such effects<sup>3, 41, 114</sup>. Strain gradients have also been utilised in halide  
582 perovskites to enhance flexoelectric responses with corresponding photoflexoelectric effects  
583 demonstrated in MAPbI<sub>3</sub> single crystals<sup>115</sup>. While engineering such exotic effects into real-  
584 world devices is yet to be realised, these reports show that strain plays a critical role in the  
585 electronic structure of halide perovskites, and that it can indeed be exploited. Strain  
586 management can lead to better material quality by suppressing the formation of defects due to  
587 ion migration in the material.

588 A robust understanding of the origins and effects of strain in halide perovskite-based devices  
589 is needed to facilitate future breakthroughs. This will be especially critical in the construction  
590 of perovskite/Si and perovskite/perovskite tandems, where many layers may each contribute  
591 both stress and strain components, influencing overall efficiency and stability. As the halide  
592 perovskite community continues to push the envelope on long-term operational stability, strain  
593 will be a key parameter for optimisation and a versatile lever for control in more exotic devices.

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### 606 **Author contributions**

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### 609 **Competing interests**

610 S.D.S. is a co-founder of Swift Solar, Inc. All other authors declare no competing interest.

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