

THE MONOTONE WRAPPED FUKAYA CATEGORY AND THE OPEN-CLOSED STRING MAP

ALEXANDER F. RITTER AND IVAN SMITH

ABSTRACT. We build the wrapped Fukaya category $\mathcal{W}(E)$ for any monotone symplectic manifold E , convex at infinity. We define the open-closed and closed open-string maps, $\text{OC} : \text{HH}_*(\mathcal{W}(E)) \rightarrow SH^*(E)$ and $\text{CO} : SH^*(E) \rightarrow \text{HH}^*(\mathcal{W}(E))$. We study their algebraic properties and prove that the string maps are compatible with the $c_1(TE)$ -eigenvalue splitting of $\mathcal{W}(E)$. We extend Abouzaid's generation criterion from the exact to the monotone setting. We construct an acceleration functor $\mathcal{AF} : \mathcal{F}(E) \rightarrow \mathcal{W}(E)$ from the compact Fukaya category which on Hochschild (co)homology commutes with the string maps and the canonical map $c^* : QH^*(E) \rightarrow SH^*(E)$. We define the $SH^*(E)$ -module structure on the Hochschild (co)homology of $\mathcal{W}(E)$ which is compatible with the string maps (this was proved independently for exact convex symplectic manifolds by Ganatra). The module and unital algebra structures, and the generation criterion, also hold for the compact Fukaya category $\mathcal{F}(E)$, and also hold for closed monotone symplectic manifolds.

As an application, we show that the wrapped category of $\mathcal{O}(-k) \rightarrow \mathbb{C}\mathbb{P}^m$ is proper (cohomologically finite) for $1 \leq k \leq m$. For any monotone negative line bundle E over a closed monotone toric manifold B , we show that $SH^*(E) \neq 0$, $\mathcal{W}(E)$ is non-trivial, and E contains a non-displaceable monotone Lagrangian torus \mathcal{L} on which OC is non-zero.

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1. INTRODUCTION

1.1. **Preamble.** The wrapped Fukaya category $\mathcal{W}(E)$ of a non-compact convex symplectic manifold E , introduced in [23, 8] in the exact setup, is of increasing importance both in symplectic topology (topology of Stein manifolds [29], the nearby Lagrangian conjecture [5]) and in Homological Mirror Symmetry [6, 44]. The wrapped category is a version of the Fukaya category, which incorporates both closed Lagrangian submanifolds and non-compact Lagrangians modelled on a Legendrian cone at infinity, and which takes into account the dynamics of the Reeb flow at infinity. Incorporating non-compact Lagrangians typically leads to a substantial increase in complexity; the wrapped Floer cohomology of a cotangent fibre is isomorphic to the homology of the based loop space [1], in particular it is of infinite rank.

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Despite its increasing prominence, there are surprisingly few cases in which the wrapped category has been computed. Essentially complete descriptions are known for cotangent bundles [4] and for punctured spheres [6], whilst we have partial information for certain Stein manifolds (given as Lefschetz fibrations [29] or plumbings [10]), and some general structural information concerning the behaviour of the category under Weinstein surgery [13].

1.2. The generation criterion. Following Abouzaid [2], we study

$$\text{OC} : \text{HH}_*(\mathcal{W}(E)) \rightarrow SH^*(E), \quad (1.1)$$

the open-closed string map from Hochschild homology to symplectic cohomology. Abouzaid's generation criterion states that for an exact convex symplectic manifold E and a collection $L_1, \dots, L_n \in \text{Ob}(\mathcal{W}(E))$ of exact Lagrangian submanifolds, if OC restricted to the subcategory generated by L_1, \dots, L_n hits the unit $1 \in SH^*(E)$, then L_1, \dots, L_n split-generate the category $\mathcal{W}(E)$ (to *split-generate* means to generate up to taking idempotent projections).

For monotone convex symplectic manifolds E , one works instead with monotone orientable Lagrangian submanifolds, and examples naturally arise in the context of (non-compact) Fano toric varieties. In this context, holomorphic discs bounding Lagrangians and holomorphic bubbles in E (and therefore, the quantum cohomology $QH^*(E)$) play a prominent role, unlike the exact setup above where these cannot arise. This complicates the Floer theory, in fact we will explain in detail in 3.7 that the wrapped category is defined as a collection of mutually orthogonal A_∞ -categories $\mathcal{W}_\lambda(E)$ indexed by the spectrum of quantum multiplication by the first Chern class $c_1(TE)$ acting by quantum product on the quantum cohomology $QH^*(E)$ (where $\lambda \in \text{Spec}(c_1(TE))$ lies in the Novikov field Λ defined in 3.3). This is necessary because if two Lagrangians L_0, L_1 involved different $m_0 = \lambda$ eigenvalues, then the differential of the Lagrangian Floer "chain complex" $CF^*(L_0, L_1)$ would not square to zero.

Recall symplectic cohomology $SH^*(E)$ is a unital graded-commutative algebra that comes with a unital algebra homomorphism $c^* : QH^*(E) \rightarrow SH^*(E)$ which decomposes

$$c^* : QH^*(E) = \bigoplus_\lambda QH^*(E)_\lambda \rightarrow \bigoplus_\lambda SH^*(E)_\lambda \subset SH^*(E)$$

into generalized λ -eigenspaces under multiplication by $c_1(TE)$. We prove that these two "eigensummand decompositions" are related, imposing a restriction on the image of (1.1):

$$\text{OC} : \text{HH}_*(\mathcal{W}_\lambda(E)) \rightarrow SH^*(E)_\lambda \subset SH^*(E).$$

Therefore the generation criterion, as stated above in terms of hitting the unit, is essentially never going to hold. We will extend Abouzaid's generation criterion [2] as follows.

Theorem 1.1. *Let $\mathcal{C} \subset \mathcal{W}_\lambda(E)$ be a full subcategory. If $\text{OC}|_{\mathcal{C}}$ hits an invertible element in $SH^*(E)_\lambda$, then \mathcal{C} split-generates $\mathcal{W}_\lambda(E)$.*

This also holds for the compact Fukaya category $\mathcal{F}_\lambda(E)$ using $\text{OC}|_{\mathcal{C}} : \mathcal{F}(E)_\lambda \rightarrow QH^(E)_\lambda$, and it also holds if we replace E by a closed monotone symplectic manifold B .*

Example 1.2. *If $QH^*(E)_\lambda$ (respectively $SH^*(E)_\lambda$) is 1-dimensional, and thus a field, then it suffices to show that the image under OC of some Hochschild cycle is non-zero. For closed monotone B we prove that $\text{OC}^0[\text{pt}] \neq 0$ in 12.6, and for E convex and monotone, $\text{OC}^0[\text{pt}] \neq 0$ when $\lambda = m_0(L) \neq 0$ (see 12.7). However, one still needs to find an actual Lagrangian L with $m_0(L) = \lambda$ for which $[\text{pt}] \in CF^*(L, L)$ (respectively $CW^*(L, L)$) is a cycle, since otherwise $[\text{pt}]$ will not survive to cohomology. For toric manifolds we have a good supply of Lagrangian tori arising from critical points of the superpotential (see 12.1).*

Example 1.3. *Theorem 1.1 applies to the Clifford torus in $B = \mathbb{C}P^m$ taken with one of $1 + m$ choices of holonomy. It also applies to the lift of the Clifford torus to a sphere bundle of the complex line bundle $E = \mathcal{O}(-k) \rightarrow \mathbb{C}P^m$ where $1 \leq k \leq m$ (to ensure monotonicity),*

taken with one of $1 + m - k$ choices of holonomy. More generally, for any closed Fano variety B for which the superpotential W is Morse with distinct critical values, we show that $\mathcal{F}_\lambda(B)$ is split-generated by the monotone torus L_z with holonomy data associated to $z \in \text{Crit}(W)$.

The proof of Theorem 1.1 requires two ingredients. Firstly, we remark that Abouzaid’s generation criterion is a fairly general consequence of homological algebra once the wrapped Fukaya category can be constructed together with the necessary Floer machinery to define string maps and a coproduct (we recall this in Section 11.2). Thus our first contribution is a foundational one, namely we construct this machinery in the monotone setting. As discussed further in Section 1.4, this is a fairly substantial undertaking. In particular, we need a new double-telescope complex for the coproduct, which has not appeared before (see Section 10).

Secondly, we prove that $\text{HH}_*(\mathcal{W}_\lambda(E))$ is a module over $SH^*(E)$, and that $c_1(TE) - \lambda \text{Id}$ acts nilpotently on each element of $\text{HH}_*(\mathcal{W}_\lambda(E))$ (Section 9). Moreover we show that OC is an $SH^*(E)$ -module homomorphism (Section 8). Therefore, OC lands in $SH^*(E)_\lambda$ and hitting any invertible in that eigensummand is as good as hitting the identity of $SH^*(E)_\lambda$.

The module structure was constructed independently, in the exact case, by Ganatra [25], and the underlying QH^* -module structure is also used in forthcoming work of Abouzaid-Fukaya-Oh-Ohta-Ono [7] on generation in the closed case.

1.3. The acceleration functor. For convex symplectic manifolds E , one can also define the compact Fukaya category $\mathcal{F}(E)$ (Section 5). In the monotone setup, this means the objects are *closed* monotone orientable Lagrangian submanifolds (again one must in fact restrict these to all involve the same $m_0 = \lambda$ eigenvalue), and the morphism spaces are the Lagrangian Floer complexes $CF^*(L_0, L_1)$. When L_0, L_1 are transverse, this complex is generated by the intersections $L_0 \cap L_1$, and in the non-transverse case one introduces a *compactly-supported* Hamiltonian H so that the time one flow $\varphi_H^1(L_0)$ of L_0 becomes transverse to L_1 (we recall this in 5.2). Closed Lagrangians are also objects of the wrapped category, and the wrapped Floer complex $CW^*(L_0, L_1)$ can be loosely thought of (at least on cohomology) as a direct limit of Floer complexes $CF^*(\varphi_H^1(L_0), L_1)$ for a non-compactly supported Hamiltonian H which is linear at infinity, where in the direct limit we increase the slope of H via continuation maps. We remark that the substantial work required to build the wrapped category is precisely the chain-level construction of this direct limit procedure using a telescope construction that we adopt from Abouzaid-Seidel [8]. In fact, we will also need a new double telescope chain model for the direct limit when we come to the construction of the coproduct in Section 10.

Therefore one might expect there to be an A_∞ -functor, called *acceleration functor*

$$\mathcal{AF} : \mathcal{F}(E) \rightarrow \mathcal{W}(E)$$

which is the natural inclusion of objects, and the inclusion of the “slope zero” part $CF^*(L_0, L_1)$ into the above direct limit construction. However, a rigorous construction of this functor is again a non-trivial task, because “slope zero” Hamiltonians are not in general allowed in $\mathcal{W}(E)$ because for non-compact Lagrangians L_0, L_1 which have intersections at infinity, for “slope zero” H the compactly supported flow $\varphi_H^1(L_0)$ of L_0 cannot become transverse to L_1 .

Theorem 1.4. *For any monotone convex E , there is an acceleration functor which fits into commutative diagrams, which we call acceleration diagrams,*

$$\begin{array}{ccc} \text{HH}_*(\mathcal{F}_\lambda(E)) & \xrightarrow{\text{HH}_*(\mathcal{AF})} & \text{HH}_*(\mathcal{W}_\lambda(E)) \\ \text{OC} \downarrow & & \downarrow \text{OC} \\ QH^*(E)_\lambda & \xrightarrow{c^*} & SH^*(E)_\lambda \end{array} \qquad \begin{array}{ccc} \text{HH}^*(\mathcal{F}_\lambda(E)) & \xleftarrow{\text{HH}^*(\mathcal{AF})} & \text{HH}^*(\mathcal{W}_\lambda(E)) \\ \text{CO} \uparrow & & \uparrow \text{CO} \\ QH^*(E) & \xrightarrow{c^*} & SH^*(E) \end{array}$$

Moreover, the maps in the diagram are $QH^*(E)$ -module homomorphisms, and the maps in the diagram on the right are unital algebra homomorphisms.

Remark. HH^* is not in general functorial under A_∞ -functors. We construct $HH^*(\mathcal{AF})$ at the cochain level as an inclusion of a non-full subcategory, for which $HH^*(\mathcal{AF})$ makes sense.

Corollary 1.5. *Let E be a convex toric Fano variety. Let $z \in \text{Crit}(W)$ be a critical point of the superpotential with non-zero critical value $\lambda = W(z) \neq 0$. If the generalized eigensummand $QH^*(E)_\lambda$ is 1-dimensional then the Lagrangian torus L_z corresponding to z split-generates $\mathcal{F}_\lambda(E)$ and also $\mathcal{W}_\lambda(E)$, in particular $\mathcal{W}_\lambda(E)$ is then proper (cohomologically finite).*

We remark that just as for closed Fano toric varieties taken with the monotone toric symplectic form [22, Theorem 7.11] (compare also with the discussion in the non-compact case in [36, Section 7.9]) one can show that the Lagrangian tori L_z , with holonomy data, arising from critical points of the superpotential are in fact monotone. The proof of the Corollary uses the fact that the 0-part of the OC map for the compact category,

$$\text{OC}^0 : HF^*(L, L) \rightarrow QH^*(E)$$

is in fact non-zero for $L = L_z$, and therefore hits an invertible in the field $QH^*(E)_\lambda$ (using the assumption), so via the acceleration diagram, using that c^* is an algebra map, also

$$\text{OC}^0 : HW^*(L, L) \rightarrow SH^*(E)$$

for $L = \mathcal{AF}(L_z) = L_z$ will hit an invertible element in $SH^*(E)_\lambda$. In particular, observe that in general $c^* : QH^*(E)_\lambda \rightarrow SH^*(E)_\lambda$ sends unit to unit, since $c^* : QH^*(E) \rightarrow SH^*(E)$ does.

1.4. Linear versus quadratic Hamiltonians. There are two known constructions of the wrapped category of exact convex symplectic manifolds, which take advantage of Hamiltonian functions which have either linear slope at infinity or which grow quadratically at infinity. In either set-up one must deal with the fact that $CF^*(L, L)$ actually means $CF^*(\varphi_H(L), L)$ (using a Hamiltonian function to displace L slightly to achieve self-transversality) so products in Floer theory cannot be defined as $CF^*(L, L)^{\otimes k} \rightarrow CF^*(L, L)$. Rather, they are maps¹

$$CF^*(\varphi_{H_1}(L), L) \otimes CF^*(\varphi_{H_2}(L), L) \otimes \cdots \otimes CF^*(\varphi_{H_k}(L), L) \longrightarrow CF^*(\varphi_{H_1+\cdots+H_k}(L), L) \quad (1.2)$$

For non-compact L , the Floer complex on the right of (1.2) may not be isomorphic, at the chain level, to the factors on the left, even if all the H_i coincide. The Hamiltonians H_i are not compactly supported, and “large” Hamiltonians $H_1 + \cdots + H_k$ can cause $CF^*(\varphi_{H_1+\cdots+H_k}(L), L)$ to have many generators (the more we flow L , the more times it wraps around L , creating new intersections).

In the setting of linear slope Hamiltonians, as in [8], one defines the morphism groups via a “telescope” (homotopy direct limit) construction, $CW^*(L, L) = \bigoplus_k CF^*(\varphi_{kH}(L), L)[\mathbf{q}]$, involving arbitrarily large multiples of a fixed Hamiltonian of linear slope; whilst in the setting of quadratic Hamiltonians, one uses a rescaling trick (conjugating by a suitable multiple of the Liouville flow) to identify the output chain complex in (1.2) with the input complexes appearing on the left. Unfortunately, the rescaling trick is only available in the exact case, with no analogue for a general monotone convex symplectic manifold. Abouzaid [2] proved the generation criterion in the exact setup using quadratic Hamiltonians, and Ganatra uses the same setup in his discussion [25] of the module structure of the string maps.

The proof of each of the module structure of the OC-map and the generation criterion in the telescope model is surprisingly involved: the necessary modifications are numerous and non-trivial, and constitute one of the main contributions of this paper.

¹we are implicitly using canonical identifications of type $CF^*(\varphi_{H_j}(L), L) \equiv CF^*(\varphi_{H+H_j}(L), \varphi_H(L))$.

Although the quadratic Hamiltonian formalism as in [2, 25] bypasses the need in the exact setup for a telescope model, that formalism also has technical complications: a subtle perturbation scheme is needed to isolate the Hamiltonian orbits (see [25, Definition 4.10]) and the conformal rescaling scheme using the Liouville flow is rather involved. Our telescope model formalism also adapts to the exact setup, and avoids both these issues. Another benefit of our approach, is that it should adapt to general convex symplectic manifolds, without any monotonicity assumptions, at the cost of appealing to Kuranishi spaces or polyfolds.

Remark. *In the non-monotone setting, the $m_0 = \lambda$ values which index the categories $\mathcal{W}_\lambda(E)$ are not a priori tied to the eigenvalues of $c_1(TE) \in QH^*(E)$, see 3.7.*

1.5. Algebraic structure. In 8.1, we construct a unital algebra homomorphism

$$\psi : SH^*(E) \rightarrow \text{End}(\mathcal{W}(E)), \quad c \mapsto \psi_c$$

where $\text{End}(\mathcal{W}(E))$ denotes the bimodule endomorphisms of the diagonal $\mathcal{W}(E)$ -bimodule. Composing with the canonical map $\text{End}(\mathcal{W}(E)) \rightarrow \text{End}(\text{HH}_*(\mathcal{W}(E)))$, the ψ -string map defines an $SH^*(E)$ -module structure on $\text{HH}_*(\mathcal{W}(E))$.

Remark. *Ganatra [25] independently constructed the map ψ for exact convex symplectic manifolds using the quadratic Hamiltonian formalism of [2]: the construction is conceptually the same as ours, although technically quite different as explained in 1.4. Ganatra calls ψ the 2-pointed closed-open string map ${}_2\text{CO}$, due to its similarities with CO , and he calls the chain complex of bimodule pre-endomorphisms $\text{end}(\mathcal{W}(E))$ the 2-pointed Hochschild chain complex. The emphasis in [25] is to view the module structure as the cap-product action of $\text{CO}(c) \in \text{HH}^*(\mathcal{W}(E))$ on $\text{HH}_*(\mathcal{W}(E))$, whereas we have preferred to work directly with bimodule endomorphisms. The two approaches are equivalent due to the following diagram.*

Theorem 1.6. *There is a commutative diagram of unital algebra homomorphisms*

$$\begin{array}{ccccccc}
 QH^*(E) & \xrightarrow{c^*} & SH^*(E) & \xrightarrow{\psi} & \text{End}(\mathcal{W}(E)) & \xrightarrow{\text{HH}_*(\cdot)} & \text{End}(\text{HH}_*(\mathcal{W}(E))) \\
 & \searrow & & \searrow & \uparrow & & \nearrow \\
 & & & & \text{HH}^*(\mathcal{W}(E)) & & \\
 & \searrow & \text{CO} & & \uparrow \mu_{\mathcal{W}(E)} \circ_{\text{left}} & & \nearrow \text{cap product} \\
 & & & & & & \\
 & \searrow & \text{CO} & & & & \\
 & & & & & &
 \end{array}$$

using composition product for $\text{End}(\mathcal{W}(E))$ and $\text{End}(\text{HH}_*(\mathcal{W}(E)))$, cup product for $\text{HH}^*(\mathcal{W}(E))$.

The vertical map (defined in 2.16) is an isomorphism. The wrapped category is cohomologically unital, the cohomological units are $e_L = \text{CO}^0([E])(1) \in CW^*(L, L)$ where $\text{CO}^0(\cdot)(1) : QC^*(E) \rightarrow CW^*(L, L)$. Moreover $\text{CO}([E]) \in \text{HH}^*(\mathcal{W}(E))$ is a unit for the cup product.

Remark 1.7. *String maps were introduced by Seidel [39]. In the closed monotone case, Albers [11] defined maps $\tau : HF_*(H) \rightarrow HF^*(L, \varphi_H^1(L))$ and $\chi : HF_*(L, \varphi_H^1(L)) \rightarrow HF_*(H)$, identifiable with $\text{CO}^0 : QH^*(B) \rightarrow HF^*(L, L)$ and $\text{OC}^0 : HF^*(L, L) \rightarrow QH^*(B)$. Biran-Cornea [14] studied the $QH^*(B)$ -module structure of these maps. In the exact case, the $SH^*(E)$ -module structure of $HW^*(L, L)$ was studied by Ritter [34].*

We will also prove that the string maps are compatible with the module structure:

$$a \bullet \text{OC}(x) = \text{OC}(\psi_a(x)) \quad \text{CO}(a \bullet b) = \psi_a \circ \text{CO}(b) = \text{CO}(a) * \text{CO}(b)$$

where \bullet is the $\mu_{SH^*(E)}^2$ -product, $a, b \in SH^*(E)$, $x \in \text{HH}_*(\mathcal{W}(E))$, \circ is defined in 2.3, and the cup product $*$ on $\text{HH}^*(\mathcal{W}(E))$ is defined in 2.14.

1.6. Eigensummand decomposition of OC and CO.

Theorem 1.8. *The open-closed string map $\text{OC} : \text{HH}_*(\mathcal{W}_\lambda(E)) \rightarrow SH^*(E)$ must land in the generalized λ -eigenspace $SH^*(E)_\lambda$ (and analogously for $\mathcal{F}_\lambda(E)$ and $QH^*(E)$).*

The closed-open string map $\text{CO} : SH^(E)_\mu \rightarrow \text{HH}^*(\mathcal{W}_\lambda(E))$, restricted to the generalized μ -eigenspace, will vanish if $\mu \neq \lambda$ (and analogously for $\text{CO} : QH^*(E)_\mu \rightarrow \text{HH}^*(\mathcal{F}_\lambda(E))$).*

At the heart of this result is an observation due to Kontsevich, Seidel and Auroux [12]. Namely, $\text{CO}^0 : QH^*(E) \rightarrow HF^*(L, L)$ maps $c_1(TE)$ to the unit $\text{CO}^0([E])$ rescaled by the $m_0(L) = \lambda$ value (see 9.2). So the module action of $c_1(TE) - \lambda \text{Id}$ on $HF^*(L, L)$ vanishes as $\text{CO}^0(c_1(TE) - \lambda \text{Id}) = 0 \in HF^*(L, L)$. This implies that the bimodule endomorphism $\psi_{c_1(TE) - \lambda \text{Id}}$ (after subtracting an exact term, since now we work at the chain level) has vanishing $0|0$ -part. It follows that $\psi_{c_1(TE) - \lambda \text{Id}}$ acts nilpotently on any given word in $\text{CC}_*(\mathcal{M})$ (composing many times, eventually one of the maps in the composite must be taken with $0|0$ -part, as the word has finite length, see 2.12). This implies that any given cycle in $\text{CC}_*(\mathcal{M})$ is killed by a sufficiently large quantum power of $c_1(TE) - \lambda \text{Id}$, therefore OC evaluated on the cycle must land in the λ -generalized eigensummand of $SH^*(E)$ (respectively $QH^*(E)$). One can dualize this argument (more precisely, this requires a subtle filtration argument) to show that CO must also respect eigensummands.

1.7. Applications. The total space of the line bundle $\mathcal{O}(-k) \rightarrow \mathbb{C}\mathbb{P}^m$ is monotone whenever $1 \leq k \leq m$. The ring $SH^*(\mathcal{O}(-k))$ was computed explicitly in [35, 36], and we will show that

$$SH^*(\mathcal{O}(-k)) \cong \text{Jac}(W)$$

recovers the Jacobian ring of the superpotential for $\mathcal{O}_{\mathbb{P}^m}(-k)$. This is a proof of closed-string mirror symmetry in this example. Note that unlike the case of closed Fano toric manifolds, the Jacobian ring does not recover the quantum cohomology $QH^*(\mathcal{O}(-k))$ (see 12.5).

Then, making heavy use of the module structure and compatibility with eigenvalue splittings, we compute the open-closed string map explicitly, and prove that the wrapped category is proper (cohomologically finite), split-generated by closed Lagrangian submanifolds. By contrast, in all previous situations in which the wrapped Fukaya category has been computed (punctured spheres, certain plumbings, cotangent bundles), it either vanishes or is cohomologically infinite. The special case of $\mathcal{O}(-1) \rightarrow \mathbb{P}^1$ had been studied previously [46, Sec.4.4].

Local Calabi-Yau's (total spaces of canonical bundles over Fano varieties) are well-known testing grounds for many aspects of mirror symmetry, with rich symplectic topology [43, 15]. Monotone negative line bundles form another class of "local symplectic manifolds", which arise naturally in symplectic birational geometry, cf. [46]. As remarked above, whilst in known computations in the Stein case the wrapped category $\mathcal{W}(E)$ either vanishes or is extremely complicated, the line bundles $\mathcal{O}(-k) \rightarrow \mathbb{P}^n$ show qualitatively different behaviour.

Making use of the module structure, the compatibility with eigenvalue splittings and a Galois action on Floer cohomology, we prove:

Theorem 1.9.

For any monotone negative line bundle $E \rightarrow B$ over a closed monotone toric manifold B ,

- (1) *the symplectic cohomology $SH^*(E)$ is non-zero;*
- (2) *the wrapped category $\mathcal{W}(E)$ is non-trivial;*
- (3) *in particular, there is a non-displaceable Lagrangian torus $L \subset E$ with holonomy data such that $HF^*(L, L) \cong HW^*(L, L) \neq 0$, on which OC^0 is non-zero.*

The proof of the theorem involves showing that critical points $p \in \text{Crit}(W_B)$ of the superpotential of the base with non-zero critical value $W_B(p) \neq 0$ give rise to critical points of the superpotential W_E , which in turn give rise to Lagrangian tori $L \subset E$ with holonomy data

such that $HF^*(L, L) \neq 0$. In particular, these exist since a recent result of Galkin [24] implies that W_B has at least one critical point p with $W_B(p) \neq 0$. The $QH^*(E)$ -module structure on $HF^*(L, L) \neq 0$ implies that $c_1(TE)$ is non-nilpotent in $QH^*(E)$. By Ritter [35],

$$SH^*(E) \cong QH^*(E)/(\text{generalized 0-eigenspace of quantum multiplication by } \pi^*c_1(E))$$

for any monotone negative line bundle $\pi : E \rightarrow B$. As $c_1(TE)$ and $\pi^*c_1(E)$ are positively proportional, it follows that $\pi^*c_1(E)$ is also not nilpotent so $SH^*(E) \neq 0$. By the argument explained at the end of 1.3, $OC^0 : HF^*(L, L) \rightarrow QH^*(E)$ is non-zero and therefore by the acceleration diagram in Theorem 1.4 also $OC^0 : HW^*(L, L) \rightarrow SH^*(E)$ is non-zero. Therefore the category $\mathcal{W}_\lambda(E)$ containing L is cohomologically non-trivial.

It is not hard to envisage other applications. In a series of papers, Abouzaid [3, 4, 5] proved that the cotangent fibre T_q^*Q generates the wrapped category $\mathcal{W}(T^*Q)$ of the cotangent bundle of a closed manifold equipped with its canonical exact symplectic structure $d\theta$. From a dynamical point of view, there is a great deal of interest in the convex non-exact manifolds defined by cotangent bundles equipped with magnetic forms, i.e. by $(T^*Q, d\theta + \pi^*\sigma)$ where $\sigma \in \Omega^2(Q)$ is some closed but not necessarily exact form pulled back from the base. It seems likely that our results can be used to show a cotangent fibre still generates the wrapped category after turning on the magnetic potential.

1.8. Remarks and Acknowledgements.

Remark about signs and characteristic. Although our constructions rely on working in characteristic zero, for simplicity of exposition we will not discuss signs at all in this paper. A very careful treatment of signs in Floer theory is carried out in Seidel [41], Abouzaid-Seidel [8], Abouzaid [2], Ganatra [25] and Ritter [34]. In particular, in several constructions we added a remark about a “sign correction factor”. These are not strictly meaningful in the absence of a detailed discussion of signs, as they are only defined relative to a certain convention for identifying orientation lines of the inputs with those of the outputs, for which a number of choices exist. So these remarks are understood to follow the conventions of [41, 8, 2].

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Current version. *Since the first version [37] of this paper appeared (Jan. 2012) two papers have appeared with significant conceptual overlap: (1) Ganatra’s PhD and preprint [25] studies the wrapped Fukaya category and the string maps in the exact setup, (2) Sheridan’s preprint [45] studies the Fukaya category of closed monotone symplectic manifolds. Still to appear, the work in progress by Abouzaid-Fukaya-Oh-Ohta-Ono [7] will study the Fukaya category of closed toric manifolds without monotonicity assumptions. A large part of an earlier version [37] constructed a commutative diagram relating the Fukaya category of a closed monotone symplectic manifold B and of a convex (non-compact) monotone symplectic manifold E related by a Lagrangian correspondence $\Gamma \subset \overline{B} \times E$. This part was removed to keep this paper of reasonable length; we hope this material will appear elsewhere.*

2. HOCHSCHILD (CO)HOMOLOGY OF A_∞ -CATEGORIES

We summarize the general algebraic constructions that we need in the paper. For ease of reading, we accompany the A_∞ -relations of this Section with pictures of the corresponding natural bubbling phenomena arising in the context of Fukaya categories. For a thorough

treatment of the A_∞ -module theory we refer to Seidel [40], Abouzaid [2], and Ganatra [25].

Remark. *We restrict ourselves to non-curved A_∞ -categories.*

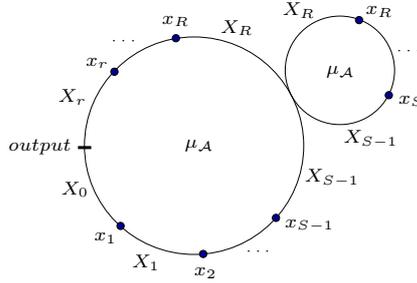
2.1. **A_∞ -categories, and grading conventions.** Recall that an A_∞ -category \mathcal{A} is: a collection of objects $\text{Ob}(\mathcal{A})$; graded vector spaces $\text{hom}_{\mathcal{A}}(X_0, X_1)$ for all $X_0, X_1 \in \text{Ob}(\mathcal{A})$, working over a chosen ground field; composition maps

$$\mu_{\mathcal{A}}^r : \mathcal{A}(X_r, \dots, X_0) \rightarrow \mathcal{A}(X_r, X_0),$$

where we abbreviate

$$\mathcal{A}(X_r, X_{r-1}, \dots, X_0) = \text{hom}_{\mathcal{A}}(X_{r-1}, X_r) \otimes \text{hom}_{\mathcal{A}}(X_{r-2}, X_{r-1}) \otimes \dots \otimes \text{hom}_{\mathcal{A}}(X_0, X_1)$$

for $r \geq 1$ (for $r = 0$ define it to equal the ground field; and we will typically denote generators by $x_r \otimes \dots \otimes x_1$); and finally the $\mu_{\mathcal{A}}^r$ must satisfy the following A_∞ -relations.



$$\sum (-1)^{\sigma_1^{S-1}} \mu_{\mathcal{A}}^r(x_r \otimes \dots \otimes x_{R+1}) \otimes \mu_{\mathcal{A}}^{R-S+1}(x_R, \dots, x_S), x_{S-1}, \dots, x_1) = 0$$

summing over all obvious choices of R, S , and we explain the sign and the gradings below. The A_∞ -relations imply that $\mu_{\mathcal{A}}^1$ is a differential and the multiplication $\mu_{\mathcal{A}}^2$ is a chain map.

Let $|x|$ be the grading of a morphism in \mathcal{A} , and define $\|x\| = |x| - 1$ the reduced grading. Then grading signs are determined by the convention that the compositions $\mu_{\mathcal{A}}^r$ are acting from the right with degree 1. Abbreviate

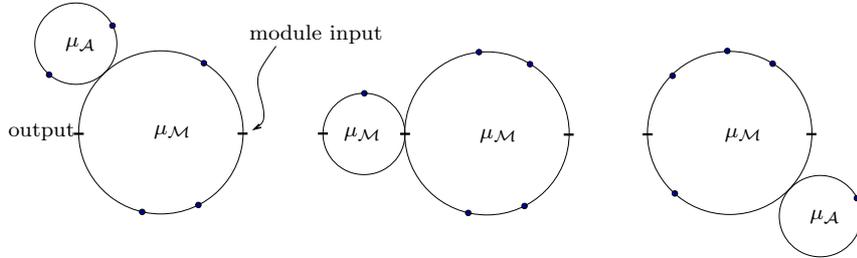
$$\sigma_i^j = \sigma(x)_i^j = \sum_{\ell=i}^j \|x_\ell\|,$$

which is the reduced grading of $x_j \otimes x_{j-1} \otimes \dots \otimes x_i$. The second expression emphasizes the letter x which is used to denote the variables indexed by $i, i+1, \dots, j$.

2.2. **Bimodules.** An \mathcal{A} -bimodule \mathcal{M} is a collection of graded vector spaces $\mathcal{M}(X, Y)$ for any objects X, Y of \mathcal{A} , together with multi-linear maps

$$\mu_{\mathcal{M}}^{r|s} : \mathcal{A}(X_r, \dots, X_0) \otimes \mathcal{M}(X_0, Y_0) \otimes \mathcal{A}(Y_0, \dots, Y_s) \rightarrow \mathcal{M}(X_r, Y_s)$$

for $r, s \geq 0$, and any objects X_i, Y_j of \mathcal{A} , satisfying the following A_∞ -relations.



$$\begin{aligned}
0 &= \sum (-1)^{\star} \mu_{\mathcal{M}}^{r-R+S|s} (x_r, \dots, x_{R+1}, \mu_{\mathcal{A}}^{R-S+1}(x_R, \dots, x_S), x_{S-1}, \dots, x_1, \underline{m}, y_1, \dots, y_s) + \\
&+ \sum (-1)^{\diamond} \mu_{\mathcal{M}}^{r-R|s-S} (x_r, \dots, x_{R+1}, \mu_{\mathcal{M}}^{R|S}(x_R, \dots, x_1, \underline{m}, y_1, \dots, y_S), y_{S+1}, \dots, y_s) + \\
&+ \sum (-1)^{\diamond} \mu_{\mathcal{M}}^{r|s-R+S} (x_r, \dots, x_1, \underline{m}, y_1, \dots, y_{R-1}, \mu_{\mathcal{A}}^{R-S+1}(y_R, \dots, y_S), y_{S+1}, \dots, y_s)
\end{aligned}$$

summing over all obvious choices of R, S , and where the signs are

$$\diamond = \sigma(y)_{S+1}^s \quad \star = \sigma(y)_1^s + \deg(\underline{m}) + \sigma(x)_1^{S-1}. \quad (2.1)$$

Note: generators of the domain of $\mu_{\mathcal{M}}^{r|s}$ are written $x_r \otimes \dots \otimes x_1 \otimes \underline{m} \otimes y_1 \otimes \dots \otimes y_s$, where we underline the bimodule element.

The above equations imply in particular that $\mu_{\mathcal{M}}^{0|0}$ is a differential, and that the left multiplication $\mu_{\mathcal{M}}^{1|0}$ and right multiplication $\mu_{\mathcal{M}}^{0|1}$ are chain maps.

Example. The diagonal bimodule $\mathcal{M} = \mathcal{A}$ has $\mathcal{M}(X_0, Y_0) = \text{hom}_{\mathcal{A}}(Y_0, X_0) = \mathcal{A}(X_0, Y_0)$ and composition maps $\mu_{\mathcal{M}}^{r|s}(x_r, \dots, x_1, \underline{m}, y_1, \dots, y_s) = (-1)^{1+\sigma(y)_1^s} \mu_{\mathcal{A}}^{r+1+s}(x_r, \dots, x_1, \underline{m}, y_1, \dots, y_s)$.

2.3. Convenient abbreviations: \circ and μ . The abbreviation $A \circ B$ will mean: compose in all possible ways, respecting the Koszul sign convention (operators act from the right)

$$A \circ B(x_n, \dots, x_1) = \sum (-1)^{\|B\| \sigma_1^{s-1}} A^{n-r+s}(x_n, \dots, x_{r+1}, B^{r-s+1}(x_r, \dots, x_s), x_{s-1}, \dots, x_1), \quad (2.2)$$

where σ_1^{s-1} uses reduced degrees unless one of the x_j is a module input (then the unreduced $\deg(x_j)$ is used) and $\|B\|$ is the degree of B (to be specified). In particular, the sum includes insertions of B^0 with no inputs, which means $B^0 : \mathbb{K} \rightarrow \mathcal{A}(X, X)$ evaluated at $1 \in \mathbb{K}$.

In the context of operators which by definition must receive a module input, it is understood that $A \circ B(x_r, \dots, x_1, \underline{m}, y_1, \dots, y_s)$ means

$$\sum (-1)^{\|B\| \sigma_{S+1}^s} A^{r-R|s-S}(x_r, \dots, x_{R+1}, \underline{B^{R|S}(x_R, \dots, x_1, \underline{m}, y_1, \dots, y_S)}, x_{S+1}, \dots, y_s). \quad (2.3)$$

In particular, in this case the lowest order insertion of B is $B^{0|0}(\underline{m})$ which receives one input.

The operator $B = \mu$ (with $\|B\| = 1$) will mean applying both $\mu_{\mathcal{M}}^{r|s}$ and $\mu_{\mathcal{A}}^r$ as appropriate: that is depending on whether or not a module input is present in the sub-word that B receives.

Example. The A_{∞} -relations for an A_{∞} -category \mathcal{A} are: $\mu_{\mathcal{A}} \circ \mu_{\mathcal{A}} = 0$.

The A_{∞} -relations for an \mathcal{A} -bimodule \mathcal{M} are: $\mu_{\mathcal{M}} \circ \mu = 0$. (or: $\mu \circ \mu = 0$.)

2.4. Pre-morphisms of bimodules $\text{hom}(\mathcal{M}, \mathcal{N})$. A pre-morphism $f : \mathcal{M} \rightarrow \mathcal{N}$ of \mathcal{A} -bimodules is a collection of multi-linear maps (for $r, s \geq 0$):

$$f^{r|s} : \mathcal{A}(X_r, \dots, X_0) \otimes \mathcal{M}(X_0, Y_0) \otimes \mathcal{A}(Y_0, \dots, Y_s) \rightarrow \mathcal{N}(X_r, Y_s).$$

Denote $\text{hom}(\mathcal{M}, \mathcal{N})$ the collection of these; $\text{end}(\mathcal{M}) = \text{hom}(\mathcal{M}, \mathcal{M})$ the pre-endomorphisms.

2.5. Yoneda composition product of pre-morphisms. The Yoneda product is the composition of pre-morphisms using (2.3),

$$\circ : \text{hom}(\mathcal{M}_2, \mathcal{M}_3) \otimes \text{hom}(\mathcal{M}_1, \mathcal{M}_2) \rightarrow \text{hom}(\mathcal{M}_1, \mathcal{M}_3).$$

2.6. The unit pre-morphism. The unit $1_{\mathcal{M}} \in \text{end}(\mathcal{M})$ for the ring $(\text{end}(\mathcal{M}), \circ)$ is

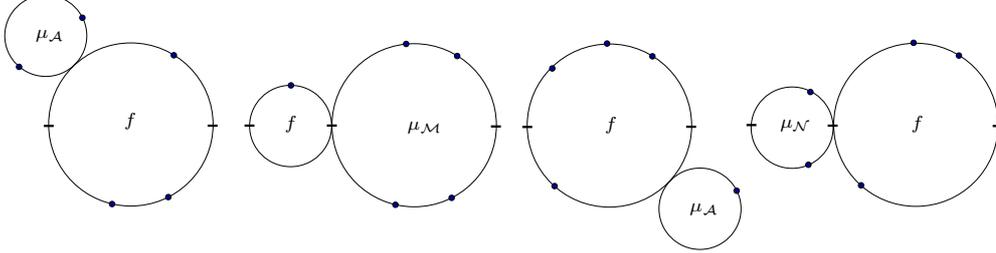
$$1_{\mathcal{M}}^{0|0} = \text{identity} : \mathcal{M}(X_0, Y_0) \rightarrow \mathcal{M}(X_0, Y_0), \quad 1_{\mathcal{M}}^{r|s} = 0 \text{ if } r \geq 1 \text{ or } s \geq 1.$$

2.7. Differential of a pre-morphism. For $f \in \text{hom}(\mathcal{M}, \mathcal{N})$, define

$$\delta f = \mu_{\mathcal{N}} \circ f - (-1)^{\|\mu\| \|f\|} f \circ \mu$$

using the conventions in 2.3 where $\|\mu\| = 1$, $\|f\| = \deg(f)$. One can now verify that the category of bimodules is a dg-category.

2.8. Morphisms of bimodules. A pre-morphism $f : \mathcal{M} \rightarrow \mathcal{N}$ is a morphism if $\delta f = 0$, equivalently $f \in \ker \delta$. Explicitly the A_∞ -relations $\delta f = 0$ are written below.



$$\begin{aligned} & \sum (-1)^* f^{r-R+S|s} (x_r, \dots, x_{R+1}, \mu_{\mathcal{A}}^{R-S+1}(x_R, \dots, x_S), x_{S-1}, \dots, x_1, \underline{m}, y_1, \dots, y_s) + \\ & + \sum (-1)^\diamond f^{r-R|s-S} (x_r, \dots, x_{R+1}, \mu_{\mathcal{M}}^{R|S}(x_R, \dots, x_1, \underline{m}, y_1, \dots, y_S), y_{S+1}, \dots, y_s) + \\ & + \sum (-1)^\diamond f^{r|s-S+R} (x_r, \dots, x_1, \underline{m}, y_1, \dots, y_{R-1}, \mu_{\mathcal{A}}^{S-R+1}(y_R, \dots, y_S), y_{S+1}, \dots, y_s) \\ & = (-1)^{\deg(f)} \sum (-1)^\diamond \mu_{\mathcal{N}}^{r-R|s-S} (x_r, \dots, x_{R+1}, f^{R|S}(x_R, \dots, x_1, \underline{m}, y_1, \dots, y_S), y_{S+1}, \dots, y_s) \end{aligned}$$

summing over the obvious choices of R, S , and where the signs are as in (2.1).

Example. For $\mathcal{M} = \mathcal{N}$, the unit pre-morphism $1_{\mathcal{M}}$ from 2.5 satisfies $\delta 1_{\mathcal{M}} = 0$.

The vector space of morphisms, modulo boundaries of pre-morphisms, will be denoted

$$\mathrm{Hom}(\mathcal{M}, \mathcal{N}) = \ker \delta / \mathrm{im} \delta \quad \mathrm{End}(\mathcal{M}) = \mathrm{Hom}(\mathcal{M}, \mathcal{M}).$$

2.9. Shift of bimodules. The shift operation $\mathcal{M}[1]$ on a bimodule \mathcal{M} is defined by shifting the degrees down by 1, so $\deg_{\mathcal{M}[1]}(\underline{m}) = \deg_{\mathcal{M}}(\underline{m}) - 1$, but also changing signs:

$$\mu_{\mathcal{M}[1]}^{r|s}(x_r, \dots, x_1, \underline{m}, y_1, \dots, y_s) = (-1)^{1+\sigma(y)_1^s} \mu_{\mathcal{M}}^{r|s}(x_r, \dots, x_1, \underline{m}, y_1, \dots, y_s)$$

That sign ensures that $\mathrm{Hom}^n(\mathcal{M}, \mathcal{N}[1]) = \mathrm{Hom}^{n+1}(\mathcal{M}, \mathcal{N})$ are identifiable chain complexes. On the other hand, for $\mathrm{Hom}^n(\mathcal{M}[1], \mathcal{N}) \cong \mathrm{Hom}^{n-1}(\mathcal{M}, \mathcal{N})$, $f \mapsto f[1]$ we require a sign change: $f[1]^{r|s}(x_r, \dots, x_1, \underline{m}, y_1, \dots, y_s) = (-1)^{\sigma(y)_1^s} f^{r|s}(x_r, \dots, x_1, \underline{m}, y_1, \dots, y_s)$.

Example. $\mathcal{M} = \mathcal{A}[1]$ has composition maps $\mu_{\mathcal{A}[1]}^{r|s} = \mu_{\mathcal{A}}^{r+1+s}$.

2.10. Cohomology, products and units. For an A_∞ -category \mathcal{A} , the cohomology category $H^*(\mathcal{A})$ has morphism spaces equal to the cohomology of the chain complex $(\mathcal{A}(X_0, X_1), \mu_{\mathcal{A}}^1)$ with (associative) composition maps equal to the μ^2 -product (up to sign):

$$[x] \cdot [y] = (-1)^{|y|} \mu_{\mathcal{A}}^2(x, y) = (-1)^{\|y\|+1} \mu_{\mathcal{A}}^2(x, y).$$

The cohomology $H^*(\mathcal{M})$ of a bimodule \mathcal{M} is the cohomology of the chain complexes $(\mathcal{M}(X, Y), \mu_{\mathcal{M}}^{0|0})$. A morphism is a quasi-isomorphism if the map on cohomology is an isomorphism. A feature of the A_∞ -theory is that quasi-isomorphisms are always invertible up to homotopy (this applies to A_∞ -categories and to bimodules).

An A_∞ -category \mathcal{A} is cohomologically unital (c-unital) if every object $X \in \mathrm{Ob}(\mathcal{A})$ has an endomorphism $e_X \in \mathrm{hom}_{\mathcal{A}}^0(X, X)$ which is a unit for $H^*(\mathcal{A})$ with the above product (note the degree $|e_X| = 0$, so $\|e_X\| = -1$). These are called cohomological units. This is weaker than strict-units $e \in \mathrm{hom}_{\mathcal{A}}^0(X, X)$, which would require $\mu_{\mathcal{A}}^1(e) = 0$, $(-1)^{\|y\|+1} \mu_{\mathcal{A}}^2(e, y) = y = \mu_{\mathcal{A}}^2(y, e)$ for all y , and higher $\mu_{\mathcal{A}}^k = 0$ if we insert e in any slot. Up to quasi-isomorphism, one can replace a cohomologically unital A_∞ -category by a strictly unital one and only consider strictly unital functors and strictly unital A_∞ -transformations. A bimodule \mathcal{M} is cohomologically unital if $(-1)^{\deg(\cdot)+1} \mu_{\mathcal{M}}^{1|0}(e_X, \cdot)$ and $\mu_{\mathcal{M}}^{0|1}(\cdot, e_X)$ are identity maps on cohomology. This is weaker than strict-unitality, which requires $(-1)^{\deg(\underline{m})+1} \mu_{\mathcal{M}}^{1|0}(e, \underline{m}) = \underline{m} = \mu_{\mathcal{M}}^{0|1}(\underline{m}, e)$ and

higher $\mu_{\mathcal{M}}^{r|s} = 0$ if we insert e in any slot. For strictly unital A_∞ -algebras, the dg category of strictly unital bimodules is quasi-equivalent to that of all bimodules.

2.11. Hochschild homology. Let \mathcal{M} be an \mathcal{A} -bimodule. Its cyclic bar complex in dimension $n \geq 0$ is the vector space

$$\mathrm{CC}_n(\mathcal{A}, \mathcal{M}) = \bigoplus \mathcal{M}(X_0, X_n) \otimes \mathcal{A}(X_n, \dots, X_0)$$

summing over all $X_j \in \mathrm{Ob}(\mathcal{A})$. A typical generator will be denoted $\underline{m} \otimes x_n \otimes \dots \otimes x_1$. We underline the element belonging to the bimodule. The grading for $\mathrm{CC}_n(\mathcal{A}, \mathcal{M})$ is:

$$\|\underline{m} \otimes x_n \otimes \dots \otimes x_1\| = \mathrm{deg}(\underline{m}) + \sigma_1^n.$$

The bar differential $b : \mathrm{CC}_n(\mathcal{A}, \mathcal{M}) \rightarrow \mathrm{CC}_{n-1}(\mathcal{A}, \mathcal{M})$ on a generator $\underline{m} \otimes x_n \otimes \dots \otimes x_1$ is

$$\begin{aligned} & \sum (-1)^{\sigma_1^{s-1}} \underline{m} \otimes x_n \otimes \dots \otimes x_{r+1} \otimes \mu_{\mathcal{A}}^{r-s+1}(x_r, \dots, x_s) \otimes x_{s-1} \otimes \dots \otimes x_1 + \\ & \sum (-1)^{\dagger} \underline{\mu_{\mathcal{M}}^{s|n-r+1}(x_s, \dots, x_1, \underline{m}, x_n, \dots, x_r)} \otimes x_{r-1} \otimes \dots \otimes x_{s+1} \end{aligned}$$

summing over all obvious choices of r, s , where $\dagger = \sigma_1^s(\mathrm{deg}(\underline{m}) + \sigma_{s+1}^n) + \sigma_{s+1}^{r-1}$.

The Hochschild homology is the resulting homology:

$$\mathrm{HH}_n(\mathcal{A}, \mathcal{M}) = H_n(\mathrm{CC}_*(\mathcal{A}, \mathcal{M}); b).$$

Example. $\mathrm{HH}_n(\mathcal{A}) = H_n(\mathrm{CC}_*(\mathcal{A}, \mathcal{A}); b)$ is the Hochschild homology of the diagonal (see 2.2).

2.12. Morphisms of bimodules induce chain maps on CC_* .

Lemma 2.1. *A morphism of \mathcal{A} -bimodules $f : \mathcal{M} \rightarrow \mathcal{N}$ induces a chain map*

$$\mathrm{CC}_*(f) : \mathrm{CC}_*(\mathcal{A}, \mathcal{M}) \rightarrow \mathrm{CC}_*(\mathcal{A}, \mathcal{N}).$$

Proof. Define

$$\mathrm{CC}_n(f) : \mathcal{M}(X_0, X_n) \otimes \mathcal{A}(X_n, \dots, X_0) \rightarrow \mathrm{CC}_*(\mathcal{A}, \mathcal{N})$$

by sending $\underline{m} \otimes x_n \otimes \dots \otimes x_1$ to

$$\sum (-1)^{\diamond} \underline{f^{r|n-s+1}(x_r, \dots, x_1, \underline{m}, x_n, \dots, x_s)} \otimes x_{s-1} \otimes \dots \otimes x_{r+1}$$

where $\diamond = \sigma_1^r(\mathrm{deg}(\underline{m}) + \sigma_{r+1}^n) + \mathrm{deg}(f)\sigma_{r+1}^{s-1}$. It is a chain map by Lemma 2.2(1). \square

Lemma 2.2.

- (1) $\mathrm{CC}_*(\delta f) = b \circ \mathrm{CC}_*(f) - (-1)^{\mathrm{deg}(f)} \mathrm{CC}_*(f) \circ b$.
- (2) $\mathrm{CC}_*(f \circ g) = \mathrm{CC}_*(f) \circ \mathrm{CC}_*(g)$.

Proof. $\mathrm{CC}_*(\delta f) = \mathrm{CC}_*(\mu_{\mathcal{N}} \circ f) - (-1)^{\|f\|} \mathrm{CC}_*(f \circ \mu) = b \circ \mathrm{CC}_*(f) - (-1)^{\|f\|} \mathrm{CC}_*(f) \circ b$. The second claim follows immediately from definitions. \square

Corollary 2.3. *The following is a unital algebra homomorphism, using composition products,*

$$\mathrm{End}(\mathcal{M}) \rightarrow \mathrm{End}(\mathrm{HH}_*(\mathcal{M})), \quad f \mapsto \mathrm{HH}_*(f). \quad (2.4)$$

Thus $\mathrm{HH}_(\mathcal{M})$ is a module over the unital ring of bimodule endomorphisms.*

Proof. By Lemma 2.2, if $\delta f = 0$ then $\mathrm{CC}_*(f)$ is a chain map (up to the sign $(-1)^{\mathrm{deg}(f)}$), so $\mathrm{HH}_*(f) \in \mathrm{End}(\mathrm{HH}_*(\mathcal{M}))$ is well-defined for $f \in \ker \delta$. The map (2.4) is well-defined by Lemma 2.2(1): $\mathrm{HH}_*(\delta f) = 0 \in \mathrm{End}(\mathrm{HH}_*(\mathcal{M}))$. Moreover, (2.4) preserves products by Lemma 2.2(2). Unitality holds already at the chain level: $1_{\mathcal{M}} \mapsto \mathrm{CC}_*(1_{\mathcal{M}}) = \mathrm{id}_{\mathrm{CC}_*(\mathcal{M})}$. \square

Lemma 2.4. *If $\psi : \mathcal{M} \rightarrow \mathcal{M}$ is a bimodule endomorphism with $\psi^{0|0} = 0$ (at the chain level), then $\mathrm{HH}_*(\psi)$ acts nilpotently on $\mathrm{HH}_*(\mathcal{M})$ (meaning, given $c \in \mathrm{HH}_*(\mathcal{M})$, there is an integer N possibly depending on c such that $\mathrm{HH}_*(\psi)^N(c) = 0$).*

Proof. An element of type $\underline{m} \otimes x_n \otimes \cdots \otimes x_1$ is called a length n word. We claim that $\psi^N = \psi \circ \cdots \circ \psi$ (composing N times at the chain level) must vanish on a word of length less than N . Indeed, consider a word $x_r \otimes \cdots \otimes x_1 \otimes \underline{m} \otimes y_1 \otimes \cdots \otimes y_s$ with $r + s < N$. If one of the ψ factors in ψ^N only receives a module input, then we get zero since $\psi^{0|0} = 0$ by assumption. If all ψ factors receive more than just a module input, then by definition of composition each ψ factor eats at least one non-module input from the $r + s$ available inputs $x_r, \dots, x_1, y_1, \dots, y_s$. But by the pigeonhole principle, there are not enough available inputs since $N > r + s$.

Given a cycle $c \in \text{CC}_*(\mathcal{M})$, let N be larger than the length of any word occurring in the finite linear combination comprising c to deduce $\text{CC}_*(\psi^N)(c) = 0$ at the chain level. Thus $\text{HH}_*(\psi)^N[c] = \text{HH}_*(\psi^N)[c] = 0$ by Corollary 2.3. \square

2.13. Hochschild cohomology. The Hochschild cochain complex is

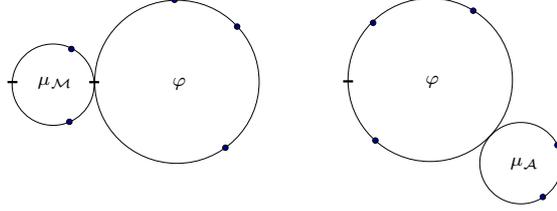
$$\begin{aligned} \text{CC}^d(\mathcal{A}, \mathcal{M}) &= \prod_{n \geq 0} \prod \text{Hom}(\mathcal{A}(X_n, \dots, X_0), \mathcal{M}(X_n, X_0)[d]) \\ &= \prod_{n \geq 0} \prod \text{Hom}^d(\mathcal{A}(X_n, \dots, X_0), \mathcal{M}(X_n, X_0)) \end{aligned}$$

where the unlabelled product is over all $X_j \in \text{Ob}(\mathcal{A})$, and we used reduced gradings for \mathcal{A} but not for \mathcal{M} (for non-reduced gradings the shift is $[d - n]$ instead of $[d]$).

We write $\varphi = (\varphi^n)_{n \geq 0}$ for a typical element of $\text{CC}^d(\mathcal{A}, \mathcal{M})$, so $\varphi^n(x_n, \dots, x_1) \in \mathcal{M}(X_n, X_0)$ for $x_j \in \text{hom}_{\mathcal{A}}(X_{j-1}, X_j)$. Using the notation 2.5, the differential is defined by

$$\delta\varphi = \mu_{\mathcal{M}} \circ \varphi - (-1)^{\|\mu_{\mathcal{A}}\| \|\varphi\|} \varphi \circ \mu_{\mathcal{A}}$$

where $\|\mu_{\mathcal{A}}\| = 1$, $\|\varphi\| = d$. Explicitly $\delta\varphi(x_n, \dots, x_1) \in \mathcal{M}(X_n, X_0)$ is as follows.



$$\begin{aligned} &\sum (-1)^d \sigma_1^{s-1} \mu_{\mathcal{M}}^{n-r|s-1}(x_n, \dots, x_{r+1}, \overline{\varphi^{r-s+1}(x_r, \dots, x_s)}, x_{s-1}, \dots, x_1) - \\ &- (-1)^d \sum (-1)^{\|\mu_{\mathcal{A}}\| \cdot \sigma_1^{s-1}} \varphi^{n-r+s}(x_n, \dots, x_{r+1}, \overline{\mu_{\mathcal{A}}^{r-s+1}(x_r, \dots, x_s)}, x_{s-1}, \dots, x_1). \end{aligned}$$

Remark. When $\mathcal{M} = \mathcal{N} = \mathcal{A}$, replacing all μ above by $\psi \in \text{CC}^*(\mathcal{A}, \mathcal{A})$ defines the Gerstenhaber bracket operation $[\varphi, \psi]$ on $\text{CC}^*(\mathcal{A}, \mathcal{A})$. More generally in that notation, $\delta = [\mu_{\mathcal{A}}, \cdot]$, and the A_∞ -relations $\mu_{\mathcal{A}} \circ \mu_{\mathcal{A}} = 0$ correspond to $\delta \circ \delta = 0$ via the Jacobi identity. Similarly, $\delta \circ \delta = 0$ in general follows from the bimodule A_∞ -relations.

The Hochschild cohomology is the resulting cohomology:

$$\text{HH}^n(\mathcal{A}, \mathcal{M}) = H^n(\text{CC}^*(\mathcal{A}, \mathcal{M}); \delta).$$

Example. $\text{HH}^n(\mathcal{A}) = H^n(\text{CC}^*(\mathcal{A}, \mathcal{A}); \delta)$ is the Hochschild cohomology of the diagonal.

Example. For cocycles $[\varphi] \in \text{HH}^*(\mathcal{A}, \mathcal{M})$, the 0-part $\varphi^0 : \mathbb{K} \rightarrow \mathcal{M}(X, X)$ has $0 = (\delta\varphi)^0(1) = \mu_{\mathcal{M}}^{0|0}(\varphi^0(1))$. So $\varphi^0(1) \in H^*(\mathcal{M}(X, X); \mu_{\mathcal{M}}^{0|0})$. So for $\mathcal{M} = \mathcal{A}$, $\varphi^0(1) \in H^*(\mathcal{A}(X, X); \mu_{\mathcal{A}}^1)$.

2.14. Cup product on Hochschild cohomology. The cup product on $\text{CC}^*(\mathcal{A}, \mathcal{A})$, denoted $\varphi * \psi$, is defined as follows on $(x_n, \dots, x_1) \in \mathcal{A}(X_n, \dots, X_0)$:

$$\sum (-1)^\# \mu_{\mathcal{A}}^\ddagger(x_n, \dots, x_{R+1}, \varphi^{R-S+1}(x_R, \dots, x_S), x_{S-1}, \dots, x_{r+1}, \psi^{r-s+1}(x_r, \dots, x_s), x_{s-1}, \dots, x_1)$$

where $\sharp = \|\varphi\| \sigma_1^{S-1} + \|\psi\| \sigma_1^{s-1}$ (and of course $\ddagger = n - R + S - r + s$). One can check this descends to a product on the cohomology $\mathrm{HH}^*(\mathcal{A})$.

Remark. *Composition of maps, as in 2.5, would not give a product on $\mathrm{HH}^*(\mathcal{A})$, indeed the differential of a composition is the Gerstenhaber bracket. The bracket terms do not occur in 2.5 because pre-morphisms of bimodules must receive module inputs.*

Example. *The 0-part $(\varphi * \psi)^0(1) \in \mathcal{A}(X, X)$ equals $\mu_{\mathcal{A}}^2(\varphi^0(1), \psi^0(1))$, the usual product.*

We thank Nick Sheridan and the referee for explaining the following analogue of Lemma 2.4.

Lemma 2.5. *If $\varphi \in \mathrm{HH}^*(\mathcal{A})$ satisfies $\varphi^0 = [e_X] \in \mathrm{Hom}(X, X)$ for all objects X , where $[e_X]$ is the cohomological unit, then cup product with φ defines an isomorphism on $\mathrm{HH}^*(\mathcal{A})$.*

Proof. Consider the length filtration on the Hochschild cochain complex

$$F^{\geq p} \mathrm{CC}^* = \text{Hochschild cochains in } \mathrm{CC}^*(\mathcal{A}) \text{ with at least } p \text{ inputs.}$$

The cup product respects the filtration:

$$* : F^{\geq p} \mathrm{CC}^* \otimes F^{\geq q} \mathrm{CC}^* \rightarrow F^{\geq p+q} \mathrm{CC}^*. \quad (2.5)$$

The E_1 -page for the filtration is the Hochschild cohomology of the cohomological category $H^*(\mathcal{A})$, which is just an ordinary linear category (compare Equation (1.13) in Seidel [41]). The length filtration is *exhaustive* (i.e. $\mathrm{CC}^* = \cup F^{\geq p} \mathrm{CC}^*$) and *complete* (i.e. $\mathrm{CC}^* = \varprojlim F^{\geq p} \mathrm{CC}^*$ taken over the reverse inclusions $F^{\geq p} \mathrm{CC}^* \supset F^{\geq p+1} \mathrm{CC}^*$). So the Eilenberg-Moore comparison theorem [49, Theorem 5.5.11] applies to the map of filtered complexes $f = \varphi * \cdot : \mathrm{CC}^*(\mathcal{A}) \rightarrow \mathrm{CC}^*(\mathcal{A})$ (the map respects the filtration due to (2.5) taking $p = 0$ and fixing the first input $\varphi \in F^{\geq 0} \mathrm{CC}^*$). Notice that on the E_1 page the map f induces the map $f^1 : H^*(\mathcal{A}) \rightarrow H^*(\mathcal{A})$ given by the identity, so it is an isomorphism. By the comparison theorem, the map f on cohomology, namely $\varphi * \cdot : \mathrm{HH}^*(\mathcal{A}) \rightarrow \mathrm{HH}^*(\mathcal{A})$, must also be an isomorphism (this version of the comparison theorem does not require the spectral sequence to converge). \square

2.15. The unit in Hochschild cohomology. The analogue of the unit map in 2.6 would not work for the cup product. If the A_∞ -category \mathcal{A} has *strict* units (see 2.10), then define

$$1_{\mathrm{CC}^*}^0 : \mathbb{K} \rightarrow \mathcal{A}(X_0, X_0), \quad 1_{\mathrm{CC}^*}^0(1) = 1_{X_0}, \quad 1_{\mathrm{CC}^*}^r = 0 \text{ for } r > 0.$$

Strict unitality then implies that 1_{CC^*} is a unit for $(\mathrm{HH}^*(\mathcal{A}), *)$.

2.16. Hochschild cohomology as endomorphisms of the diagonal bimodule. There is a chain map from the Hochschild complex to pre-morphisms

$$\mu_{\mathcal{M} \circ \text{left}} : \mathrm{CC}^*(\mathcal{A}, \mathcal{M}) \rightarrow \mathrm{hom}(\mathcal{A}, \mathcal{M}), \quad \varphi \mapsto \mu_{\mathcal{M} \circ \text{left}} \varphi.$$

Explicitly, $\mu_{\mathcal{M} \circ \text{left}} \varphi(x_r, \dots, x_1, \underline{m}, y_1, \dots, y_s)$ equals the following:

$$\sum (-1)^{\|\varphi\| \star} \mu_{\mathcal{M}}^{r-R+S|s}(x_r, \dots, x_{R-1}, \varphi^{R-S+1}(x_R, \dots, x_S), x_{S-1}, \dots, x_1, \underline{m}, y_1, \dots, y_s).$$

where $\star = \sigma(y)_1^s + \deg(\underline{m}) + \sigma(x)_1^{S-1}$.

Lemma 2.6. *The map $\mu_{\mathcal{M} \circ \text{left}}$ is a chain map, inducing*

$$\mathrm{HH}^*(\mathcal{A}, \mathcal{M}) \rightarrow \mathrm{Hom}(\mathcal{A}, \mathcal{M}).$$

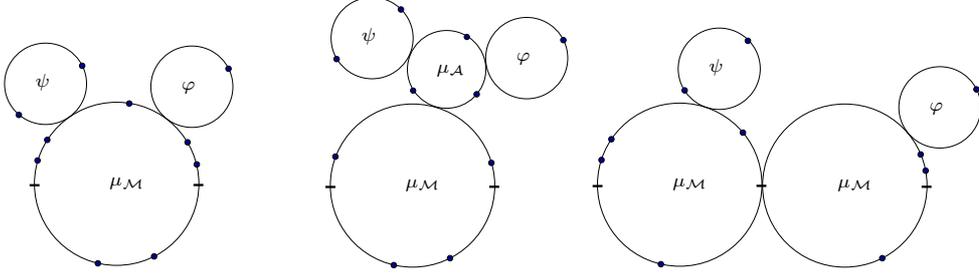
For the diagonal bimodule $\mathcal{M} = \mathcal{A}$, the map is the first term of an A_∞ -homomorphism,

$$\mu_{\mathcal{M} \circ \text{left}} : \mathrm{CC}^*(\mathcal{A}) \rightarrow \mathrm{end}(\mathcal{A}), \quad \varphi \mapsto \mu_{\mathcal{M} \circ \text{left}} \varphi,$$

(where the A_∞ -structure on $\mathrm{CC}^(\mathcal{A})$ was constructed by Getzler [26], and using the dg-algebra structure on $\mathrm{end}(\mathcal{A})$). This in turn induces an algebra homomorphism on cohomology*

$$\mu_{\mathcal{M} \circ \text{left}} : \mathrm{HH}^*(\mathcal{A}) \rightarrow \mathrm{End}(\mathcal{A}).$$

Sketch Proof. This is a purely algebraic result. We will only summarize part of the argument, by showing the bubblings that arise in the context of Fukaya categories:



Consider a 1-family of Floer solutions defined on the domain in the first picture. Such a family can break as in the last two pictures (the other breakings involve terms coming from $\delta\varphi$, $\delta\psi$, $\mu_{\mathcal{A}}$, $\mu_{\mathcal{M}}$, but these can be ignored at the cohomology level). The three small bubbles in the middle picture are the cup product $\psi * \varphi$, so the middle picture is $\mu_{\mathcal{M}} \circ_{\text{left}} (\psi * \varphi)$, the picture on the right is the Yoneda composition $(\mu_{\mathcal{M}} \circ_{\text{left}} \psi) \circ (\mu_{\mathcal{M}} \circ_{\text{left}} \varphi)$.

For the second claim, we remark that the construction of the A_{∞} -functor generalizes the formula for $F^1(\varphi) = \mu_{\mathcal{M}} \circ_{\text{left}} \varphi$ as follows: $F^n(\varphi_n, \dots, \varphi_1)(x_r, \dots, x_1, \underline{m}, y_1, \dots, y_s)$ equals

$$\sum (-1)^{\star\star} \mu_{\mathcal{M}}^{\dagger|s}(x_r, \dots, x_{k_n+1}, \varphi_n^{k_n-\ell_n+1}(x_{k_n}, \dots, x_{\ell_n}), x_{\ell_n-1}, \dots, x_{k_2+1}, \varphi_2^{k_2-\ell_2+1}(x_{k_2}, \dots, x_{\ell_2}), x_{\ell_2-1}, \dots, x_{k_1+1}, \varphi_1^{k_1-\ell_1+1}(x_{k_1}, \dots, x_{\ell_1}), x_{\ell_1-1}, \dots, x_1, \underline{m}, y_1, \dots, y_s).$$

where $\dagger = r - \sum k_j + \sum \ell_j$, and $\star\star = \sum \|\varphi_j\|(\sigma(y)_1^s + \deg(\underline{m}) + \sigma(x)_1^{\ell_j-1})$. \square

Lemma 2.7. *If \mathcal{A} is cohomologically unital (see 2.10), then $\mu_{\mathcal{M}} \circ_{\text{left}}$ is a quasi-isomorphism.*

The result $\text{HH}^*(\mathcal{A}) \cong \text{End}(\mathcal{A})$ is often stated in the literature; the only detailed proof in the literature that we are aware of is due to Ganatra [25, (2.201)] (Ganatra considers algebras, but the proof extends immediately to categories, and note that Ganatra places φ on the right of the module input rather than the left).

Corollary 2.8. *If \mathcal{A} is c -unital, $\text{HH}^*(\mathcal{A}) \cong \text{End}(\mathcal{A})$ is a unital algebra isomorphism.*

Proof. This follows from the previous two Lemmas. In particular, a posteriori $(\text{HH}^*(\mathcal{A}), *)$ has a unit: the preimage of $1_{\mathcal{M}} \in \text{End}(\mathcal{A})$. If \mathcal{A} is strictly unital, one can see directly that $1_{\text{CC}^*} \mapsto 1_{\mathcal{A}}$ at the chain level (using the units from 2.6 and 2.15). \square

Remark. *There is an analogous isomorphism $\text{HH}_*(\mathcal{A}, \mathcal{M}) \cong H^*(\mathcal{M} \otimes_{\mathcal{A}^{\text{opp}} \otimes \mathcal{A}} \mathcal{A})$ of vector spaces (viewing \mathcal{M}, \mathcal{A} respectively as right, left modules over $\mathcal{A}^{\text{opp}} \otimes \mathcal{A}$, see 2.19-2.20), in particular $\text{HH}_*(\mathcal{A}) \cong H^*(\mathcal{A} \otimes_{\mathcal{A}^{\text{opp}} \otimes \mathcal{A}} \mathcal{A})$.*

2.17. Cap product action of HH^* on HH_* .

Corollary 2.9. *The map $\mu_{\mathcal{M}} \circ_{\text{left}}$ induces an algebra homomorphism*

$$\text{HH}_*(\mu_{\mathcal{M}} \circ_{\text{left}}) : \text{HH}^*(\mathcal{A}) \rightarrow \text{End}(\mathcal{A}) \rightarrow \text{End}(\text{HH}_*(\mathcal{A})),$$

therefore $\text{HH}_(\mathcal{A})$ is an $\text{HH}^*(\mathcal{A})$ -module.*

Proof. Combine Lemma 2.6 (for the diagonal bimodule $\mathcal{M} = \mathcal{A}$) and Lemma 2.1. \square

More generally for any \mathcal{A} -bimodule \mathcal{M} , that composite

$$\text{CC}_*(\mu_{\mathcal{M}} \circ_{\text{left}}) : \text{CC}^*(\mathcal{A}, \mathcal{A}) \rightarrow \text{hom}(\mathcal{A}, \mathcal{M}) \rightarrow \text{End}(\text{CC}_*(\mathcal{A}, \mathcal{M})),$$

is the *cap product*. Explicitly $\varphi \in \text{CC}^*(\mathcal{A}, \mathcal{A})$ acts on $c = \underline{m} \otimes x_n \otimes \cdots \otimes x_1 \in \text{CC}_*(\mathcal{A}, \mathcal{M})$ by:

$$\sum (-1)^{\square} \mu_{\mathcal{M}}^{r-R+S|n-s+1}(x_r, \dots, x_{R+1}, \varphi^{R-S+1}(x_R, \dots, x_S), x_{S-1}, \dots, x_1, \underline{m}, x_n, \dots, x_s) \otimes x_{s-1} \otimes \cdots \otimes x_{r+1}$$

where $\square = \sigma_1^r(\deg(\underline{m}) + \sigma_{r+1}^n) + \|\varphi\|(\sigma_{r+1}^n + \deg(\underline{m}) + \sigma_1^{S-1}) + \sigma_{r+1}^{s-1}$. One can check that cocycles φ act as chain maps on $\text{CC}_*(\mathcal{A}, \mathcal{M})$.

2.18. Functorial properties. The Hochschild homology $\text{HH}_*(\mathcal{A}, \mathcal{M})$ is covariantly functorial in \mathcal{A} and \mathcal{M} . Indeed, let $f : \mathcal{A} \rightarrow \mathcal{B}$ be an A_∞ -functor. Recall this means a map on objects $f : \text{Ob}(\mathcal{A}) \rightarrow \text{Ob}(\mathcal{B})$ together with a grading-preserving multi-linear map

$$f^n : \mathcal{A}(X_n, \dots, X_0) \rightarrow \mathcal{B}(fX_n, fX_0) = \text{hom}_{\mathcal{B}}(fX_0, fX_n)$$

for each $n \geq 1$ (using reduced gradings), satisfying the A_∞ -relations:

$$0 = \sum \mu_{\mathcal{B}}^m(f^{k_m}(x_n, \dots, x_{n-k_m+1}), \dots, f^{k_1}(x_{k_1}, \dots, x_1)) - \sum (-1)^{\sigma_1^{s-1}} f^{n-r+s}(x_n, \dots, x_{r+1}, \mu_{\mathcal{A}}^{r-s+1}(x_r, \dots, x_s), x_{s-1}, \dots, x_1)$$

summing over the obvious r, s , and over all partitions $k_1 + \cdots + k_m = n$ and all m .

Lemma 2.10. *An A_∞ -functor $f : \mathcal{A} \rightarrow \mathcal{B}$ induces a change of rings for bimodules:*

$$\mathcal{B}\text{-Bimod} \rightarrow \mathcal{A}\text{-Bimod}, \quad M \mapsto f^*M,$$

which turns a \mathcal{B} -bimodule \mathcal{M} into an \mathcal{A} -bimodule $f^*\mathcal{M}$. There is a tautological chain map

$$\text{taut}_f : \text{CC}_*(\mathcal{A}, f^*\mathcal{M}) \rightarrow \text{CC}_*(\mathcal{B}, \mathcal{M}).$$

Moreover, f induces a canonical morphism $\text{diag}_f : \mathcal{A} \rightarrow f^*\mathcal{B}$ of \mathcal{A} -bimodules (where \mathcal{A}, \mathcal{B} are the diagonal bimodules), yielding a chain map

$$\text{taut}_f \circ \text{CC}_*(f_{\text{diag}}) : \text{CC}_*(\mathcal{A}, \mathcal{A}) \rightarrow \text{CC}_*(\mathcal{A}, f^*\mathcal{B}) \rightarrow \text{CC}_*(\mathcal{B}, \mathcal{B}).$$

Proof. For any $X, Y \in \text{Ob}(\mathcal{A})$, define $f^*\mathcal{M}(X, Y) = \mathcal{M}(fX, fY)$ and let

$$\mu_{f^*\mathcal{M}}^{r|s} = \sum \mu_{\mathcal{M}}^{\tilde{r}|\tilde{s}}(f^{k_{\tilde{r}}}(x_r, \dots, x_{r-k_{\tilde{r}}+1}), \dots, f^{k_{\tilde{s}}}(x_{\tilde{s}}, \dots, x_{\tilde{s}-k_{\tilde{s}}+1}), \underline{\text{Id}} \otimes f^{\ell_1}(x_{\tilde{s}+1}, \dots, x_{\tilde{s}+\ell_1}), \dots, f^{\ell_{\tilde{n}}}(x_{\tilde{s}+\ell_{\tilde{n}}}, \dots, x_{\tilde{s}+\ell_{\tilde{n}}+1}))$$

summing over all partitions $k_1 + \cdots + k_{\tilde{r}} = r$ and $\ell_1 + \cdots + \ell_{\tilde{s}} = s$, and all \tilde{r}, \tilde{s} , where $\underline{\text{Id}}$ is the identity map on the module input. The fact that $f^*\mathcal{M}$ is an A_∞ -bimodule is a consequence of the A_∞ -relations satisfied by f . Define taut_f on $\text{CC}_n(\mathcal{A}, f^*\mathcal{M})$ by

$$\text{taut}_f = \sum \underline{\text{Id}} \otimes f^{\ell_1}(x_{\tilde{s}+1}, \dots, x_{\tilde{s}+\ell_1}) \otimes \cdots \otimes f^{\ell_{\tilde{n}}}(x_{\tilde{s}+\ell_{\tilde{n}}}, \dots, x_{\tilde{s}+\ell_{\tilde{n}}+1})$$

summing over all partitions $\ell_1 + \cdots + \ell_{\tilde{n}} = n$ and all \tilde{n} . The A_∞ -relations again imply that taut_f determines a chain map. The canonical bimodule map in the claim is:

$$\begin{aligned} \text{diag}_f : \mathcal{A}(X_r, \dots, X_0) \otimes \mathcal{A}(X_0, Y_0) \otimes \mathcal{A}(Y_0, \dots, Y_s) &\rightarrow \underline{f^*\mathcal{B}}(X_r, Y_s) = \underline{\mathcal{B}}(fX_r, fY_s) \\ \text{diag}_f(x_r, \dots, x_1, \underline{m}, y_1, \dots, y_s) &= f^{r+s+1}(x_r, \dots, x_1, m, y_1, \dots, y_s). \end{aligned}$$

So explicitly, the composite map $\text{CC}_*(\mathcal{A}, \mathcal{A}) \rightarrow \text{CC}_*(\mathcal{B}, \mathcal{B})$ in the claim is:

$$\sum (-1)^{\diamond} \underline{f^{r+n-s+1}}(x_r, \dots, x_1, \underline{m}, x_n, \dots, x_s) \otimes f^{\ell_1}(x_{s-1}, \dots, x_{s-\ell_1}) \otimes \cdots \otimes f^{\ell_{\tilde{s}}}(x_{r+\ell_{\tilde{s}}}, \dots, x_{r+1})$$

summing over r, s and all partitions $\ell_1 + \cdots + \ell_{\tilde{s}} = s$, where $\diamond = \sigma_1^r(\deg(\underline{m}) + \sigma_{r+1}^n)$. \square

The Hochschild cohomology $\text{HH}^*(\mathcal{A}, \mathcal{M})$ is covariantly functorial in \mathcal{A} and contravariantly functorial in \mathcal{M} . However, unlike $\text{HH}_*(\mathcal{A}, \mathcal{A})$ above, the Hochschild cohomology $\text{HH}^*(\mathcal{A}, \mathcal{A})$ of the diagonal bimodule is not in general functorial in \mathcal{A} with respect to A_∞ -functors.

Lemma 2.11. *The inclusion of a full subcategory $\mathcal{A} \subset \mathcal{B}$, induces a chain map*

$$\text{restriction} : \text{CC}^*(\mathcal{B}, \mathcal{B}) \rightarrow \text{CC}^*(\mathcal{A}, \mathcal{A}).$$

Proof. For a cochain $\varphi \in \text{CC}^*(\mathcal{B}, \mathcal{B})$, we take the composite

$$\mathcal{A}(X_n, \dots, X_0) = \mathcal{B}(X_n, \dots, X_0) \rightarrow \underline{\mathcal{B}(X_n, X_0)} = \underline{\mathcal{A}(X_n, X_0)}$$

where the middle map is φ^n . \square

Theorem 2.12. *If an A_∞ -functor $f : \mathcal{A} \rightarrow \mathcal{B}$ is a quasi-equivalence² then $\text{HH}_*(\mathcal{A}) \cong \text{HH}_*(\mathcal{B})$ and $\text{HH}^*(\mathcal{B}) \cong \text{HH}^*(\mathcal{A})$ for the diagonal bimodules.*

The Theorem is essentially the A_∞ -analogue of Morita invariance (see Seidel [41, Sec.(2d), Remark 2.7]), and we will not prove it. One approach to proving the Theorem, when \mathcal{A} is c-unital, is via Corollary 2.8: $\text{HH}^*(\mathcal{A}) \cong \text{End}(\mathcal{A})$. Indeed, one checks that the map $\mathcal{B}\text{-Bimod} \rightarrow \mathcal{A}\text{-Bimod}$, induced by f via Lemma 2.10, is an equivalence of categories which sends the diagonal bimodule to the diagonal bimodule, and thus $\text{End}(\mathcal{B}) \cong \text{End}(\mathcal{A})$. For HH_* one uses $\text{HH}_*(\mathcal{A}) \cong H^*(\mathcal{A} \otimes_{\mathcal{A}^{\text{opp}} \otimes \mathcal{A}} \mathcal{A})$ (see the Remark after Corollary 2.8).

2.19. Left modules and right modules. We follow Seidel [41, Sec.(1j)] and [40, page 94]. Given two A_∞ -categories $\mathcal{A}_1, \mathcal{A}_2$, there is a generalization of the definition of \mathcal{A} -bimodules \mathcal{M} in 2.2 to $(\mathcal{A}_1, \mathcal{A}_2)$ -bimodules \mathcal{M} . Such a bimodule comprises a collection of graded vector spaces $\mathcal{M}(X, Y)$ for any objects $X \in \mathcal{A}_1, Y \in \mathcal{A}_2$, together with multi-linear maps

$$\mu_{\mathcal{M}}^{r|s} : \mathcal{A}_1(X_r, \dots, X_0) \otimes \mathcal{M}(X_0, Y_0) \otimes \mathcal{A}_2(Y_0, \dots, Y_s) \rightarrow \mathcal{M}(X_r, Y_s)$$

for $r, s \geq 0$, and any objects $X_i \in \mathcal{A}_1, Y_j \in \mathcal{A}_2$ satisfying the A_∞ -relations written in 2.2 except we distinguish between using $\mu_{\mathcal{A}_1}$ (on the x elements) and $\mu_{\mathcal{A}_2}$ (on the y elements).

Now consider the two A_∞ -categories \mathcal{A}, \mathbb{K} . Here the ground field \mathbb{K} is viewed as an A_∞ -category with one object $X_{\mathbb{K}}$, with $\text{hom}_{\mathbb{K}}(X_{\mathbb{K}}, X_{\mathbb{K}}) = \mathbb{K} \cdot \text{id}_{X_{\mathbb{K}}}$ in degree 0, and all $\mu_{\mathbb{K}}^r = 0$.

A *left \mathcal{A} -module* \mathcal{L} is an $(\mathcal{A}, \mathbb{K})$ -bimodule \mathcal{L} with $\mu_{\mathcal{L}}^{r|s} = 0$ for $s \geq 1$. We abbreviate the graded vector spaces by $\mathcal{L}(X) = \mathcal{L}(X, X_{\mathbb{K}})$, and the multi-linear maps by

$$\mu_{\mathcal{L}}^{r|} = \mu_{\mathcal{L}}^{r|0} : \mathcal{A}(X_r, \dots, X_0) \otimes \mathcal{L}(X_0) \rightarrow \mathcal{L}(X_r).$$

A *right \mathcal{A} -module* \mathcal{R} is a $(\mathbb{K}, \mathcal{A})$ -bimodule \mathcal{R} with $\mu_{\mathcal{R}}^{r|s} = 0$ for $r \geq 1$. Let $\mathcal{R}(X) = \mathcal{R}(X_{\mathbb{K}}, X)$,

$$\mu_{\mathcal{R}}^{s|} = \mu_{\mathcal{R}}^{0|s} : \mathcal{R}(Y_0) \otimes \mathcal{A}(Y_0, \dots, Y_s) \rightarrow \mathcal{R}(Y_s).$$

2.20. Tensor products. Let \mathcal{A} be an A_∞ -category, and \mathcal{L}, \mathcal{R} a left and a right \mathcal{A} -module (see 2.19). We will define a chain complex $\mathcal{R} \otimes_{\mathcal{A}} \mathcal{L}$ and an \mathcal{A} -bimodule $\mathcal{L} \otimes \mathcal{R}$ such that

$$\text{HH}_*(\mathcal{A}, \mathcal{L} \otimes \mathcal{R}) \cong H^*(\mathcal{R} \otimes_{\mathcal{A}} \mathcal{L}).$$

The chain complex is given by

$$\mathcal{R} \otimes_{\mathcal{A}} \mathcal{L} = \bigoplus_{X_i \in \text{Ob}(\mathcal{A})} \mathcal{R}(X_n) \otimes \mathcal{A}(X_n, \dots, X_0) \otimes \mathcal{L}(X_0)$$

graded by the sum of the degrees, using reduced degrees for \mathcal{A} , and the differential (of degree +1) on a generator $\mathbf{x} \otimes x_n \otimes \dots \otimes x_1 \otimes \mathbf{l}$ is defined by

$$\begin{aligned} & \sum (-1)^{\text{deg}(\mathbf{l}) + \sigma_1^{R-1}} \frac{\mu_{\mathcal{R}}^{n-R+1}(\mathbf{x}, x_n, \dots, x_R)}{\mu_{\mathcal{R}}^{n-R+1}(\mathbf{x}, x_n, \dots, x_R)} \otimes x_{R-1} \otimes \dots \otimes x_1 \otimes \mathbf{l} + \\ & + \sum (-1)^{\text{deg}(\mathbf{l}) + \sigma_1^{S-1}} \mathbf{x} \otimes x_n \otimes \dots \otimes x_{R+1} \otimes \mu_{\mathcal{A}}^{R-S+1}(x_R, \dots, x_S) \otimes x_{S-1} \otimes \dots \otimes x_1 \otimes \mathbf{l} + \\ & + \sum \mathbf{x} \otimes x_n \otimes \dots \otimes x_{S+1} \otimes \mu_{\mathcal{L}}^{S|}(x_S, \dots, x_1, \mathbf{l}). \end{aligned}$$

summing over the obvious R, S . The \mathcal{A} -bimodule is defined by

$$(\mathcal{L} \otimes \mathcal{R})(X, Y) = \mathcal{L}(X) \otimes \mathcal{R}(Y)$$

²meaning the underlying cohomology level functor is an equivalence. For example, a quasi-isomorphism.

with composition maps

$$\mu_{\mathcal{L} \otimes \mathcal{R}}^{r|s} : \mathcal{A}(X_r, \dots, X_0) \otimes \mathcal{L}(X_0) \otimes \mathcal{R}(Y_0) \otimes \mathcal{A}(Y_0, \dots, Y_s) \rightarrow \mathcal{L}(X_r) \otimes \mathcal{R}(Y_s)$$

defined on generators by

$$\begin{aligned} \mu_{\mathcal{L} \otimes \mathcal{R}}^{0|s}(\mathbb{1}, \mathfrak{r}, y_1, \dots, y_s) &= \mathbb{1} \otimes \underline{\mu}_{\mathcal{R}}^{0|s}(\mathfrak{r}, y_1, \dots, y_s) \\ \mu_{\mathcal{L} \otimes \mathcal{R}}^{r|0}(x_r, \dots, x_1, \mathbb{1}, \mathfrak{r}) &= (-1)^{\deg(\mathfrak{r})} \underline{\mu}_{\mathcal{L}}^{r|0}(x_r, \dots, x_1, \mathbb{1}) \otimes \mathfrak{r} \\ \mu_{\mathcal{L} \otimes \mathcal{R}}^{0|0}(\mathbb{1}, \mathfrak{r}) &= \mathbb{1} \otimes \underline{\mu}_{\mathcal{R}}^{0|0}(\mathfrak{r}) + (-1)^{\deg(\mathfrak{r})} \underline{\mu}_{\mathcal{L}}^{0|0}(\mathbb{1}) \otimes \mathfrak{r} \end{aligned}$$

and $\mu_{\mathcal{L} \otimes \mathcal{R}}^{r|s}$ is zero whenever r, s are both non-zero.

Define $\mathcal{T} : \text{CC}_*(\mathcal{L} \otimes \mathcal{R}) \rightarrow \mathcal{R} \otimes_{\mathcal{A}} \mathcal{L}$ by reordering the positions of \mathfrak{r} and $\mathbb{1}$:

$$\mathcal{T}(\mathbb{1}, \mathfrak{r}, y_1, \dots, y_s) = (-1)^{\deg(\mathbb{1})(\deg(\mathfrak{r}) + \sigma_1^{\mathfrak{r}})} \mathfrak{r} \otimes y_1 \otimes \dots \otimes y_s \otimes \mathbb{1}.$$

After the reordering, the bar differential becomes precisely the differential defined for $\mathcal{R} \otimes_{\mathcal{A}} \mathcal{L}$, so \mathcal{T} is a chain isomorphism inducing the isomorphism $\text{HH}_*(\mathcal{A}, \mathcal{L} \otimes \mathcal{R}) \cong H^*(\mathcal{R} \otimes_{\mathcal{A}} \mathcal{L})$.

2.21. Grading conventions. We summarize here all grading conventions in the paper, and we do not discuss these further (see [41, 2, 8] for a detailed discussion of gradings).

- (0) The grading group in general is $\mathbb{R}/2N\mathbb{Z}$ where N is the minimal Chern number (so $N\mathbb{Z} = c_1(TE)(\pi_2(E))$), as in general it is not possible to assign a grading to a Hamiltonian chord or orbit without additional choices. When working with a single Lagrangian, one can refine the grading group to \mathbb{R} ;
- (1) The Novikov variable t for E has cohomological degree $|t| = 2\lambda_E \in \mathbb{R}$;
- (2) Unreduced gradings for Lagrangian Floer cohomology are as in Seidel [41], although in general the grading group will only be $\mathbb{R}/2\mathbb{Z}$. This improves to an \mathbb{R} -grading in some situations, such as for the endomorphism algebra of a single Lagrangian, assuming monotonicity, and grading the Novikov variable as above;
- (3) For Hamiltonian Floer cohomology, our grading conventions are determined by requiring $QH^*(M) \cong HF^*(H)$ to have degree 0 for closed monotone manifolds;
- (4) and so $c^* : QH^*(E) \rightarrow SH^*(E)$ has degree 0;
- (5) Products preserve unreduced gradings, units lie in unreduced grading 0;
- (6) We never reduce gradings of module inputs/outputs or of $HF^*(H)$, $SH^*(E)$;
- (7) We reduce gradings of morphisms of A_∞ -categories, so all μ^d have degree +1;
- (8) $\mathcal{AF} : \text{HH}_*(\mathcal{F}(E)) \rightarrow \text{HH}_*(\mathcal{W}(E))$ has degree 0;
- (9) $\text{CO} : SH^*(E) \rightarrow HW^*(K, K)$ has degree 0 (unreduced gradings);
- (10) $\text{OC} : \text{HH}_*(\mathcal{W}(E)) \rightarrow SH^*(E)$ has degree $\dim_{\mathbb{C}} E$;
- (11) $\Delta^{r|s} : \mathcal{W}(L_r, \dots, L_0) \otimes \underline{\mathcal{W}(L_0, L'_0)} \otimes \mathcal{W}(L'_0, \dots, L'_s) \rightarrow \underline{\mathcal{W}(L_r, K, L'_s)}$ has degree $\dim_{\mathbb{C}} E$ (the underlined terms are viewed as modules, so their degrees are not reduced).

3. THE WRAPPED FUKAYA CATEGORY

3.1. Symplectic manifolds conical at infinity. (E, ω) will always denote a (non-compact) convex symplectic manifold. Here *convex* means that E is conical at infinity: outside of a compact domain $E^{\text{in}} \subset E$ there is a symplectomorphism

$$\psi : (E \setminus E^{\text{in}}, \omega|_{E \setminus E^{\text{in}}}) \cong (\Sigma \times [1, \infty), d(R\alpha)).$$

where (Σ, α) is a contact manifold, and R is the coordinate on $[1, \infty)$. We call $E \setminus E^{\text{in}}$ the *conical end* of E . The conical condition implies that outside of E^{in} the symplectic form becomes exact: $\omega = d\Theta$ where $\Theta = \psi^*(R\alpha)$. It also implies that the Liouville vector field $Z = \psi^*(R\partial_R)$ (defined by $\omega(Z, \cdot) = \Theta$) will point strictly outwards along ∂E^{in} , and that ψ is induced by the flow of Z for time $\log R$, so we can write $\Sigma = \partial E^{\text{in}}$, $\alpha = \Theta|_{\Sigma}$ (pull-back).

By a *conical structure* J we mean an ω -compatible almost complex structure on E (so $\omega(\cdot, J\cdot)$ is a J -invariant metric) satisfying the *contact type condition* $J^*\Theta = dR$ for large R . On Σ this implies $JZ = Y$ where Y is the Reeb vector field for (Σ, α) defined by $\alpha(Y) = 1$, $d\alpha(Y, \cdot) = 0$. By choosing α or Σ generically, one ensures that the periods of the closed orbits of Y form a countable closed subset of $[0, \infty)$.

3.2. Monotonicity assumptions, Maslov index. We always assume (E, ω) is *monotone*:

$$c_1(TE) = \lambda_E [\omega]$$

for some constant $\lambda_E > 0$ in \mathbb{R} . A submanifold $L \subset (E, \omega)$ is Lagrangian if $\dim_{\mathbb{R}} L = \dim_{\mathbb{C}} E$ and $\omega|_L = 0$. We will always assume that the Lagrangians are *monotone*, meaning:

$$\omega[u] = \lambda_L \mu([u]) \quad \text{for } [u] \in \pi_2(E, L)$$

for some constant $\lambda_L > 0$, where μ is the Maslov index and we abbreviated $\omega[u] = \int_{\mathbb{D}} u^* \omega$. Recall the Maslov index $\mu(u) \in \mathbb{Z}$ is associated to any continuous disc $u : (\mathbb{D}, \partial\mathbb{D}) \rightarrow (E, L)$, and $\mu(u)$ is even for orientable L . The Maslov index only depends on the $\pi_2(E, L)$ class of u . ‘‘Gluing in a sphere’’ $v : S^2 \rightarrow E$ causes $\mu(v\#u) = 2c_1(TE)([v]) + \mu(u)$. This gluing relation forces E to be monotone, and the monotonicity constants are therefore related by $\lambda_L = \frac{1}{2\lambda_E}$.

3.3. The Novikov field. We will always work over the field

$$\Lambda = \left\{ \sum_{i=0}^{\infty} a_i t^{n_i} : a_i \in \mathbb{K}, n_i \in \mathbb{R}, \lim n_i = \infty \right\},$$

called the *Novikov field*, where \mathbb{K} is some chosen ground field, and t is a formal variable. The choice of \mathbb{K} will not matter, except in the toric examples at the end of the paper, where we will specialize to $\mathbb{K} = \mathbb{C}$. The Novikov field is graded by placing t in (real) degree

$$|t| = 2\lambda_E.$$

Floer/holomorphic solutions will be counted with Novikov weights, e.g. for the quantum product on $QH^*(E)$, holomorphic spheres u are counted with weight $\pm t^{\omega[u]}$, lying in degree

$$2\lambda_E \omega[u] = 2c_1(TE)[u].$$

Recall by 2.21, the grading is understood to lie in $\mathbb{R}/2N$, because there is an ambiguity in gradings in Floer theory (arising from attaching spheres which causes gradings to change by $c_1(TE)$ integrated over spheres). For simply-connected E (and as usual L, E monotone), a part of the Floer theory, e.g. $QH^*(E)$, $SH^*(E)$, $HF^*(L, L)$, $HW^*(L, L)$, is in fact \mathbb{R} -graded.

3.4. The choice of Lagrangians in $\text{Ob}(\mathcal{W}(\mathbf{E}))$. The objects of $\mathcal{W}(E)$ are the closed monotone orientable Lagrangian submanifolds of E , together with all non-compact monotone orientable Lagrangian submanifolds $L \subset E$ which are *conical at infinity*, meaning

- (1) L intersects Σ transversally;
- (2) the pull-back $\Theta|_L = 0$ near Σ and on the conical end.

Conditions (1)-(2) ensure that $\ell = L \cap \Sigma$ is a Legendrian submanifold of the contact manifold $(\Sigma, \Theta|_{\Sigma})$, and that L has the form $\ell \times [1 - \epsilon, \infty) \subset \Sigma \times [1 - \epsilon, \infty)$ near the conical end. Conversely, a monotone Lagrangian $L \subset E^{\text{in}}$ with $\Theta|_L = 0$ near Σ can be extended to an L of the above form. In fact, any monotone Lagrangian $L \subset E^{\text{in}}$ which intersects Σ in a Legendrian submanifold, can be Hamiltonianly isotoped in E^{in} relative to ∂E^{in} so that $\Theta|_L$ will vanish near Σ . This can be proved as in [8, Lemma 3.1] (it is a local argument, and the Liouville flow is defined near Σ). Finally, as will be explained in detail in 3.7, one must in fact restrict $\text{Ob}(\mathcal{W}(E))$ to Lagrangians having the same m_0 -value.

Remark 3.1 (Grading, characteristic). For $\mathcal{W}(E)$ to be graded one would require $2c_1(E, L) = 0 \in H^2(E, L; \mathbb{Z})$. For $\mathcal{W}(E)$ to be defined over a field of characteristic $\neq 2$ one would require $w_2(L) = 0 \in H^2(L; \mathbb{Z}/2)$ and then an L should come with a choice of Pin structure. One also needs to make a choice of grading for the Lagrangians. For a discussion of this brane data, we refer to [41, 8]: we will not discuss gradings or orientation signs for sake of brevity.

3.5. Reeb chords and Hamiltonian chords. Fix a Hamiltonian $H : E \rightarrow [0, \infty)$ which on the conical end has the form $H(\sigma, R) = R$ for $(\sigma, R) \in \Sigma \times [1, \infty)$. Later, in 4.2 (see the Convention), we will make a global rescaling of H , but the reader can ignore this for now.

Let X be the Hamiltonian vector field: $\omega(\cdot, X) = dH$. The Reeb vector field on Σ is $X|_{\Sigma}$. Fix a countable collection $L_i \in \text{Ob}(\mathcal{W}(E))$, so they give rise to $\ell_i \subset \Sigma$.

A *Reeb chord of period w* is a trajectory $[0, 1] \rightarrow \Sigma$ of $wX|_{\Sigma}$ with ends on ℓ_i, ℓ_j (for some, possibly non-distinct, indices i, j). It is an *integer Reeb chord* if w is an integer. By rescaling ω (or by choosing Σ generically), there are no integer Reeb chords. An *integer X -chord of weight w* is a flowline $[0, 1] \rightarrow E$ of wX with ends on L_i, L_j for integer values $w \geq 1$. These lie in E^{in} (otherwise, there would be one in the conical end projecting to an integer Reeb chord in Σ). For a generic choice of H (see [8]), integer X -chords are non-degenerate and any point of L_i is not both an endpoint and a starting point of two (possibly non-distinct) integer X -chords (in particular, no critical point of H may lie in $L_i \cap L_j$). So only finitely many integer X -chords have ends on given L_i, L_j , and they lie in E^{in} .

3.6. The Floer differential \mathfrak{d} . Given $L_i, L_j \in \text{Ob}(\mathcal{W}(E))$, let $CF^*(L_i, L_j; wH)$ denote the free Λ -module generated by the (finitely many) integer X -chords of weight w with ends on L_i, L_j . Fix a conical structure J on E (we remark that the contact type condition in 3.1 on the conical end is equivalent to $JZ = X$). The chain differential

$$\mathfrak{d} : CF^*(L_i, L_j; wH) \rightarrow CF^{*+1}(L_i, L_j; wH)$$

is defined as follows on a generator y . Let $u : \mathbb{R} \times [0, 1] \rightarrow E$ be a non-constant isolated (modulo \mathbb{R} -translation) solution of Floer's equation $\partial_s u + J(\partial_t u - wX) = 0$, satisfying: Lagrangian boundary conditions $u(\cdot, 0) \in L_i$ and $u(\cdot, 1) \in L_j$; and asymptotic conditions $u \rightarrow x, y$ as $s \rightarrow -\infty, +\infty$ respectively, where x is an X -chord. Then u contributes

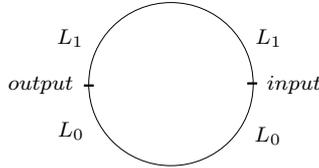
$$\pm t^{E_{\text{top}}(u)} x$$

to $\mathfrak{d}y$, where $E_{\text{top}}(u) \geq 0 \in \mathbb{R}$ is defined by:

$$E_{\text{top}}(u) = \int_S u^* \omega - d(u^*(wH) \wedge dt) = \int_S u^* \omega + wH(x) - wH(y) = \frac{1}{2} \int_S \|du - wX \otimes dt\|^2 ds \wedge dt.$$

We will always omit discussing orientation signs \pm , which is done carefully in [8] by replacing each Λx by an orientation line.

The strip $\mathbb{R} \times [0, 1]$ can also be viewed as the disc with two boundary punctures at ± 1 :



This domain is unstable: there is a residual \mathbb{R} -translation symmetry in $\text{PSL}(2, \mathbb{R})$ fixing ± 1 . Non-constant solutions arise in \mathbb{R} -families, and we quotient moduli spaces of non-constant solutions by this action, and only then count the 0-dimensional part to define \mathfrak{d} .

If $\mathfrak{d}^2 = 0$, then the Floer cohomology of L_i, L_j using wH is defined as

$$\boxed{HF^*(L_i, L_j; wH) = H^*(CF^*(L_i, L_j; wH); \mathfrak{d})}$$

Since E is monotone, $\mathfrak{d}^2 = 0$ precisely when L_i, L_j have the same m_0 -value (see 3.7).

3.7. The role of monotonicity. We keep the discussion quite general by supposing that we have a 1-dimensional family of “Floer solutions” and we study the possible bubblings.

First observe that sphere bubbling is ruled out for dimension reasons, since by monotonicity a non-constant holomorphic sphere will increase the index of the family by at least 2.

In the proof that $\mathfrak{d}^2 = 0$, a non-constant J -holomorphic sphere of Maslov index 2 (Chern number 1) may bubble off from an s -independent strip $u(s, t) = x(t)$, where x is an integer X -chord of weight w . But this bubbling is ruled out for generic H, J , as follows. Recall that for generic J , the J -holomorphic spheres with $c_1 = 1$ come in a family of real dimension $2n - 4$, so they sweep out a real codimension 2 set. In principle, this set could hit our moduli space of strips. But for index reasons, the only bubbling-off allowed is that of a constant strip. Now, the image of a constant strip is a Hamiltonian chord, which is 1-dimensional, so generically this does avoid that moduli space of spheres.

Now consider disc bubbling. By disc bubble, we mean a non-constant holomorphic disc $u : \mathbb{D} \rightarrow E$, bounding a single Lagrangian $u(\partial\mathbb{D}) \subset L$ say, with one boundary marked point $z_0 \in \mathbb{D}$ landing at a prescribed generic point $p_L = u(z_0) \in L$ (so z_0 plays the role of the node when our family breaks). The index of families is additive, so only isolated discs will bubble off in our 1-dimensional family and so, by monotonicity and a dimension count, the disc must have Maslov index $\mu(u) = 2$.

Technical Remark about transversality. *Recall that transversality can be achieved for simple discs (i.e. somewhere injective discs). By Lazzarini [28], for any holomorphic disc $u : (\mathbb{D}, \partial\mathbb{D}) \rightarrow (X, \mathcal{L})$, where X is a smooth almost complex manifold and \mathcal{L} is a totally real submanifold, there is an underlying simple holomorphic disc $v : (\mathbb{D}, \partial\mathbb{D}) \rightarrow (X, \mathcal{L})$ with image contained in $u(\mathbb{D})$, and the relative homology class $[u] \in H_2(X, \mathcal{L})$ is a linear combination $\sum n_i [v_i]$ of such simple discs, where $n_i \in \mathbb{Z}_{>0}$. In our case, $X = E$, $\mathcal{L} = L$ are monotone, and the Maslov index is additive, so Maslov index 2 discs (which is the minimal Maslov index, as L is orientable) must be simple: otherwise u has strictly larger area than v_i and thus strictly larger Maslov index. It follows that transversality can be achieved for Maslov 2 discs.*

Now consider the moduli space $\mathcal{M}_1(\beta)$ of all such Maslov 2 discs in the class $\beta \in \pi_2(E, L)$ but do not impose a condition on the value at the marked point. Let $\text{ev} : \mathcal{M}_1(\beta) \rightarrow L$ be the evaluation at the marked point. Then $\text{ev}_*[\mathcal{M}_1(\beta)]$ determines a locally finite cycle in L (it is a cycle and not just a chain, because $\mathcal{M}_1(\beta)$ cannot have boundary degenerations since, by monotonicity, 2 is the minimal Maslov index of non-constant disc bubbles). For dimension reasons this lf-cycle must be a scalar multiple of $[L]$ (viewed as a locally finite cycle since L may be non-compact - we explain this in 3.8). Then define

$$\mathfrak{m}_0(L) = \sum t^{\omega[\beta]} \text{ev}_*[\mathcal{M}_1(\beta)] = m_0(L) [L] \in C_{\dim(L)}^{\text{lf}}(L; \Lambda), \text{ where } m_0(L) \in \Lambda.$$

When there are no such $\mu = 2$ discs, then $\mathfrak{m}_0(L) = 0$, and if this holds for all objects L then the A_∞ -category is called *unobstructed*, and no disc bubbling discussion is necessary. In the applications we have in mind, we typically expect $\mu = 2$ discs. So we briefly explain why in the monotone case all bubbling contributions will in fact cancel out when we restrict the category to Lagrangians having the same m_0 -value. In other words, monotonicity allows us to pretend that our A_∞ -categories $\mathcal{W}(E, \lambda)$ are all unobstructed.

Firstly, a standard Gromov compactness argument, looking at the possible degenerations of Maslov 2 discs with two boundary punctures, with Lagrangian boundary conditions L_i, L_j between the punctures, shows that:

$$\mathfrak{d} \circ \mathfrak{d}(x) = (m_0(L_i) - m_0(L_j)) x.$$

For this argument, we refer the reader to Oh [30, 31]. Thus, one cannot even define the cohomology $HF^*(L_i, L_j; wH)$ when $m_0(L_i) \neq m_0(L_j)$, let alone the category. If we restrict to only Lagrangians with equal $m_0(L)$ values, then $\mathfrak{d} \circ \mathfrak{d} = 0$. For higher A_∞ -operations, the bubbling of such Maslov 2 discs is ruled out for dimension reasons (the index of the main component would have negative virtual dimension). We summarise this in the following Lemma (we will explain the claim about eigenvalues of c_1 in 9.2).

Lemma 3.2. *$\mathcal{W}(E)$ is a disjoint union of A_∞ -categories $\mathcal{W}_\lambda(E)$ indexed by $\lambda \in \Lambda$, where all $L \in \text{Ob}(\mathcal{W}_\lambda(E))$ satisfy $m_0(L) = \lambda$. Since E is monotone, and assuming $\text{char } \mathbb{K} = 0$, the non-triviality of $\mathcal{W}_\lambda(E)$ implies that λ is an eigenvalue of quantum cup product by c_1 :*

$$c_1(TE) : QH^*(E) \rightarrow QH^*(E).$$

The analogous statement holds for the compact category $\mathcal{F}(E)$ or the Fukaya category $\mathcal{F}(B)$ of a closed (rather than convex) monotone manifold B .

Convention. We always treat $\mathcal{W}(E)$ as a non-curved A_∞ -category which splits into non-interacting pieces, not as a curved category. When we write $\mathcal{W}(E)$ we actually always mean $\mathcal{W}_\lambda(E)$ for some fixed m_0 -value λ . The same comment pertains to $\mathcal{F}(B)$ and $\mathcal{F}(E)$.

3.8. The construction of \mathfrak{m}_0 in the non-compact setting.

Lemma 3.3. *Let E be conical at infinity. Let $L \subset E$ be any Lagrangian submanifold conical at infinity (3.4). Then all pseudo-holomorphic discs $u : (\mathbb{D}, \partial\mathbb{D}) \rightarrow (E, L)$ have $u(\partial\mathbb{D})$ lying either in E^{in} or in a contact hypersurface $\Sigma \times \{R\}$ of the conical end.*

Proof. For general reasons of convexity [8, Sec.7.3], which essentially reduce to Green's formula, any pseudo-holomorphic map $v : S \rightarrow E$ defined on a Riemann surface S (which can have corners) with $v(S) \subset \Sigma \times [R, \infty)$ landing in the conical end of E , with $v(\partial S) \subset L \cup (\Sigma \times \{R\})$, must in fact lie entirely in $\Sigma \times \{R\}$. This maximum principle uses the fact that $\Sigma \times \{R\}$ is a contact hypersurface in E on which the Liouville flow is outward-pointing. The claim follows immediately by taking v to be the restriction of u to $u^{-1}(\Sigma \times [R, \infty))$ (for generic R so that $\Sigma \times \{R\}$ intersects the image of u transversely). \square

By the Lemma, given a point p_L of L , all holomorphic discs bounding L satisfying $p_L \in u(\partial\mathbb{D})$ must lie entirely in a compact region of E determined by p_L . Thus all the constructions relating to the moduli space of holomorphic discs and in particular the construction of $\mathfrak{m}_0(L)$ reduce to the constructions in the closed monotone case, provided we replace cycles/chains by locally finite cycles and locally finite chains. Indeed, for $L \subset E$ a monotone Lagrangian in $\text{Ob}(\mathcal{W}_\lambda(E))$, on each compact subset $C \subset E$ the chain \mathfrak{m}_0 represents (up to the scalar multiple $\lambda = m_0(L)$) the fundamental cycle for the pair $(C \cap L, \partial C \cap L)$. Exhausting E by a sequence of larger and larger compact subsets C proves that $\mathfrak{m}_0(L)$ is, up to the scalar multiple $\lambda = m_0(L)$, an infinite sum of chains which is precisely the locally finite cycle $[L]$.

3.9. The homotopy \mathfrak{K} . Denote \mathfrak{K} the Floer continuation map

$$\boxed{\mathfrak{K} : CF^*(L_i, L_j; wH) \rightarrow CF^*(L_i, L_j; (w+1)H)}$$

determined by a monotone homotopy from $(w+1)H$ to wH . This involves counting isolated solutions similar to those in the definition of \mathfrak{d} , except now J, H depend on the domain coordinate s and there is no \mathbb{R} -reparametrization. The choice of homotopy J_s, H_s interpolating the given data on the ends is chosen once and for all. Section 3.11 describes the auxiliary data carefully (for \mathfrak{K} one uses: $d = 1$, $F = \{1\}$, $w_0 = w+1$, $w_1 = w$, and notice that the \mathbb{R} -translational freedom of the domain is fixed by imposing that the preferred point lies at 0).

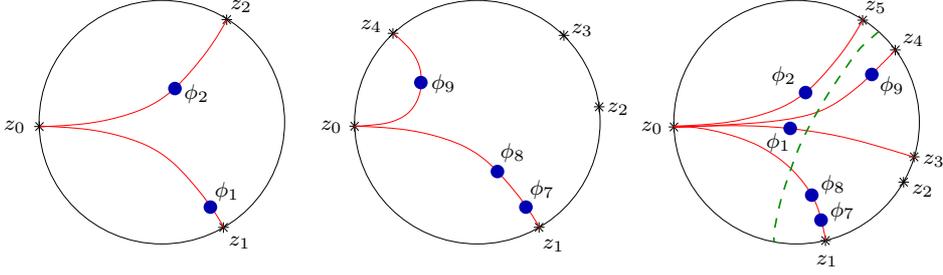


FIGURE 1. Left: $S = \text{pair-of-pants surface}$, $d = 2$, $F = \{1, 2\}$, $p_1 = 1$, $p_2 = 2$. Middle: $d = 4$, $F = \{7, 8, 9\}$, $p_7 = p_8 = 1$, $p_9 = 4$. Right: the dotted line is a cut, which defines a possible breaking.

The above maps \mathfrak{K} are the chain maps which at the level of cohomology are used to define the *cohomology direct limit* as $w \rightarrow \infty$:

$$HW^*(L_i, L_j) = \varinjlim HF^*(L_i, L_j; wH)$$

3.10. **Telescope construction of the wrapped complex.** Following [8, Sec.3.7, Sec.5.1],

Definition 3.4. *The wrapped complex is*

$$CW^*(L_i, L_j) = \bigoplus_{w=1}^{\infty} CF^*(L_i, L_j; wH)[\mathbf{q}]$$

where \mathbf{q} is a formal variable of degree -1 satisfying $\mathbf{q}^2 = 0$. This is a module over the exterior algebra $\mathbb{C}[\partial_{\mathbf{q}}]$, where the differentiation operator $\partial_{\mathbf{q}}$ has degree $+1$. The chain differential is chosen to respect this module structure:

$$\mu^1(x + \mathbf{q}y) = (-1)^{|x|}\partial x + (-1)^{|y|}(\mathbf{q}\partial y + \mathfrak{K}y - y) \quad (\text{using unreduced degrees})$$

Observe that for a \mathfrak{d} -cocycle y , the role of $\mathbf{q}y$ is to identify y and $\mathfrak{K}y$ at the cohomology level, as expected in the cohomology direct limit; whereas the role of the subcomplex ($\partial_{\mathbf{q}} = 0$) is to yield representatives of $\bigoplus HF^*(L_i, L_j; wH)$. We refer to [8, Sec.3.7] for the proof that $H^*(CW^*(L_i, L_j); \mu^1) \cong HW^*(L_i, L_j)$ canonically.

3.11. **The auxiliary data.** By [8], there is a universal family $\mathcal{R}^{d+1, \mathbf{p}} \rightarrow \mathcal{R}^{d+1}$ over the moduli space of stable discs, with elements (S, F, \mathbf{p}, ϕ) where (Figure 1):

- (1) F is a finite set of indices f (with $d \geq 1$, $d + |F| \geq 2$).
Remark. Later, when we define μ^d in terms of a sum of contributions $\mu^{d, \mathbf{p}, \mathbf{w}}$, we will require F to be a subset of $\{1, \dots, d\}$;
- (2) A map $\mathbf{p} : F \rightarrow \{1, \dots, d\}$ defining labels p_f (possibly not distinct).
Remark. When we define μ^d , we will require \mathbf{p} to be an inclusion $F \subset \{1, \dots, d\}$;
- (3) **Disc with $d + 1$ ends.** S is a Riemann surface isomorphic to \mathbb{D} with $d + 1$ distinct boundary punctures z_0, z_1, \dots, z_d (of which z_0 is distinguished, and the others are ordered according to the orientation of ∂S);
- (4) **Strip-like parametrizations.** $\phi = (\phi_f)_{f \in F}$ are holomorphic $\phi_f : S \rightarrow Z$ extending to smooth isomorphisms on the compactifications $\overline{S} \rightarrow \overline{Z}$, with $\phi_f(z_0) = -\infty$, $\phi_f(z_{p_f}) = +\infty$ (where Z is the strip $\mathbb{R} \times [0, 1]$ viewed as a disc with 2 ends). This is the same as a choice of a **preferred point** on the hyperbolic geodesic joining z_0, z_{p_f} (corresponding via ϕ_f to the point $(0, \frac{1}{2})$ and to the geodesic $(t = \frac{1}{2}) \subset Z$);

We will omit the discussion of finite symmetries [8, Sec.2.3, 2.6, Lemma 3.7], but we mention that all auxiliary data constructions are done so as to preserve the symmetries which exchange preferred points lying on the same geodesic.

The above family is compactified by boundary strata described by the natural bubbling configurations one expects when a subset of the preferred points of (4) converges towards a boundary puncture [8, Fig.2, Fig.3]. We omit discussing gluing [8, Sec.2.4], but we mention one important technical aspect:

Remark 1. *One actually allows the geodesics in (4) to be deformed (except near the ends). We explain why in an example. Suppose we want to glue the first two discs in Figure 1 at one puncture. Call the discs S^-, S^+ and we want to glue $z_1 \in \overline{S^-}$ with $z_0 \in \overline{S^+}$. Gluing gives a 6-punctured disc with a geodesic and 2 curves which are not geodesics (they overlap along a common curve carrying the preferred point ϕ_1) but they can be deformed into two geodesics. The “glued” F is $F = F^- \sqcup F^+ = \{1, 2, 7, 8, 9\}$. There are four choices for the “glued” \mathbf{p} : z_{p_1} can be any of the z_1, z_2, z_3, z_4 coming from S^+ , and in the cases z_2 and z_3 we would need to draw a new geodesic before gluing, so that we obtain glued data as in (4) (up to deforming). Now observe the third disc in Figure 1. The dotted line is a “cut”: we consider a 1-parameter family of such discs in which the auxiliary data contained on one side of the cut converges towards a common boundary puncture. In the limit, the disc breaks into two, namely S^- glued with S^+ with “glued” \mathbf{p} determined by $\mathbf{p}(1) = 3$, and the new geodesic we draw on S^+ is remembering that ϕ_1 came from the geodesic joining z_0 to z_3 before breaking.*

Now assign further auxiliary data:

- (1) **Weights.** $\mathbf{w} = (w_0, \dots, w_d)$ are positive integers with $w_0 = \sum_{k \geq 1} w_k + |F|$;
- (2) **Strip-like ends.** Holomorphic embeddings $\epsilon_k : Z_+ \rightarrow S$, for $k = 1, \dots, d$, of $Z_+ = [0, \infty) \times [0, 1] \subset Z$ mapping the boundaries $t = 0, 1$ to ∂S and converging to z_k as $s \rightarrow +\infty$. Similarly $\epsilon_0 : Z_- \rightarrow S$ for the negative strip $Z_- = (-\infty, 0] \times [0, 1]$;
- (3) **Closed 1-forms.** Closed 1-forms α_k on S for $k = 1, \dots, d$ pulling back to: 0 on ∂S ; dt via ϵ_k, ϵ_0 on $(|s| \gg 0) \subset Z_{\pm}$; 0 via the other ϵ_j on $(s \gg 0) \subset Z_+$.
- (4) **Subclosed 1-forms.** 1-forms β_f on S with $d\beta_f \leq 0$ (for the complex orientation of S), $d\beta_f = 0$ near ∂S , and β_f pulls back to: 0 on ∂S ; dt via ϵ_0 on $s \ll 0$; 0 via the other ϵ_k on $s \gg 0$. One can choose β_f so that $d\beta_f = 0$ except on a small neighbourhood of the preferred point ϕ_f , however this choice is not necessary (see 3.13).

The above data determines $\gamma = \sum_{k \geq 1} w_k \alpha_k + \sum_{f \in F} \beta_f$, a 1-form on S with $d\gamma \leq 0$, $d\gamma = 0$ near ∂S . Also γ pulls-back to: 0 on ∂S ; $w_j dt$ via ϵ_j on $(|s| \gg 0) \subset Z_{\pm}$. Observe that this is consistent with Stokes’s theorem, $0 \leq -\int d\gamma = w_0 - \sum_{k \geq 1} w_k = |F|$.

The choices of α, β can be made to vary smoothly over the moduli space [8, Sec.2.5, 2.6] consistently with the boundary strata, by an inductive argument (we discuss this further in 3.13). We make one important technical remark about gluing:

Remark 2. *When gluing a disc S^- at z_k^- with a disc S^+ at z_0^+ , we require the weights to agree: $w_k^- = w_0^+$. The α forms glue naturally. For the β forms: let $\beta_{f^+} = \alpha_k^-$ on S^- , $\beta_{f^-} = 0$ on S^+ , these glue naturally to yield β_f ’s on the gluing indexed by $f \in F = F^- \sqcup F^+$.*

The last piece of auxiliary data is the almost complex structure:

- (1) **Almost complex structures.** A family $(I_z)_{z \in S}$ of conical structures which pull-back via ϵ_j on $|s| \gg 0$ to almost complex structures I_t^w depending only on t, w (which are chosen before-hand for each integer w). This choice can be made smoothly over the moduli space, and compatibly with gluing;
- (2) **Infinitesimal deformations.** Certain deformation data K for I [8, Sec.3.2], namely: K is an endomorphism of TE which anti-commutes with I , such that $\omega(\cdot, K\cdot)$ is

symmetric (together with a condition at infinity which ensures that $J = I \exp(-IK)$ is of contact type at infinity);

- (3) **Regular almost complex structures.** The I, K give rise to a family of perturbed almost complex structures $J = I \exp(-IK)$, agreeing with I asymptotically, and which vary smoothly over the moduli space.

3.12. Domain reparametrizations and moduli. There is a (real 3-dimensional) $PSL(2, \mathbb{R})$ reparametrization action on discs by biholomorphisms which moves the boundary punctures z_0, z_1, \dots, z_d . So we may for example fix the choice of parametrization by choosing the location of three of the boundary marked points, say $z_0 = -1, z_1 = e^{-\pi i/3}, z_2 = e^{+\pi i/3}$. Indeed, the position of the remaining boundary points z_3, \dots, z_d on the anti-clockwise boundary arc from z_2 to z_0 determine a choice of $d-3$ real local coordinates for \mathcal{R}^{d+1} . All the auxiliary data above is chosen for a given moduli element $[S] \in \mathcal{R}^{d+1}$, rather than a parametrized choice of disc S , meaning: the data is constructed consistently with reparametrizations of the underlying domain S (the use of strip-like parametrizations ϕ_f ensure that the data is constructed independently of the choice of parametrization of the domain). Given a parametrized disc S , there are new parameters in the auxiliary data involved in the choice of the location of the preferred points, but we emphasize that a biholomorphism of the parametrized domain with data (so pulling back all auxiliary data) gives rise to the same moduli element. So later, when we count Floer solutions for a given parametrized domain with data, the two solutions which are related by a biholomorphism of the domain with data are considered to be the same solution. In particular, for \mathfrak{d} and \mathfrak{R} we emphasized that there is only one moduli element, so we fixed the domain with data once and for all (input at $+1$, output at -1 , and for \mathfrak{R} we put the preferred point at 0), and we counted solutions (for \mathfrak{d} we quotiented by the residual \mathbb{R} -translational freedom on non-constant solutions and we ignored constant solutions).

3.13. Existence of universal choices of auxiliary data. We mentioned that the auxiliary data from 3.11 needs to be constructed consistently over the universal family $\mathcal{R}^{d+1, \mathbf{p}}$. If one were to attempt this construction explicitly, it would mean: given a Riemann surface S with strip-like ends and some additional data ϕ_f (the preferred points), one needs to stipulate explicit rules for constructing the one-forms α_k and β_f consistently as we vary S and ϕ_f . For example, one typically picks the support of $d\beta_f$ to be located near the preferred point ϕ_f , although this is not necessary in general. This construction needs to be done compatibly with gluing/breaking, i.e. over the compactification $\overline{\mathcal{R}}^{d+1, \mathbf{p}}$, so that when disc bubbling occurs the auxiliary data converges to the correct auxiliary data on the various components. This amounts to constructing a section of the bundle

$$\mathcal{A}ux \rightarrow \overline{\mathcal{R}}^{d+1, \mathbf{p}} \tag{3.1}$$

described below. It is sufficient to prove the existence of the section, without giving an explicit construction. We keep the discussion quite general, since the same argument will apply to the construction of consistent auxiliary data when we define operations, etc. We want the fibre bundle (3.1) to satisfy the following properties:

- (1) the fibre over (S, ϕ_f) is the space of all choices of auxiliary data (so α_k, β_f , etc.),
- (2) the bundle is locally trivial,
- (3) the fibre is non-empty,
- (4) the fibre is contractible.

If this holds, then it follows that there exists a section $\overline{\mathcal{R}}^{d+1, \mathbf{p}} \rightarrow \mathcal{A}ux$, using a standard local-to-global argument [47, Sec.12.2] (the obstruction to extending a section of a fibration lives in the relative cohomology of the base but with coefficients in the homotopy groups of

the fibre [27, Sec.6.4], so we have no obstructions by (4)). The section is precisely a choice of universally consistent auxiliary data. The local-to-global argument gets applied inductively over the strata of $\overline{\mathcal{R}}^{d+1, \mathbf{P}}$, so that the section constructed on the lower strata gets extended when one defines the section on the higher strata, i.e. compatibly with gluing/breaking.

The problem of finding universal data therefore reduces in practice to checking conditions (3) and (4). For the data α_k, β_f : (3) follows easily from de Rham theory (e.g. compare [34, Appendix A]), and (4) follows because the space of choices of α_k, β_f is convex: linear interpolations of two sets of choices will continue to satisfy all the requirements. Two such choices of universal data can be linearly interpolated, so a standard homotopy argument implies that the constructions we make in Floer theory for those two choices are chain homotopic (and thus homological constructions will be independent of the choice of auxiliary data).

The local-to-global argument [47, Sec.12.2] is very general, and includes the above infinite dimensional fibre bundle. If one wanted to be more explicit, one can appeal to explicit coordinates. The space $\overline{\mathcal{R}}^{d+1, \mathbf{P}}$ carries a canonical structure of a smooth manifold with corners [8, Cor 6.5]. In particular the marked points and strip-like ends can be trivialized, so one can locally identify the fibres of the fibre bundle near the boundary strata so that it is easy to extend smooth sections from a boundary stratum to a higher stratum by ‘‘averaging’’.

3.14. Moduli space of pseudo-holomorphic maps. Let x_k be X -chords of weight w_k with ends on $(L_0, L_d), (L_0, L_1), (L_1, L_2), \dots, (L_{d-1}, L_d)$, respectively for $k = 0, \dots, d$, and where $L_i \in \text{Ob}(\mathcal{W}(E))$. Denote by $\mathcal{R}^{d+1, \mathbf{P}}(\mathbf{x})$ the moduli space of solutions $u : S \rightarrow E$ of

$$(du - X \otimes \gamma)^{0,1} \equiv \frac{1}{2}[(du - X \otimes \gamma) + J \circ (du - X \otimes \gamma) \circ j] = 0$$

with Lagrangian boundary conditions $u(\partial_i S) \subset L_i$ where $\partial_i S$ is the component of ∂S between z_i and z_{i+1} ; and asymptotic conditions $u \circ \epsilon_k \rightarrow x_k$ as $|s| \rightarrow \infty$. For a generic choice of K the $\mathcal{R}^{d+1, \mathbf{P}}(\mathbf{x})$ are smooth manifolds of the expected dimension [8, Thm 3.5].

In the exact setup of [8], there is a natural compactification of $\mathcal{R}^{d+1, \mathbf{P}}(\mathbf{x})$ modeled on the compactification of $\mathcal{R}^{d+1, \mathbf{P}}$ [8, Sec.3.5]. The compactness relied on two ingredients: (1) a maximum principle forcing solutions to lie in E^{in} [8, Sec.7.4]; (2) an a priori energy estimate:

$$E(u) = \int_S \frac{1}{2} \|du - X \otimes \gamma\|^2 \equiv \int_S u^* \omega - u^* dH \wedge \gamma \leq \int_S u^* \omega - d(u^* H \wedge \gamma) \equiv E_{\text{top}}(u)$$

using $H \geq 0, d\gamma \leq 0$, so by Stokes’s theorem the topological energy E_{top} is bounded (in the exact setup $\omega = d\Theta$) in terms of the asymptotics \mathbf{x} using action functionals [8, Sec.7.2].

In our non-exact setup, (1) still holds since the argument only relied on exactness of ω on the conical end. We no longer have an a priori estimate for E_{top} since the action functionals are multi-valued. However, we will count solutions with weight $t^{E_{\text{top}}(u)}$ (which is a homotopy invariant of u relative to the ends), so we only need compactness for $\mathcal{R}^{d+1, \mathbf{P}}(\mathbf{x}; C) \subset \mathcal{R}^{d+1, \mathbf{P}}(\mathbf{x})$, which is the union of the components which have $E_{\text{top}} \leq C$. This follows as in the exact case (2) as we artificially imposed an a priori bound on $E(u)$ via $E(u) \leq E_{\text{top}}(u) \leq C$.

Remark. By 3.7, monotonicity ensures that moduli spaces of Floer solutions have well-behaved compactifications and eliminates the issues with multiple covers of the general theory [21].

3.15. The wrapped A_∞ -structure. We now define, for $d \geq 2$, the maps $\mu^{d, \mathbf{P}, \mathbf{w}}$:

$$CF^*(L_{d-1}, L_d; w_d H)[\mathbf{q}] \otimes \dots \otimes CF^*(L_0, L_1; w_1 H)[\mathbf{q}] \rightarrow CF^*(L_0, L_d; w_0 H)[\mathbf{q}]$$

First we determine the \mathbf{q}^0 coefficient of the image. An isolated $u \in \mathcal{R}^{d+1, \mathbf{P}}(x_0, x_1, \dots, x_d)$ will contribute $\pm t^{E_{\text{top}}(u)} x_0$ to $\mu^{d, \mathbf{P}, \mathbf{w}}(\mathbf{q}^{i_d} x_d \otimes \dots \otimes \mathbf{q}^{i_1} x_1)$ where recall

$$E_{\text{top}}(u) = \int_S u^* \omega - d(u^* H \wedge \gamma)$$

and where $i_k = 1$ for $k \in \mathbf{p}(F)$, $i_k = 0$ for $k \notin \mathbf{p}(F)$, and it will contribute 0 otherwise (for orientation signs, see [8, Sec.3.8] and the Remark below). The conditions on $\mathbf{p}(F)$ ensure that when $i_k = 1$ the formal variable \mathbf{q}^1 plays the algebraic role of remembering that there is a preferred point on the geodesic in S connecting z_0, z_k . For symmetry reasons [8, Lemma 3.7], there is a cancellation of the counts of pseudo-holomorphic maps involving auxiliary data where more than one preferred point lies on a geodesic, so we only need to keep track of whether there is or there isn't a preferred point on each geodesic (encoded by $\mathbf{q}^1, \mathbf{q}^0$). Encoding this information in the algebra is necessary, so that the breaking analysis for the family $\mathcal{R}^{d+1, \mathbf{p}} \rightarrow \mathcal{R}^{d+1}$ will indeed yield the A_∞ -relations for the μ^d defined below (see Remarks 3 and 4 below, or see [8, Equation (61)]).

The \mathbf{q}^1 coefficient of the image is determined by the requirement that the map $\mu^{d, \mathbf{p}, \mathbf{w}}$ respects the $\partial_{\mathbf{q}}$ operator as follows, for $d \geq 2$:

$$\partial_{\mathbf{q}} \mu^{d, \mathbf{p}, \mathbf{w}}(c_d \otimes \cdots \otimes c_1) = \sum_{k=1}^d (-1)^{\sigma(c)_{k+1}^d} \mu^{d, \mathbf{p}, \mathbf{w}}(c_d \otimes \cdots \otimes \partial_{\mathbf{q}} c_k \otimes \cdots \otimes c_1),$$

where $\sigma(\cdot)$ is defined in 2.1. From this equation, it follows that also the \mathbf{q}^1 coefficient of $\mu^{d, \mathbf{p}, \mathbf{w}}(\mathbf{q}^{i_d} x_d \otimes \cdots \otimes \mathbf{q}^{i_1} x_1)$ is now determined (for an example, see [8, Ex.3.14]). The geometrical meaning of $\partial_{\mathbf{q}}$ -equivariance is explained in Remark 3 below. Finally, recall that $\mathbf{q}^2 = 0$, so we have determined the image.

Summing up all $\mu^{d, \mathbf{p}, \mathbf{w}}$ as \mathbf{p}, \mathbf{w} vary, we obtain the A_∞ -operations for $d \geq 2$:

$$\mu^d : CW^*(L_{d-1}, L_d) \otimes \cdots \otimes CW^*(L_0, L_1) \rightarrow CW^*(L_0, L_d).$$

Remark. *The sum is weighted by signs [8, Equation (75)]: $\mu^{d, \mathbf{p}, \mathbf{w}}(\mathbf{q}^{i_d} x_d \otimes \cdots \otimes \mathbf{q}^{i_1} x_1)$ is weighted $(-1)^\dagger$ where $\dagger = \sum_{j=1}^d j(\|x_j\| + 1) + \sum_{j \in F} \sigma(x)_{j+1}^d$ (the second sum is the Koszul sign expected from viewing \mathbf{q} as an operator of degree -1 acting from the left, the first sum is as in [2, (4.8)]).*

Remark about symmetry. *In view of the symmetry argument mentioned above, it is always understood that, when “summing up”, we only let $\{\mathbf{p}, F\}$ vary among inclusions of subsets*

$$\mathbf{p} : F \subset \{1, \dots, d\}.$$

For example, the first picture in Figure 1 contributes to the sum, whereas if we swapped the labels ϕ_1, ϕ_2 then the picture would not be part of the sum (by symmetry, the count of those solutions would contribute the same except for a reversal in orientation sign [8, Lemma 9.1]). At a more technical level, consider the associahedra in [8, Figure 2 and 3]: the bottom edge, which involves a bubbling with two preferred points on one geodesic, seems like a legitimate contribution to the A_∞ -equations. The key is that the preferred points are labelled, and they can converge at different speeds to the puncture where breaking occurs. But the auxiliary data (in particular the β_f) are constructed to be independent of the choice of labelling. So the broken Floer solutions contributing to the bottom edge of that associahedron will arise in pairs, depending on the order in which we place the labels ϕ_f . The symmetry argument implies that these pairs of contributions will cancel for orientation reasons (essentially because we are reordering the tensor of two orientation lines associated to the two preferred points). A genericity argument [8, Theorem 3.6] rules out that the two preferred points coincide, so the solutions come in genuine pairs.

We conclude with two technical remarks relating to gluing:

Remark 3. *In Remark 1 of 3.11 we described the gluings of the first two discs S^-, S^+ in Figure 1 for four choices of “glued” data \mathbf{p} . Let us pretend ϕ_7 is not present. The discs S^-, S^+ define two operators φ^-, φ^+ , whose \mathbf{q}^0 -output is non-zero only if we input terms of the*

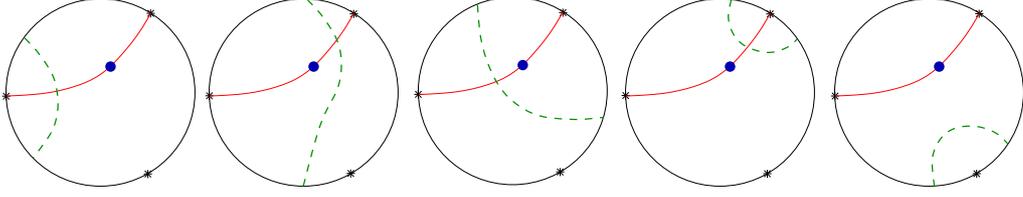


FIGURE 2. $\mu^1 \mu^2(\mathbf{q}x_2, x_1) + (-1)^{\|x_1\|} \mu^2(\mu^1 \mathbf{q}x_2, x_1) + \mu^2(\mathbf{q}x_2, \mu^1 x_1) = 0$.

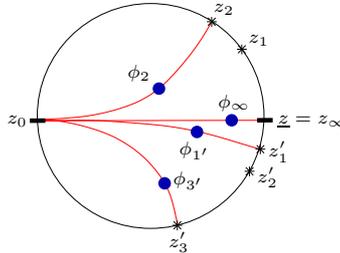
form $(\mathbf{q}y_2, \mathbf{q}y_1)$, $(\mathbf{q}x_4, x_3, x_2, \mathbf{q}x_1)$ respectively. The \mathbf{q}^0 -outputs x_0^-, x_0^+ of the two operators in these respective cases are then a count of isolated solutions for the relevant data. Gluing means composing operators, but S^- requests a \mathbf{q}^1 -input at the gluing puncture (namely $\mathbf{q}y_1$). The \mathbf{q}^1 -output of φ^+ is determined by $\partial_{\mathbf{q}}$ -equivariance: for example $\varphi^+(\mathbf{q}x_4, \mathbf{q}x_3, x_2, \mathbf{q}x_1) = \pm \mathbf{q}x_0^+$ (ignore signs for simplicity). So in the composition, we take $\mathbf{q}y_1 = \pm \varphi^+(\mathbf{q}x_4, \mathbf{q}x_3, x_2, \mathbf{q}x_1) \pm \varphi^+(\mathbf{q}x_4, x_3, \mathbf{q}x_2, \mathbf{q}x_1)$: those correspond to two of the four gluing choices (the other two involve $\mathbf{q}^2 = 0$). For example, $\varphi^+(\mathbf{q}x_4, \mathbf{q}x_3, x_2, \mathbf{q}x_1)$ corresponds to the choice $\mathbf{p}(1) = 3$, and it arises from the breaking of Figure 1 (Right). In particular, at the level of composing operators, each term of the composition corresponds to precisely one type of breaking. So the choices of possible “glued” data arising from the auxiliary data are consistent with the possible breakings of a family of solutions. So distinct operator composites do correspond to distinct breakings.

Remark 4. We now explain the A_∞ -equations for the μ^d in an example, ignoring signs. Recall $\mu^1(x) = \mathfrak{d}x$ and $\mu^1(\mathbf{q}x) = (\partial_{\mathbf{q}} + \mathfrak{K}\partial_{\mathbf{q}})(\mathbf{q}x) + (\mathbf{q}\mathfrak{d}\partial_{\mathbf{q}})(\mathbf{q}x)$. Consider the A_∞ -equation $\mu^1 \circ \mu^2 + \mu^2(\mu^1, \cdot) + \mu^2(\cdot, \mu^1) = 0$. Evaluate this equation on $(\mathbf{q}x_2, x_1)$. First, the \mathbf{q}^1 -coefficient is: $\mathfrak{d}\mu^2(x_2, x_1) + \mu^2(\mathfrak{d}x_2, x_1) + \mu^2(x_2, \mathfrak{d}x_1)$, which is the usual A_∞ -equation without \mathbf{q} 's coming from the 3 possible cuts of a 3-punctured discs with no geodesics. Secondly, the \mathbf{q}^0 -coefficient: this has 7 terms of which two cancel (the cancelling terms are $\partial_{\mathbf{q}}\mu^2(\mathbf{q}x_2, x_1)$, $\mu^2(x_2, x_1)$ arising respectively from $\mu^1\mu^2(\mathbf{q}x_2, x_1)$, $\mu^2(\mu^1(\mathbf{q}x_2), x_1)$, ignoring signs). This leaves the five terms $\mathfrak{d}\mu^2(\mathbf{q}x_2, x_1) + \mathfrak{K}\mu^2(x_2, x_1) + \mu^2((\mathfrak{K}\partial_{\mathbf{q}})\mathbf{q}x_2, x_1) + \mu^2((\mathbf{q}\mathfrak{d}\partial_{\mathbf{q}})\mathbf{q}x_2, x_1) + \mu^2(\mathbf{q}x_2, \mathfrak{d}x_1)$. These are the breakings coming from the cuts shown in Figure 2. In general, $\mathfrak{K}\partial_{\mathbf{q}}$ corresponds to the disc with 2 punctures, 1 geodesic, 1 preferred point; $\mathbf{q}\mathfrak{d}\partial_{\mathbf{q}}$ arises as an incoming disc with 2 punctures, 1 geodesic, no preferred point (i.e. the \mathfrak{d} disc with a new geodesic drawn).

3.16. The diagonal bimodule composition maps $\mu^{r|s}$. Abbreviate by $\mathcal{A} = \mathcal{W}(E)$ the A_∞ -category, and by $\mathcal{M} = \mathcal{W}(E)$ the diagonal bimodule (see 2.2). The composition maps

$$\mu^{r|s} : \mathcal{W}(E)(L_r, \dots, L_0) \otimes \mathcal{M}(L_0, L'_0) \otimes \mathcal{W}(E)(L'_0, \dots, L'_s) \rightarrow \mathcal{M}(L_r, L'_s),$$

for $r, s \geq 0$ are determined by 2.2. For example: $\mu^{0|0} = -\mu^1$ is the Floer differential. So the moduli spaces of domains with auxiliary data are the same as those for the maps μ^d . We do have an artistic preference in how we represent the parametrized domains pictorially:



for μ^d we fixed the output at -1 leaving a $PSL(2, \mathbb{R})/\mathbb{R}$ -reparametrization freedom, now we prefer to fix two of the boundary punctures, drawn as dashes: the module output at -1 and

the module input at $+1$ leaving a residual \mathbb{R} -translational freedom (fixing ± 1). In the picture, the map $\mathbf{p} : F \subset \{3', 2', 1', \infty, 1, 2\}$ is defined by the inclusion of the subset $\{3', 1', \infty, 2\}$.

3.17. Wrapped category with local systems. Let $\Lambda_0 \subset \Lambda$ be the subring of series with only non-negative powers of t (see 3.3). Let $\Lambda_0^\times \subset \Lambda_0$ be the multiplicative group of units: the series with non-zero t^0 -coefficient. Λ_0 is a local ring with maximal ideal $\Lambda_+ = \ker(\text{evaluation } \Lambda_0 \rightarrow \mathbb{K} \text{ of the } t^0\text{-coefficient})$ consisting of the series with only strictly positive powers of t .

The A_∞ -category $\underline{\mathcal{W}}(E)$ equips objects $L \in \text{Ob}(\mathcal{W}(E))$ with additional data: a local system of coefficients $\underline{\Lambda}^L$ with fibre Λ and structure group Λ_0^\times . So:

- (1) $\underline{\Lambda}_x^L \cong \Lambda$ over each $x \in L$;
- (2) for any path γ in L from x to y , there is an isomorphism $P_\gamma^L : \underline{\Lambda}_x^L \rightarrow \underline{\Lambda}_y^L$ given by multiplication by an element in Λ_0^\times depending only on $[\gamma] \in \pi_1(L, x, y)$;
- (3) concatenating paths yields composed P^L 's.

For such L_0, L_1 and Hamiltonian wH the Floer complex is now

$$\underline{CF}^*(L_0, L_1; wH) = \bigoplus \text{Hom}_\Lambda(\underline{\Lambda}_{x(0)}^{L_0}, \underline{\Lambda}_{x(1)}^{L_1})$$

summing over all integer X -chords x of weight w , and a Floer solution $u : (\mathbb{D} \setminus \{\pm 1\}, \partial\mathbb{D}) \rightarrow (E, L_0 \cup L_1)$ asymptotic to Hamiltonian chords x, y at $-1, +1$, which contributed $\pm t^{\omega[u]}x$ to the old differential $\mathfrak{d}y$, now contributes

$$\text{Hom}_\Lambda(\underline{\Lambda}_{y(0)}^{L_0}, \underline{\Lambda}_{y(1)}^{L_1}) \ni \phi_y \mapsto \pm t^{\omega[u]} P_{\gamma_1}^{L_1} \circ \phi_y \circ P_{\gamma_0}^{L_0} \in \text{Hom}_\Lambda(\underline{\Lambda}_{x(0)}^{L_0}, \underline{\Lambda}_{x(1)}^{L_1}),$$

where γ_0, γ_1 are the paths in L_0, L_1 swept by $u|_{\partial\mathbb{D}}$ along the oriented arcs in $\partial\mathbb{D}$ connecting -1 to $+1$, and $+1$ to -1 . This defines the new differential $\underline{\mathfrak{d}}$.

Equivalently, one considers flat Λ_0^\times -connections on the trivial line bundle $L \times \Lambda$ over L . So the holonomy is a homomorphism $\pi_1(L) \rightarrow \Lambda_0^\times$, which determines a class $[b_L] \in H^1(L; \Lambda_0^\times)$, and the parallel transport $P_\gamma^L : \{x\} \times \Lambda \rightarrow \{y\} \times \Lambda$ for a path γ from x to y is given by multiplication by $b_L[\gamma] \in \Lambda_0^\times$ (evaluation of cocycles on chains). Independence of the homotopy class γ is recorded by the fact that b_L is a cocycle.

The Λ -module $\underline{CF}^*(L_0, L_1; wH) = CF^*(L_0, L_1; wH) = \bigoplus \Lambda x$ is generated by X -chords of weight w , except now the above u contributes $\pm t^{\omega[u]} b_{L_0}[\gamma_0] b_{L_1}[\gamma_1] x$ to the new differential $\underline{\mathfrak{d}}y$. For zero b_{L_0}, b_{L_1} , which is a trivial local system, this recovers the usual CF^* . If we consider a 1-family of discs u as above with fixed asymptotics x, y , then the homotopy class of the arcs γ_0, γ_1 is constant in the family, so $t^{\omega[u]} b_{L_0}[\gamma_0] b_{L_1}[\gamma_1]$ is constant. Thus $\underline{\mathfrak{d}}^2 = 0$ if and only if the old $\mathfrak{d}^2 = 0$. For the analogue $\underline{\mathfrak{k}}$ of 3.9: $\pm t^{\omega[u]} b_{L_0}[\gamma_0] b_{L_1}[\gamma_1] x$ is a contribution to $\underline{\mathfrak{k}}y$ precisely when $\pm t^{\omega[u]} x$ is a contribution to $\mathfrak{k}y$. Together with $\underline{\mathfrak{d}}$ this defines $\mu_{\underline{\mathcal{W}}(E)}^1$.

Similarly, $\mu_{\underline{\mathcal{W}}(E)}^n$ now counts solutions u with weight $\pm t^{\omega[u]} b_{L_0}[\gamma_0] \cdot b_{L_1}[\gamma_1] \cdots b_{L_n}[\gamma_n]$ where γ_i are the images via u of the oriented arcs in $\partial\mathbb{D}$ which land in L_i .

Mimicking 3.7, the new $\underline{\mathfrak{m}}$'s now depend on $[b_L] \in H^1(L; \Lambda_0^\times)$:

$$\underline{\mathfrak{m}}_0(L, b_L) = \sum t^{\omega[\beta]} b_L[\partial\beta] \text{ev}_*[\mathcal{M}_1(\beta)] \in C_{\dim L}^{\text{lf}}(L; \underline{\Lambda}^L)$$

summing over the homotopy classes $\beta \in \pi_2(E, L)$ of Maslov index 2, where $\mathcal{M}_1(\beta) \subset \mathcal{M}_1$ corresponds to the disc bubbles u in class $[u] = \beta$ (see 3.7). As argued in 3.7, $\mathcal{M}_1(\beta)$ is a $\dim(L)$ -cycle so $\text{ev}_*[\mathcal{M}_1(\beta)]$ is a multiple of $[L]$, and finally we define $\underline{\mathfrak{m}}_0(L, b_L) \in \Lambda_0$ by

$$\underline{\mathfrak{m}}_0(L, b_L) = \underline{\mathfrak{m}}_0(L, b_L) [L].$$

Just like in the absence of local systems, $\underline{\mathcal{W}}(E)$ splits up according to m_0 -values (see 3.7). Even though typically the m_0 -values have changed by introducing local systems (i.e. $m_0(L) \neq \underline{\mathfrak{m}}_0(L, b_L)$), the possible m_0 -values are still constrained to lie in the spectrum of eigenvalues

of $c_1(TE) * \cdot$ because Lemma 3.2 holds also in the presence of local systems. The advantage of local systems is that a wider range of eigenvalues can now arise geometrically as m_0 -values.

Local systems will only appear in the toric applications in Section 12, and we suppress them from the notation elsewhere.

4. SYMPLECTIC COHOMOLOGY

4.1. Hamiltonian Floer cohomology. We refer to Salamon [38] for a precise definition of the Floer cohomology $HF^*(H^B)$ of a closed monotone symplectic manifold (B, ω_B) , for a Hamiltonian $H^B : B \rightarrow \mathbb{R}$. [38, Lecture 3] discusses the role of monotonicity. We recall that at the chain level, $CF^*(H^B)$ is generated, over the Novikov field Λ , by all 1-periodic Hamiltonian orbits of H^B . This is not entirely accurate, as for time-independent Hamiltonians the 1-orbits are not isolated: they arise in S^1 -families, due to the freedom in the choice of basepoint. So one must make a time-dependent perturbation of H^B , discussed in 4.5. The differential $\partial : CF^*(H^B) \rightarrow CF^{*+1}(H^B)$ is defined as follows on generators: there is a contribution

$$\pm t^{E_{\text{top}}(u)} x_{\pm}$$

to $\partial(x_{\pm})$ for each isolated (up to \mathbb{R} -translation) map $u : \mathbb{R} \times S^1 \rightarrow B$ solving Floer's equation

$$\partial_s u + J^B(\partial_t u - X_{H^B}) = 0,$$

with asymptotic conditions $u(s, t) \rightarrow x_{\pm}(t)$ as $s \rightarrow \pm\infty$, where J^B is a fixed almost complex structure on B compatible with ω_B . Above, $E_{\text{top}}(u) \geq 0 \in \mathbb{R}$ is defined by:

$$E_{\text{top}}(u) = \frac{1}{2} \int_S \|du - X_{H^B} \otimes dt\|^2 ds \wedge dt = \int_S u^* \omega - d(u^*(H^B)) \wedge dt = \int_S u^* \omega + H^B(x_-) - H^B(x_+).$$

Up to isomorphism, the $HF^*(H^B)$ do not depend on the choice of H^B, J^B . Indeed, a homotopy H_s^B, J_s^B which is constant for $|s| \gg 0$, induces a *continuation map* $HF^*(H_{+\infty}^B, J_{+\infty}^B) \rightarrow HF^*(H_{-\infty}^B, J_{-\infty}^B)$ by counting solutions u of Floer's equation above where now J_s^B, H_s^B depend on the coordinate $s \in \mathbb{R}$ (and we no longer have an \mathbb{R} -reparametrization action). This continuation map is always an isomorphism (the inverse is defined using the ‘‘reversed homotopy’’ H_{-s}^B, J_{-s}^B). When H^B is C^2 -small and Morse, and H^B, J^B are time-independent, then $CF^*(H^B)$ reduces to the Morse complex for H^B , in particular all 1-periodic Hamiltonian orbits are constant, and so by invariance we have an isomorphism

$$QH^*(B) \cong HF^*(H^B),$$

which by PSS [32] (see 4.4) intertwines the quantum product and the pair-of-pants product.

4.2. Symplectic cohomology. For symplectic manifolds E conical at infinity (see 3.1) $HF^*(H^E)$ depends on the choice of $H^E : E \rightarrow \mathbb{R}$ because the larger H^E is on the conical end the likelier it is that there are many 1-periodic Hamiltonian orbits there. The continuation maps are no longer isomorphisms, because they can no longer be reversed: only certain homotopies H_s^E with $\partial_s H_s^E \leq 0$ are guaranteed to be well-defined. To construct an invariant one needs to choose a class of Hamiltonians which grow in a controlled way at infinity. We consider H^E which at infinity are linear of positive slope in the radial coordinate R (see 3.1). One then takes a direct limit over the continuation maps which increase the slope of H^E at infinity: $HF^*(H_+^E) \rightarrow HF^*(H_-^E)$ where $\partial_s H_s^E \leq 0$. This defines the *symplectic cohomology*

$$SH^*(E) = \varinjlim HF^*(H^E),$$

first constructed in the non-exact setup by Ritter [33, 35]. We refer the reader also to the original construction in the exact setup due to Viterbo [48]; the survey by Seidel [42]; and the construction of algebraic structures such as the pair-of-pants product by Ritter [34]. We

follow the conventions of [35], except we use the simpler Novikov field defined in 3.3 as we work in the monotone setting.

When H^E is C^2 -small on E^{in} , Morse, time-independent, and of small enough slope on the conical end, the chain complex $CF^*(H^E)$ reduces to the Morse complex for E . Since this is part of the direct limit, there is a canonical map (which need not be injective/surjective)

$$c^* : QH^*(E) \rightarrow SH^*(E).$$

Remark 4.1. c^* is a ring homomorphism, so it splits into eigensummands $c^* : QH^*(E)_\lambda \rightarrow SH^*(E)_\lambda$ indexed by the eigenvalues of quantum product by $c_1(TE)$ acting on $QH^*(E)$.

Warning (Grading). $SH^*(E)$ is typically not \mathbb{Z} -graded. Floer cohomology is graded by the Conley-Zehnder index, which is a $\mathbb{Z}/2N$ -grading where $N\mathbb{Z} = c_1(TE)(\pi_2(E))$.

For the purpose of defining the open-closed string map in Section 6, it is actually convenient to view also $SH^*(E)$ as a homotopy direct limit (recall 3.10). Namely, define

$$SC^*(E) = \bigoplus_{w=1}^{\infty} CF^*(wH)[\mathbf{q}]$$

with differential $\mu^1 = \mu_{SC^*(E)}^1$ defined by the same formula as in Definition 3.4,

$$\mu^1(x + \mathbf{q}y) = (-1)^{|x|} \mathfrak{d}x + (-1)^{|y|} (\mathbf{q}\mathfrak{d}y + \mathfrak{K}y - y) \quad (4.1)$$

where $\mathfrak{d} : CF^*(wH) \rightarrow CF^{*+1}(wH)$ is the differential defining Hamiltonian Floer cohomology and $\mathfrak{K} : CF^*(wH) \rightarrow CF^*((w+1)H)$ is the Floer continuation map analogous to 3.9. Explicitly, the auxiliary β_f form that defines these maps can be chosen to be of the form

$$\beta_f = c(s) dt \text{ where } c : \mathbb{R} \rightarrow [0, 1], c'(s) \leq 0, \text{ with } c = 1 \text{ for } s \ll 0 \text{ and } c = 0 \text{ for } s \gg 0 \quad (4.2)$$

and the natural choice would be to pick $c(s) \neq 0, 1$ only near the circular s -slice of $\mathbb{R} \times S^1$ which contains the preferred point ϕ_f (although this is not strictly necessary). Recall in general we do not necessarily need the support of $d\beta_f$ to lie near ϕ_f . Notice the above β_f are constructed S^1 -invariantly, and we will use these also for the A_∞ -structure in 4.6.

We define the product μ^2 on $SC^*(E)$ and the A_∞ -operations μ^n later, in 4.6.

Convention. In 3.5, we chose H to have slope 1 at infinity, but we could have rescaled H by a constant $\epsilon > 0$ (once and for all). From now on, we assume H is Morse and C^2 -small on E^{in} , and linear on the conical end with slope $\epsilon > 0$ smaller than the minimal Reeb period of the closed Reeb orbits in Σ . Then the only 1-periodic orbits of H lie in E^{in} and are constant (whilst wH typically will have non-constant 1-periodic orbits on the conical end for large w), and the complex $CF^*(H)$ becomes the Morse complex for H , so $HF^*(H) \cong QH^*(E)$.

Lemma 4.2. $SH^*(E) \cong H^*(SC^*(E); \mu^1)$ and the canonical map $c^* : HF^*(H) \rightarrow SH^*(E)$ arises as the inclusion of the subcomplex $CF^*(H) \rightarrow SC^*(E)$, where $CF^*(H)$ is the \mathbf{q}^0 -part of the weight $w = 1$ summand of $SC^*(E) = \bigoplus CF^*(wH)[\mathbf{q}]$.

Proof. The first claim follows by the same arguments as in 3.10, the second claim follows by definition (and required the rescaling in the Convention). \square

4.3. Contractible vs Non-contractible orbits. The Floer differential preserves the free homotopy class of the orbits, so one can restrict the chain complexes $CF^*(H^B)$ and $SC^*(E)$ to only contractible orbits. Denote their cohomologies by $HF_0^*(H^B)$ and $SH_0^*(E)$. By invariance, and considering small H^B as mentioned in 4.1, we deduce $HF_0^*(H^B) = HF^*(H^B)$ for any H^B . Similarly when H^E has small slope, such as the H in 4.2, we have $HF^*(H^E) =$

$HF_0^*(H^E)$, however this may not hold for general H^E . Moreover $SH_0^*(E) \subset SH^*(E)$ is a subalgebra containing $c^*(QH^*(E))$ and the unit $1 = c^*(1)$, and there is a projection map $SH^*(E) \rightarrow SH_0^*(E)$ (which arises from a homomorphism of chain complexes, but it is not an algebra homomorphism). When $\pi_1(E) = 1$, $SH^*(E) = SH_0^*(E)$.

4.4. The PSS-isomorphisms. There are mutually inverse PSS-isomorphisms [32]

$$\psi^+ : QH^*(B) \rightarrow HF_0^*(H^B) \quad \psi^- : HF_0^*(H^B) \rightarrow QH^*(B)$$

defined by counting spiked planes as follows. Parametrize $\mathbb{C}^* = \{z \in \mathbb{C} : z \neq 0\}$ by $z = e^{-2\pi(s+it)}$ for $(s, t) \in \mathbb{R} \times S^1$. Consider solutions $u : \mathbb{C}^* \rightarrow B$ of Floer's equation

$$\partial_s u + J^B(\partial_t u - c(s)X_{H^B}) = 0$$

where $c : \mathbb{R} \rightarrow [0, 1]$ is a decreasing cut-off function which equals 1 for $s \leq -2$ and equals 0 for $s \geq -1$. Notice u is J^B -holomorphic near the puncture $z = 0$.

The *geometric energy* of u is $E_{\text{geom}}(u) = \int_{\mathbb{C}^*} |\partial_s u|_{J^B}^2 ds \wedge dt$. If it is finite, then u converges to a Hamiltonian orbit of H^B as $s \rightarrow -\infty$, and u extends J^B -holomorphically over $z = 0$.

For a generic cycle $\alpha \in C_{2 \dim_{\mathbb{C}} B - *}(B)$, define $\psi^+(\alpha)$ by counting 0-dimensional moduli spaces of such finite energy solutions u with $u(0) \in \alpha$.

The definition of ψ^- is similar, except parametrize \mathbb{C}^* by $z = e^{+2\pi(s+it)}$. In this case, solutions u converge to a Hamiltonian orbit as $s \rightarrow +\infty$, and u is J^B -holomorphic for $s \ll 0$ and extends J^B -holomorphically over the puncture $z = 0$.

In [35, Theorem 37], it is proved that there are mutually inverse PSS-isomorphisms for E ,

$$QH^*(E) \rightarrow HF_0^*(H^E) \quad HF_0^*(H^E) \rightarrow QH^*(E),$$

where $H^E : E \rightarrow \mathbb{R}$ is radial at infinity, $H^E = h^E(R)$, such that h^E is C^2 -small and concave (in particular H^E is globally bounded). Homotoping H to H^E , where H is as in 4.2, defines a continuation map $HF_0^*(H^E) \rightarrow HF^*(H)$ which is an isomorphism and the composite $QH^*(E) \rightarrow HF_0^*(H^E) \rightarrow HF^*(H) \rightarrow SH^*(E)$ equals the c^* of 4.2 (for details, see [34]).

4.5. Floer datum to break the S^1 -symmetry. Choosing H generically, we may assume the 1-orbits of H are transversally non-degenerate, but we need a time-dependent perturbation to make them non-degenerate in the S^1 -direction. The Floer equation is $(du - X_H \otimes dt)^{0,1} = 0$, and the time-dependent perturbation corresponds to replacing this equation by

$$(du - X_H \otimes dt - X_{H_{\text{pert}}} \otimes c_-(s)dt - X_{H_{\text{pert}}} \otimes c_+(s)dt)^{0,1} = 0$$

where the *Floer datum* H_{pert} is a small time-dependent Hamiltonian supported near the 1-orbits for H , ensuring that the 1-orbit breaks into two isolated 1-orbits [18, Lemma 2.1]. The cut-off functions $c_-, c_+ : \mathbb{R} \rightarrow [0, 1]$ are supported respectively in the region $s \ll 0$ and $s \gg 0$. By standard arguments in Floer theory, one can show that these additional choices do not affect $HF^*(H)$ up to isomorphism. In the non-compact setup of 4.2, we do not need to impose that $c_{\pm} \leq 0$ (the analogue of the condition $d\gamma \leq 0$). This is because H_{pert} is compactly supported, so the maximum principle still applies at infinity since we have not changed H there. So we avoid the complications [25, Definition 4.10] arising when Hamiltonians grow quadratically. We now justify the analogue of the a priori energy estimate in 3.14. Suppose we are counting Floer solutions u of $(du - X_H \otimes \gamma)^{0,1} = 0$ defined on a punctured Riemann surface S . Consider an end of S where this equation becomes $(du - X_H \otimes w dt)^{0,1} = 0$, then we need to modify $X_H \otimes w dt = X_{wH} \otimes dt$ so that it becomes $X_{w(H+H_{\text{pert},w})} \otimes dt$, where $H_{\text{pert},w} : S^1 \times E \rightarrow \mathbb{R}$ is a small time-dependent function depending on the weight w and supported near the 1-orbits for wH . We can extend this to a function $H_{\text{pert}} : S \times E \rightarrow \mathbb{R}$ supported near the ends of S , where it equals $c_{\pm}(s)H_{\text{pert},w}$ if the end has weight w and sign

\pm . Finally, the perturbed Floer equation becomes: $(du - X_{\tilde{H}} \otimes \gamma)^{0,1} = 0$ where $\tilde{H} = H + H_{\text{pert}}$ (which now depends on S). Then the a priori energy estimate in 3.14 becomes

$$\begin{aligned} E(u) &\equiv \int_S u^* \omega - u^* d\tilde{H} \wedge \gamma \\ &\leq E_{\text{top}}(u) + \sum_{\text{ends}} \text{measure}(\text{support}(c'_\pm)) \cdot \sup |c'_\pm| \cdot \|H_{\text{pert},w}\|_\infty \end{aligned}$$

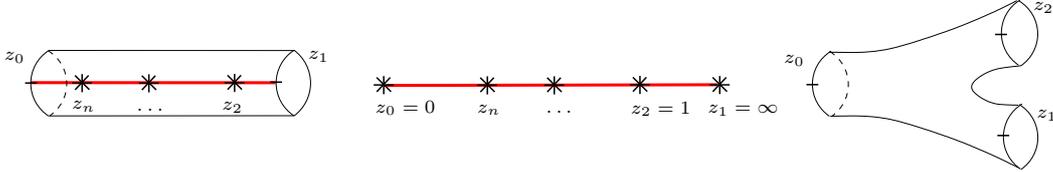
where $E_{\text{top}}(u) = \int_S u^* \omega - d(u^* \tilde{H} \wedge \gamma)$, and the support of c'_\pm is a compact subset of S whose measure is finite and independent of u . Above, we used that $\tilde{H} \geq 0$ and $d\gamma \leq 0$ and that on an end, for very large $|s|$, $\gamma = w dt$ and \tilde{H} is t -dependent but not s -dependent.

When we define chain level operations on $SC^*(E) = \oplus CF^*(wH)[\mathbf{q}]$, it is understood that:

- $CF^*(wH)$ in fact means $CF^*(w(H + H_{w,\text{pert}}))$ where $H_{w,\text{pert}}$ is the Floer datum chosen for the weight $w \geq 1$, as described above.
- Whenever we define a chain level map (e.g. the A_∞ -structure on $SC^*(E)$, or operations involving $SC^*(E)$), we will construct an auxiliary one-form γ determined by auxiliary forms α_k, β_f . It is always understood that the Floer equation $(du - X_H \otimes \gamma)^{0,1} = 0$ defining the counts of solutions $u : S \rightarrow E$, in fact needs to be perturbed, so $X_H \otimes \gamma$ is replaced by $X_{H+H_{\text{pert}}} \otimes \gamma$, which equals $X_{wH} \otimes dt + X_{H_{w,\text{pert}}} \otimes c_\pm(s)dt$ at the ends.
- The data H_{pert} depends on S , so this new auxiliary data also needs to be constructed universally and consistently, following the recipe of 3.13.

We will not mention these perturbations again, so that the notation stays under control.

4.6. The A_∞ -structure on $SC^*(E)$. Consider the moduli space of decorated domains as in the first picture below: cylinders $\mathbb{R} \times S^1$ with finitely many ordered positive punctures z_n, \dots, z_2 along the $t = 0$ line.



One can equivalently view this as the moduli space of punctured spheres, as in the second picture, where there is exactly one negative puncture z_0 and one distinguished positive puncture z_1 , and ordered positive punctures z_n, \dots, z_2 along the geodesic connecting z_0 to z_1 , where it is understood that the cylindrical ends at the punctures must have $t = 0$ lining up with the geodesic (each interior puncture carries an *asymptotic marker* which determines the $t = 0$ direction in the end parametrization). If there are three or more punctures, we can remove all parametrization freedom by considering only $\mathbb{C}P^1$ and fixing $z_0 = 0, z_2 = 1, z_1 = \infty$ (letting z_n, \dots, z_2 be ordered punctures on the real line segment $(0, 1) \subset \mathbb{C}P^1 = \mathbb{C} \cup \{\infty\}$).

If there are only two punctures, then there is a residual \mathbb{R} -translation reparametrization freedom. The third picture above is an equivalent way to view the 3-punctured case, $n = 1$, and is called a pair-of-pants. Roughly speaking, given Hamiltonian 1-orbits x_j , the x_0 -coefficient of $\mu^n(x_n, \dots, x_1)$ is a count of Floer solutions on the domain above, taking x_j to be the asymptotic condition at the puncture z_j . As usual, we allow the presence of preferred points ϕ_f lying on the geodesic line $t = 0$ (cylinder model) joining z_0 to a positive puncture z_{p_f} . It is always understood that the point ϕ_f belongs to the compactified cylinder where all punctures except z_0, z_{p_f} have been filled in. When constructing auxiliary data, it is always understood that the β_f forms are required to vanish at all positive punctures.

Rotational symmetry trick. We insist that the β_f forms are constructed S^1 -equivariantly

in the cylinder model, universally and consistently as in 3.13. (Note this means that the support of $d\beta_f$ is not located near the preferred point ϕ_f , but rather at best near a “preferred circle”.) So we do not need to constrain the ϕ_f to the $t = 0$ line, provided we work with moduli of domains: we only care about location of the ϕ_f up to rotating each ϕ_f individually.

The differential μ^1 counts solutions $u : \mathbb{R} \times S^1 \rightarrow E$ of Floer’s equation $(du - X \otimes \gamma)^{0,1} = 0$ where γ depends on the auxiliary data analogous to 3.11, except the ends are cylindrical rather than strip-like (as we replaced the time-interval $[0, 1]$ by S^1). So we have a closed form α which pulls back to dt in the cylindrical ends near the punctures; and a preferred point ϕ_f (if present) on the geodesic $t = 0$ corresponds to the translation $\mathbb{R} \times S^1 \rightarrow \mathbb{R} \times S^1$, $\phi_f \mapsto (0, 0)$. These translations allow one to consistently build 1-forms β_f with $d\beta_f \leq 0$, and by the above convention we build these S^1 -invariantly, for example as in (4.2).

We now discuss the construction of the pair-of-pants product μ^2 , so the case $n = 2$, but the same argument with additional punctures defines more generally the A_∞ -structure maps $\mu^n : SC^*(E)^{\otimes n} \rightarrow SC^*(E)$ for $n \geq 2$. In analogy with $\mu_{\mathcal{W}(E)}^2$, one constructs

$$\mu^2 = \mu_{SC^*(E)}^2 : SC^*(E) \otimes SC^*(E) \rightarrow SC^*(E),$$

except the domain is now the 3-punctured sphere (equivalently the cylinder with a positive puncture at $(s, t) = (0, 0)$), rather than a 3-punctured disc.

The two positive punctures 1 and ∞ receive inputs from $CF^*(w_1H)[\mathbf{q}]$, $CF^*(w_\infty H)[\mathbf{q}]$, whereas the negative puncture 0 receives the output from $CF^*(w_0H)[\mathbf{q}]$. The count of solutions on these domains, with auxiliary data, determine the \mathbf{q}^0 -contribution of

$$\mu^2 : CF^*(w_1H)[\mathbf{q}] \otimes CF^*(w_\infty H)[\mathbf{q}] \rightarrow CF^*(w_0H)[\mathbf{q}].$$

Finally, requiring μ^2 to be $\partial_{\mathbf{q}}$ -equivariant determines the \mathbf{q}^1 contributions.

As usual, the auxiliary form $\gamma = \sum w_k \alpha_k + \sum \beta_f$ which determines the Floer equation pulls back to $w_1 dt$, $w_\infty dt$, $w_0 dt$ in the parametrizations near 1, ∞ , 0 (where $w_0 = w_1 + w_\infty + |F|$). There are two closed forms α_1 , α_∞ associated to 1, ∞ satisfying $\alpha_\infty = 0$ near 1 and $\alpha_\infty = dt$ near 0 and ∞ ; $\alpha_1 = 0$ near ∞ , and $\alpha_1 = dt$ near 0 and 1. Along the $t = 0$ line in the cylinder model, there are two geodesics: one connecting 0 to 1, and the other connecting 0 to ∞ . These geodesics may carry preferred points ϕ_f , to which one associates S^1 -invariant β_f forms with $d\beta_f \leq 0$ which satisfy: $\beta_f = 0$ near 1, ∞ and $\beta_f = dt$ near 0.

As usual, the auxiliary data is built universally and consistently for all μ^n , $n \geq 1$, following the argument in 3.13. The bundle $\mathcal{A}ux$ of auxiliary data is over the compactification of the moduli space of all pairs-of-pants. The additional strata of the compactification arise because of the bubbling off of a twice-punctured sphere (i.e. a cylinder) at some end. As usual one constructs the auxiliary data inductively over the strata of the compactification so that the universal auxiliary data that we already constructed for μ^1 will be extended to define the universal auxiliary data for μ^2 . This is necessary to ensure that μ^2 will be a chain map (i.e. the cylinders which break off in a 1-family of pairs-of-pants, carry the correct universal auxiliary data chosen for μ^1). Just as in 3.13, the fibre of $\mathcal{A}ux$ is non-empty by de Rham theory (a proof can be found in [34, Appendix A]) and it is contractible, indeed convex, because one can linearly interpolate the data whilst preserving all the conditions. Linear interpolation of two sets of universal auxiliary data and a homotopy argument implies that the μ^2 structure on the cohomology $SH^*(E)$ does not depend on the auxiliary choices.

That μ^2 is a chain map follows by considering the possible breaking of a 1-dimensional moduli space with the above data. Namely, one cylinder can break off at one of the three punctures, carrying an \mathbb{R} -symmetry (this is an index 1 cylinder, and for index reasons there cannot be more than one cylinder breaking off). If the cylinder carries no preferred points then this count contributes to the \mathfrak{d} -part of μ^1 , if it carries one preferred point then this

contributes to the \mathfrak{K} -part of μ^1 . In the definition of μ^2 and in this breaking analysis, one can always forget configurations where there is more than one preferred point on a geodesic as the count of such configurations will cancel for the same reason as in [8, Lemma 3.7].

Technical Remark. *The symmetry argument [8, Lemma 3.7] relies on showing that the transposition of two preferred points lying on the same geodesic switches the orientation sign with which the isolated solution is counted [8, Lemma 9.1]. This works generally both for moduli spaces of punctured discs and for moduli spaces of punctured genus zero Riemann surfaces, using that the auxiliary data has been constructed universally and consistently.*

That μ^2 is graded-commutative on cohomology follows by considering a suitable family of domains in which one makes the punctures $1, \infty$ move in $\mathbb{CP}^1 \setminus \{0\}$ until their positions are interchanged (we need to break the S^1 -invariance of the β_f forms in this argument, and recall by the homotopy argument above that on cohomology we do not care whether the universal auxiliary data before/after the movement coincides).

4.7. Associativity and unitality. For general reasons, as in 3.10, the telescope construction for $SC^*(E)$ corresponds on cohomology to the direct limit construction $SH^*(E) = \varinjlim HF^*(H)$. In particular, \mathfrak{D} -cocycles $y \in HF^*(wH)$ generate $SH^*(E)$ (allowing w to vary), and cocycles $y, \mathfrak{K}y$ are identified on cohomology via the boundary of $\mathfrak{Q}y$, which encodes the fact that generators of the direct limit are identified if they are related by applying continuation maps \mathfrak{K} . Thus on cohomology μ^2 is determined by the usual pair-of-pants product $HF^*(w_1H) \otimes HF^*(w_\infty H) \rightarrow HF^*(w_0H)$ constructed as in [34]. Thus cohomological properties of μ^2 follow from [34] by general TQFT arguments, in particular: the product is associative and unital, where the cohomological unit for $SH^*(E)$ is the image $c^*(1)$ of 1 under $c^* : QH^*(E) \rightarrow SH^*(E)$. Indeed, $c^* : QH^*(E) \rightarrow SH^*(E)$ is a unital ring homomorphism (this was proved in [34, 35]). A clean way to prove the latter is carried out in 7.1, where we enlarge the chain complex $SC^*(E) \subset SC^*_\diamond(E)$ to also include a $QC^*(E)$ -summand, so that c^* becomes the inclusion of this summand at the chain level.

Associativity can also be proved directly, by verifying the A_∞ -relation relating μ^1, μ^2, μ^3 , obtained from the bubbling analysis for the domains in 4.6 for $n = 3$.

We sketch the argument from [34] that the product μ^2 is unital. The unit is represented by the image of 1 under the PSS isomorphism $QH^*(E) \rightarrow HF^*(H)$. This corresponds to a count of Floer solutions $\mathbb{C} \rightarrow E$ with a tautological intersection condition $[E]$ at the origin, using a form γ such as $\gamma = c(s) dt$ where $c : \mathbb{R} \rightarrow [0, 1]$ is an increasing cut-off function with $c = 0$ near $s = -\infty$ (so the Floer solution is J -holomorphic near the origin), $c = 1$ near $s = \infty$, and where $e^{2\pi(s+it)}$ parametrizes $\mathbb{C} \setminus \{0\}$. The particular choice of γ only affects the representative by a chain homotopy, so choices do not matter on cohomology. We call these solutions the *E-spiked discs*. Gluing the *E-spiked discs* onto a pair-of-pants at a $CF^*(H)$ input puncture, yields the Floer solutions defined on a cylinder which on cohomology determine the continuation map $HF^*(w_\infty H) \rightarrow HF^*(w_0 H)$. But this map represents the identity map in the direct limit $SH^*(E)$ (generators get identified via continuation maps), which shows that $\mu^2(1, \cdot)$ represents the identity map on $SH^*(E)$, and similarly for $\mu^2(\cdot, 1)$.

4.8. $SH^*(E)$ is a module. $SH^*(E)$ is an $SH^*(E)$ -module by acting by μ^2 , and a $QH^*(E)$ -module by using the ring homomorphism $c^* : QH^*(E) \rightarrow SH^*(E)$ and then acting by μ^2 .

5. THE COMPACT FUKAYA CATEGORY

5.1. The compact Fukaya category \mathcal{E} . Let E be a monotone symplectic manifold conical at infinity. Denote the (compact) Fukaya category by

$$\mathcal{A} = \mathcal{F}(E).$$

The objects form the subset of $\text{Ob}(\mathcal{W}(E))$ consisting of the *closed* monotone Lagrangians (and recall from 3.7 that we fix a common m_0 -value) . Since the Lagrangians are compact, the maximum principle from 3.14 forces all Floer solutions to lie in a compact subset determined by the Lagrangians, so the non-compactness of E is unproblematic and the constructions for $\mathcal{F}(E)$ also carry over to $\mathcal{F}(B)$ for closed monotone symplectic manifolds B .

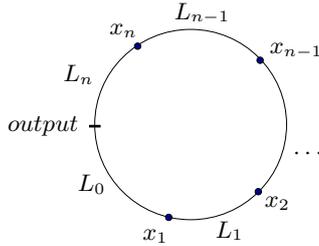
Lemma 5.1. $HF^*(L_0, L_1) \cong HW^*(L_0, L_1)$ for any closed Lagrangians $L_0, L_1 \in \text{Ob}(\mathcal{F}(E))$.

Proof. $HW^*(L_0, L_1) \cong \varinjlim HF^*(L_0, L_1; H)$ taking the direct limit over continuation maps $HF^*(L_0, L_1; H) \rightarrow HF^*(L_0, L_1; H')$ for linear Hamiltonians, where H' has larger slope at infinity than H . For L_0, L_1 closed, we can choose a cofinal sequence of Hamiltonians, so that each H' is obtained from the previous H by increasing the slope at infinity outside of a large compact set C containing L_0, L_1 . By the maximum principle, the Floer continuation solutions are trapped in C . But in C , the homotopy H_s from H' to H is s -independent, so Floer solutions would have an \mathbb{R} -symmetry in s if they were not constant. Since they are rigid they must be constant. So these continuation maps are identity maps at the chain level. \square

Suppose first that $L_0, \dots, L_n \in \text{Ob}(\mathcal{A})$ are pairwise transverse. Then no Hamiltonians arise in the definition of

$$\mu_{\mathcal{A}}^n : \mathcal{A}(L_n, \dots, L_0) \rightarrow \mathcal{A}(L_n, L_0).$$

Namely, $\mu_{\mathcal{A}}^n(x_n, \dots, x_1)$ counts J -holomorphic discs $u : S \rightarrow E$ modulo reparametrization (weighing counts by $t^{E_{\text{top}}(u)}$ as usual, with signs as in 3.15), where S is illustrated below.



On the boundary there are marked points (not punctures) which are free to move subject to being distinct and ordered anticlockwise; they receive inputs $x_j \in CF^*(L_{j-1}, L_j)$. We chose to fix the position of the output at -1 (which gives the output in $CF^*(L_0, L_n)$), and following 3.12 we count moduli of solutions, under the residual (real 2-dimensional) $PSL(2, \mathbb{R})/\mathbb{R}$ -reparametrization freedom which moves the input boundary marked points whilst fixing -1 .

Compared with the wrapped case, formally this would correspond to taking $w = 0$ and allowing only \mathbf{q}^0 terms (in the wrapped case we only allow weights $w \geq 1$). So moduli spaces are defined using the equation $du^{0,1} = 0$ rather than $(du - X \otimes \gamma)^{0,1} = 0$, and $CF^*(L_i, L_j)$ is generated by the intersections $L_i \cap L_j$. So there are no auxiliary α, β, γ forms, nor do we have strip-like end parametrizations. Just as for \mathfrak{d} in 3.6, the domain defining $\mu_{\mathcal{A}}^1$ has two boundary marked points which we fix at ± 1 leaving a residual \mathbb{R} -symmetry, and we count only non-constant solutions modulo this symmetry. The A_∞ -equations follow by a standard bubbling argument depicted in the picture of 2.1.

5.2. The Floer datum: non-transverse Lagrangians. The ordinary Lagrangian Floer complexes are related to the wrapped versions via

$$CF^*(L_i, L_j; H) \equiv CF^*(\varphi_H^1(L_i), L_j)$$

where φ_H^1 is the time 1 flow of X_H , and where one must choose suitable almost complex structures (for example [34] explains this). When L_i, L_j are not transverse, then $CF^*(L_i, L_j)$ is in fact *defined* as above using a compactly supported H so that $\varphi_H^1(L_i), L_j$ are transverse.

Setting up the A_∞ -category to keep track of the Hamiltonian perturbation data requires substantial work, carried out in detail by Seidel [41, Sec.(8e),(8f),(8j)]. So we only recall the key ideas. This perturbation was not required in the wrapped case as the weights $w \neq 0$.

Remark. For $L = L_i = L_j$ one can use a Morse-Bott model using $C_*(L) = CF^*(L, L)$ as carried out by Auroux [12]. However it would be a non-trivial task to implement such a Morse-Bott theory approach more generally for the whole A_∞ -category $\mathcal{F}(E)$.

In the Floer datum approach [41], loosely speaking, the A_∞ -operations μ_A^k “carry” the Hamiltonian perturbation data which is used to perturb the J -holomorphic curve equation which defines the moduli spaces, to achieve transversality. The equation is defined on the associahedron and a point on the associahedron determines the following data: boundary marked points; strip-like ends; Hamiltonians on each strip-like end such that the flow moves L to $\varphi_H^1(L)$ so that transversality is achieved. The Floer chain complexes $CF^*(L_0, L_1) \equiv CF^*(L_0, L_1; H_{L_0, L_1}, J_{L_0, L_1})$ are generated by the Hamiltonian chords using auxiliary data

$$H_{L_0, L_1} = H_{L_1, L_0} \quad J_{L_0, L_1} = J_{L_1, L_0}$$

called *Floer datum*. This is chosen generically once and for all for each pair L_0, L_1 (the generators then correspond bijectively to the transverse intersections $\varphi_H^1(L_0) \cap L_1$). Outside of a compact set containing L_0, L_1 we choose $H_{L_0, L_1} = 0$ and $J_{L_0, L_1} = J$ of contact type, so the maximum principle applies at infinity (this would not work in the wrapped setup for non-compact L_0, L_1 because a non-compactly supported H_{L_0, L_1} is usually needed to make $\varphi_H^1(L_0), L_1$ transverse). For transverse L_0, L_1 one can pick $H_{L_0, L_1} = 0$ and $J_{L_0, L_1} = J$. Even if one restricted the category to a subcollection L_i of pairwise transverse Lagrangians, one would still need a Floer datum for pairs of type L, L .

The *perturbation datum* is a universally consistent choice (over the associahedron) of interpolating data H, J for each operation μ_A^k with given Lagrangians L_0, \dots, L_k , which agrees with the Floer datum on the strip-like ends. So in fact we are still working, in general, with boundary punctures rather than boundary marked points; choices of strip-like ends; and the equation $(du - X \otimes \gamma)^{0,1} = 0$ rather than $du^{0,1} = 0$. However, now we can always choose $\gamma = dt$ near the strip-like ends of the domain of u , we do not require $d\gamma \leq 0$ and we can choose γ to be zero away from the strip-like ends (this would have been impossible in the wrapped setup where $d\gamma \leq 0$ is required by the maximum principle). On the strip-like ends we can ensure the equation is $(du - X_{H_{L_0, L_1}} \otimes c(s)dt)^{0,1} = 0$ for the given Floer datum, for a cut-off function $c : \mathbb{R} \rightarrow [0, 1]$ which is 1 for $|s| \gg 0$ and equals 0 for $|s|$ close to 0, so explicitly:

$$\partial_s u + J_s(\partial_t u - c(s)X_{H_{L_0, L_1}}) = 0,$$

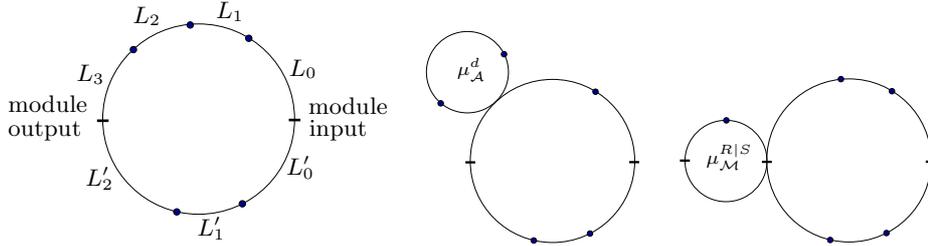
where $J_s = J_{L_0, L_1}$ for $|s| \gg 0$, $J_s = J$ for $|s|$ close to 0. An a priori energy estimate follows by the same argument as in 4.5. Unlike the wrapped setup, we do not need a telescope construction for the Floer chain complex (compare Lemma 5.1).

Remark. In the wrapped setup, for non-compact L_0, L_1 the maximum principle requires H_{L_0, L_1} to have non-zero slope ϵ at infinity. Stokes’s theorem and $d\gamma \leq 0$ force the output of μ^n to involve a Hamiltonian of slope at least $n\epsilon$. But for large n , $CF^*(L_0, L_1; nH_{L_0, L_1})$ is typically not quasi-isomorphic to $CF^*(L_0, L_1; H_{L_0, L_1})$. So the telescope construction is necessary.

5.3. The bimodule structure. Just as in 3.16, μ^d determines via 2.2 the composition maps

$$\mu_{\mathcal{M}}^{r|s} : \mathcal{A}(L_r, \dots, L_0) \otimes \mathcal{M}(L_0, L'_0) \otimes \mathcal{A}(L'_0, \dots, L'_s) \rightarrow \mathcal{M}(L_r, L'_s).$$

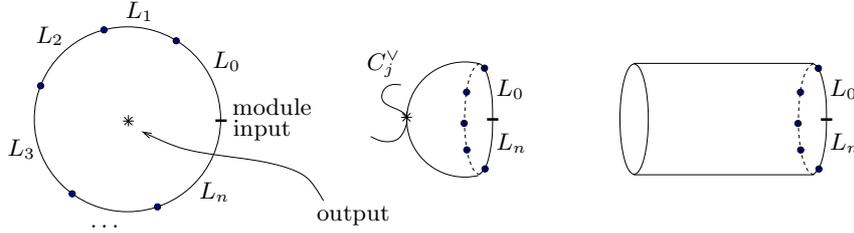
for the diagonal \mathcal{A} -bimodule $\mathcal{M} = \mathcal{A}$. Below, the first picture is $\mu_{\mathcal{M}}^{3|2} : \mathcal{A}(L_3, L_2, L_1, L_0) \otimes \mathcal{M}(L_0, L'_0) \otimes \mathcal{A}(L'_0, L'_1, L'_2) \rightarrow \mathcal{M}(L_3, L'_2)$, and the bubbling pictures show the A_∞ -relations.



5.4. **The open-closed string map.** To define $OC : CC_*(\mathcal{A}, \mathcal{M}) \rightarrow QC^*(E)$, we let

$$OC^n : \mathcal{M}(L_0, L_n) \otimes \mathcal{A}(L_n, \dots, L_0) \rightarrow QC^*(E)$$

be the count of maps as in the first picture (weighing counts by $t^{E_{\text{top}}(u)}$, with signs as in 6.1):



The underlying parametrized domains are the same as those used to define $\mu_{\mathcal{A}}^n$ except we have an additional marked point q_c in the interior, which determines the output in $QH^*(E)$.

Remark 5.2. *We cannot simply pull back the $PSL(2, \mathbb{R})$ -invariant auxiliary data used for $\mu_{\mathcal{A}}$ via the forgetful map, which forgets the new marked point q_c . Consider the moduli spaces of domains (modulo the $PSL(2, \mathbb{R})$ -reparametrization action, whereas later we will fix this choice by placing q_c at 0 and the module input at +1). If q_c converges to a boundary point which is not a puncture, then a J -holomorphic disc bubble appears which is attached to the main component at a boundary node rather than a boundary puncture. This is not consistent with the required breaking, which is a bubble carrying a module input puncture (counted by OC^0). So given the auxiliary data obtained via the forgetful map, we must deform the data when q_c gets close to a boundary point which is not a puncture (equivalently, if we fix the domain parametrization so that $q_c = 0$, this occurs when all boundary punctures converge to the module input puncture). As in 3.13, the universal and consistent construction of auxiliary data then requires the data on those bubbles to be the data inductively chosen for OC^0 .*

We will fix the parametrization of the domain, removing all reparametrization freedom, by fixing at 0 the interior marked point, and at +1 the positive puncture which receives the module input from $\mathcal{M}(L_0, L_n)$. The other n positive boundary punctures can move freely subject to being distinct and ordered; they receive the $\mathcal{A}(L_n, \dots, L_0)$ inputs of OC^n . Thus OC^n counts isolated maps $u : \mathbb{D} \setminus \{\text{punctures}\} \rightarrow E$ solving Floer's equation, subject to an intersection condition at 0. If all L_i, L_{i+1} were transverse, Floer's equation is simply the J -holomorphic equation $du^{0,1} = 0$, and we could use boundary marked points rather than boundary punctures. In general, however, as discussed in 5.2 on each strip-like end of a boundary puncture the equation is of type $(du - X \otimes \gamma)^{0,1} = 0$ using the relevant Floer datum, with $\gamma = dt$ near infinity and $\gamma = 0$ away from the end.

The intersection condition at 0 follows the usual routine: we fix a generic choice of intersection-dual cycles C_j^v in $C_*(E)$ for the basis of lf-cycles C_j that we chose to define $QC^*(E) = QH^*(E)$; then $OC(\text{input}) = \sum n_j C_j$ where n_j is the weighted count of solutions above (with boundary conditions determined by the input) which pass through C_j^v at 0. The

second picture above shows what the image of u looks like in the target manifold E .

Technical remark about transversality. *The condition that the solution “passes” through C_j^\vee is an intersection condition between C_j^\vee and the evaluation map at the centre of the disk. However, some care is needed because the evaluation does not sweep a pseudo-cycle (the moduli space has a boundary of real codimension one). If a disk or a sphere bubbles off (simple or not), then by monotonicity the index of the main component drops at least by 2, whilst remaining generically regular. Hence generically that main component will not hit C_j^\vee .*

To define

$$\text{OC}^n : \mathcal{M}(L_0, L_n) \otimes \mathcal{A}(L_n, \dots, L_0) \rightarrow \text{HF}^*(H^E)$$

we modify the above setup by making 0 a negative puncture (here negative means: with a choice of cylindrical parametrization $(-\infty, 0] \times S^1$ near the puncture, where the S^1 is parametrized by the argument of z and the puncture corresponds to the limit at $-\infty$). The equation on the cylindrical end near the interior puncture is also of type $(du - X \otimes \gamma)^{0,1} = 0$, indeed $(du - X_{H^E} \otimes c(s)dt)^{0,1} = 0$, where recall the Hamiltonian H^E at infinity is linear of positive slope. Explicitly this is the Floer continuation equation

$$\partial_s u + J(\partial_t u - c(s)X_{H^E}) = 0$$

in this cylindrical parametrization, where $c : \mathbb{R} \rightarrow [0, 1]$ is a decreasing cut-off function which equals 1 for $s \ll 0$ and equals 0 for s close to 0 (so the equation near $s = 0$ is $\partial_s u + J\partial_t u = 0$ and thus agrees with the condition that u is J -holomorphic away from the punctures). The third picture above is an equivalent way of viewing the map $\text{OC} : \text{CC}_*(\mathcal{A}, \mathcal{M}) \rightarrow \text{CF}^*(H^E)$: the domain is a half-infinite cylinder $(-\infty, 0] \times S^1$ with $n+1$ punctures on the circle $\{0\} \times S^1$, one of which is a distinguished puncture at $(0, 0)$, and we require the maps to satisfy the above Floer equation, so the left asymptotic circle $\{-\infty\} \times S^1$ plays the role of the original puncture at 0 from the first picture.

That OC is a chain map follows by considering the possible degenerations of the 1-dimensional moduli spaces: (1) if boundary punctures move together, we obtain $\mu_{\mathcal{A}}^d$ or $\mu_{\mathcal{M}}^{r|s}$ contributions; and (2) if the Floer trajectory breaks (on the region where $c(s) \equiv 1$), we obtain contributions described by the Floer differential $\mathfrak{d} : \text{CF}^*(H^E) \rightarrow \text{CF}^{*+1}(H^E)$ (in the above case, when we map OC into $\text{QC}^*(E)$, type (2) breaking would have corresponded to a limiting solution u which intersects the boundary chain ∂C_j^\vee , but in our setup C_j^\vee are cycles so this does not contribute). Thus we obtain the following two maps on homology:

$$\text{OC} : \text{HH}_*(\mathcal{A}) \rightarrow \text{QH}^*(E) \quad \text{OC} : \text{HH}_*(\mathcal{A}) \rightarrow \text{HF}^*(H^E). \quad (5.1)$$

Lemma 5.3. *The two maps in (5.1) commute via the PSS map $c^* : \text{QH}^*(E) \rightarrow \text{HF}^*(H^E)$.*

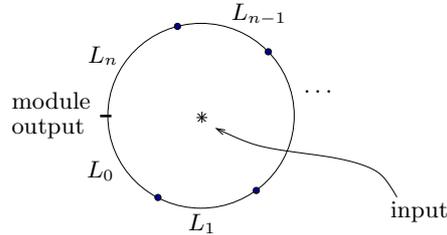
Proof. This is a standard gluing argument. At the chain level, a chain homotopy is involved because the glued data from the composite needs to be homotoped to the universal data which defines the second map. A clean way to implement this argument, is to place a new preferred point $\phi_f \in (0, \frac{1}{2}]$ on the geodesic $(0, +1)$ with an associated β_f form (as in the wrapped constructions), so that the new γ satisfies $\gamma = \beta_f = dt$ near 0 in the coordinates of the cylindrical end. As $\phi_f \rightarrow 0$ we can ensure that the support of β_f itself also converges to 0 (so $\beta_f = 0$ except near 0), and we can ensure that β_f becomes the S^1 -invariant form that we use to define the PSS map (compare the unitality argument at the end of 4.7). Consider a 1-family of Floer solutions on those domains. Suppose $\phi_f \rightarrow 0$. Then a Floer continuation cylinder breaks off at 0 involving Hamiltonians H^E at one end and 0 at the other ($\beta_f = 0$ on the main component). Since the solution has finite energy, a standard removal of singularities argument implies that the Floer continuation solution extends holomorphically over the puncture at the 0 end, and thus these cylinders are counted by a PSS map $\text{QC}^*(E) \rightarrow \text{CF}^*(H^E)$. The main

component in this degeneration has $\beta_f = 0$, so these define $OC : CC_*(\mathcal{A}, \mathcal{M}) \rightarrow QC^*(E)$. Other possible degenerations of the 1-family are: a disc bubble not carrying ϕ_f which is counted by the bar differential on $CC_*(\mathcal{A}, \mathcal{M})$, or a Floer cylinder not carrying ϕ_f which breaks off at 0 which is counted by the Floer differential. In both these cases the main component (which carries ϕ_f) is a rigid solution on a domain as above, and the count of these rigid solutions is the chain homotopy. If $\phi_f \rightarrow \frac{1}{2}$ and no bubbling or breaking occurs, we obtain solutions for the above domains with $\phi_f = \frac{1}{2}$ and this is (one of the possible definitions of) the map $OC : CC_*(\mathcal{A}, \mathcal{M}) \rightarrow CF^*(H^E)$. \square

5.5. **The closed-open string map.** In the notation of 2.13,

$$CO^d : QC^*(E) \rightarrow \prod_{n \geq 0} \prod \text{Hom}(\mathcal{A}(L_n, \dots, L_0), \mathcal{M}(L_n, L_0)[d])$$

counts Floer solutions (with weight $t^{E_{\text{top}}(u)}$ and sign as in 6.2) on the following domain.



Except for the following two modifications, this is the same setup as for OC (in particular, as in Remark 5.2, one can pull back the auxiliary data from $\mu_{\mathcal{A}}$ and then make a perturbation of the data when all boundary punctures converge to the module output).

- (1) to define $CO^d(C_j)$, the intersection condition at the marked point 0 is now with the 1f-cycle C_j (as opposed to the dual cycle C_j^V used for OC);
- (2) there is a negative boundary puncture at -1 for the module output in $\mathcal{M}(L_n, L_0)$ (instead of the module input puncture at $+1$ used for OC).

The proof that CO is a chain map is analogous to the OC case. So

$$CO : QH^*(E) \rightarrow HH^*(\mathcal{A}).$$

6. THE OPEN-CLOSED AND CLOSED-OPEN STRING MAPS

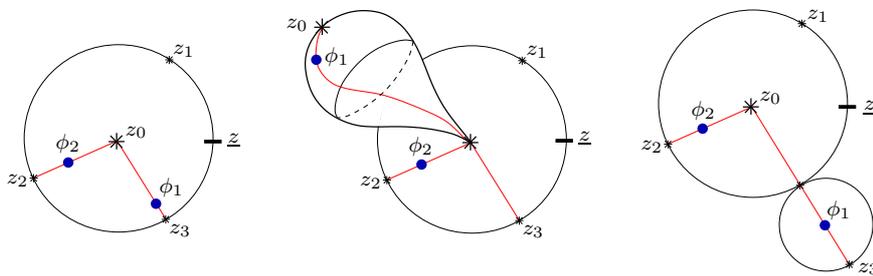


FIGURE 3. Left: $n = 3$, $F = \{1, 2\}$, $p_1 = 3$, $p_2 = 2$. Right: two broken configurations caused because the preferred point ϕ_1 converged to a puncture.

6.1. The open-closed string map. We will use the notation from 2.1-2.2: we denote the wrapped category and its diagonal bimodule by

$$\mathcal{A} = \mathcal{W}(E) \quad \mathcal{M} = \mathcal{A} = \mathcal{W}(E). \quad (6.1)$$

As always, E is monotone or exact, and in the monotone setting we always mean $\mathcal{W}_\lambda(M)$ for a fixed m_0 -value λ . We now define the analogue of 5.4 for $\mathcal{W}(E)$,

$$\text{OC} : \text{CC}_*(\mathcal{A}, \mathcal{M}) \rightarrow \text{SC}^*(E).$$

We first construct

$$\begin{aligned} \text{OC}^{n, \mathbf{p}, \mathbf{w}} : \text{CF}^*(L_n, L_0; w_{n+1}H)[\mathbf{q}] \otimes \text{CF}^*(L_{n-1}, L_n; w_n H)[\mathbf{q}] \otimes \cdots \otimes \text{CF}^*(L_0, L_1; w_1 H)[\mathbf{q}] \\ \rightarrow \text{CF}^*(w_0 H)[\mathbf{q}]. \end{aligned}$$

Observe Figure 3. Consider discs S as in 3.11 for $d+1 = n+2$ except now, mimicking 5.4:

- (1) We fix the parametrization of the domain by taking $S = \mathbb{D} \setminus \{\text{punctures}\}$, $\underline{z} = z_{n+1} = +1$ and $z_0 = 0$ (compare 5.4 and 3.12).
- (2) The last positive boundary puncture $z_{n+1} = \underline{z}$ is distinguished (it receives the module input from $\text{CC}_*(\mathcal{A}, \mathcal{M})$). We denote it by a dash in Figure 3.
- (3) The interior puncture z_0 is *negative* (it receives the output from $\text{CF}^*(w_0 H)[\mathbf{q}] \subset \text{SC}^*(E)$), meaning we make a choice $\epsilon_0 : (-\infty, 0] \times S^1 \rightarrow S$ of a cylindrical end parametrization, with $s \rightarrow -\infty$ corresponding to z_0 . We fix $\epsilon_0(s, t) = \frac{1}{2}e^{2\pi(s+it)}$ to be the parametrization by polar coordinates of $\{z \in \mathbb{C} : 0 < |z| \leq 1/2\} \subset \mathbb{D} \setminus \{\text{punctures}\}$. This fixes the S^1 parametrization at the asymptote (generators of $\text{CF}^*(w_0 H)$ are *parametrized* 1-orbits).
- (4) We no longer interpret preferred points ϕ_f in terms of maps (compare 3.11). The map formalism is not useful when geodesics connect to an interior puncture because the location of a boundary puncture together with an interior puncture fix the disc's reparametrization freedom, leaving no freedom to move the preferred point.
- (5) The preferred point ϕ_f can move along the geodesic connecting z_0 to z_{p_f} . By the Rotational symmetry trick in 4.6 it does not matter at which angle the ϕ_f come into z_0 (the β_f are S^1 -invariant near 0). When compactifying the moduli spaces, a cylinder can break off at 0 carrying preferred points ϕ_f . By 4.6, the location of each ϕ_f on the cylinder is only defined up to rotation, but once we glue the cylinder back onto the main component we fix these rotational-freedoms by placing ϕ_f on the geodesic connecting z_0, z_{p_f} .
- (6) Just as in Remark 5.2, if one uses the auxiliary data pulled back from the data for $\mu_{\mathcal{A}}$, one must deform this data when all punctures converge to the module input. Consistent universal auxiliary data is constructed as in 3.13, inductively over strata of the compactification of the above moduli space of decorated domains. On disc bubbles which do not carry the interior puncture z_0 , we require that the auxiliary data agrees with the data defining $\mu_{\mathcal{A}}$ and $\mu_{\mathcal{M}}$. On sphere bubbles at z_0 we require that the data agrees with that defining $\mu_{\text{SC}^*(E)}^1$. (These requirements will ensure that OC is a chain map.)

Analogously to the $\mathcal{R}^{d+1, \mathbf{p}}(\mathbf{x})$ defined in 3.14, except for the aforementioned modifications, we obtain a moduli space of maps solving $(du - X \otimes \gamma)^{0,1} = 0$ which we denote $\mathcal{OC}^{n, \mathbf{p}}(\mathbf{x})$. In particular, note that $(du - X \otimes \gamma)^{0,1} = 0$ turns into $\partial_s u + J(\partial_t u - w_0 X) = 0$ via ϵ_0 near z_0 .

To define $\text{OC}^{n, \mathbf{p}, \mathbf{w}}$ we mimic the definition of the $\mu^{n, \mathbf{p}, \mathbf{w}}$ in 3.15. For the same symmetry reasons [8, Lemma 3.7], there is a cancellation of the counts of solutions when more than one preferred point lies on a geodesic, so we only need to keep track of whether there is or there isn't a preferred point on each geodesic (encoded by $\mathbf{q}^1, \mathbf{q}^0$). First we determine the \mathbf{q}^0 coefficient of the image. An isolated $u \in \mathcal{OC}^{n, \mathbf{p}}(x_0, x_1, \dots, x_n, \underline{x})$ will contribute $\pm t^{E_{\text{top}}(u)} x_0$ to $\text{OC}^{n, \mathbf{p}, \mathbf{w}}(\mathbf{q}^{i_{n+1}} \underline{x} \otimes \mathbf{q}^{i_n} x_n \otimes \cdots \otimes \mathbf{q}^{i_1} x_1)$ where $i_k = 1$ for $k \in \mathbf{p}(F)$, $i_k = 0$ for $k \notin \mathbf{p}(F)$, and it will contribute 0 otherwise (for signs mimic [8, Sec.3.8], see below). The \mathbf{q}^1 coefficient of the image is determined by the requirement that the map $\mu^{d, \mathbf{p}, \mathbf{w}}$ respects the $\partial_{\mathbf{q}}$ operator.

Finally, sum up the $\text{OC}^{n,\mathbf{p},\mathbf{w}}$ as \mathbf{p}, \mathbf{w} vary to obtain:

$$\text{OC}^n : \mathcal{M}(L_0, L_n) \otimes \mathcal{A}(L_n, \dots, L_0) \rightarrow SC^*(E)$$

Example. In the first picture of Figure 3, after relabeling ϕ_1, ϕ_2 to ϕ_3, ϕ_2 (as we only sum over inclusions $\mathbf{p} : F \subset \{1, \dots, n+1\}$), an isolated solution u with asymptotic conditions x_j at z_j will contribute $\pm t^{E(u)} x_0 \in CF^*(w_0 H) \subset SC^*(E)$ to $\text{OC}^3(\underline{x} \otimes \mathbf{q}x_3 \otimes \mathbf{q}x_2 \otimes x_1)$.

Remark. The sum is weighted by signs [2, (5.24)] like in 3.15: $\text{OC}^{n,\mathbf{p},\mathbf{w}}(\mathbf{q}^{i_{n+1}} \underline{x} \otimes \mathbf{q}^{i_n} x_n \otimes \dots \otimes \mathbf{q}^{i_1} x_1)$ is weighted $(-1)^{\dagger\dagger}$ where

$$\dagger\dagger = (n+1)(\deg(\underline{x}) + 1) + \sum_{j=1}^n j(\|x_j\| + 1) + \sum_{j \in F} (\sigma(x)_{j+1}^n + \deg(\underline{x}))$$

(the second sum is the Koszul sign expected from viewing \mathbf{q} as an operator of degree -1 acting from the left, the first sum is as in [2, (5.24)], and $\deg(\underline{x}) = \|\underline{x}\| + 1$ is unreduced).

Lemma 6.1. OC is a chain map, so it induces $\text{OC} : \text{HH}_*(\mathcal{W}(E)) \rightarrow SH^*(E)$.

Sketch Proof. The proof that OC is a chain map is similar to the proof of the A_∞ -relations for $\mu_{\mathcal{A}}$ (see Remarks 3 and 4 in 3.15). We illustrate the argument in a concrete example. Consider picture 1 in Figure 3. If ϕ_1 moves towards z_3 , the limit (picture 3) contributes to $\text{OC}^3(\underline{x} \otimes (\mathcal{K}\partial_{\mathbf{q}})(\mathbf{q}x_3) \otimes \mathbf{q}x_2 \otimes x_1)$ which is a part of $\text{OC}^3 \circ (\text{id} \otimes \mu_{\mathcal{A}}^1 \otimes \text{id}^{\otimes 2})$. If ϕ_1 moves towards z_0 , in the limit (picture 2) a 2-punctured sphere bubbles off with ϕ_1 lying on the geodesic joining the two punctures. This breaking contributes to $\mu_{SC^*(E)}^1(\text{OC}^3(\underline{x} \otimes x_3 \otimes \mathbf{q}x_2 \otimes x_1))$. Now consider the 1-dimensional components of the compactified moduli space $\mathcal{OC}^{n,\mathbf{p}}(\mathbf{x})$. The oriented sum of the possible boundary degenerations sums up to zero; this corresponds to the equation $\text{OC} \circ b = \mu_{SC^*(E)}^1 \circ \text{OC}$ where b is the bar differential. \square

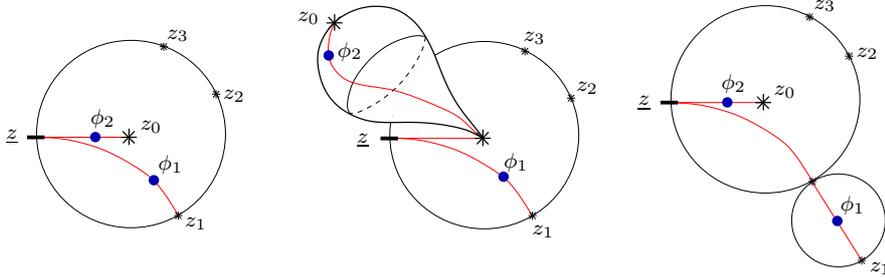


FIGURE 4. Left: $n = 3$, $F = \{1, 2\}$, $p_1 = 1$, $p_2 = 0$. Right: two broken configurations in which a preferred point converged to a puncture.

6.2. The closed-open string map. We now define the analogue of 5.5 for $\mathcal{W}(E)$,

$$\text{CO} : SC^*(E) \rightarrow \text{CC}^*(\mathcal{A}, \mathcal{M}). \quad (6.2)$$

We first construct

$$\text{CO}^{n,\mathbf{p},\mathbf{w}} : CF^*(w_0 H)[\mathbf{q}] \longrightarrow \text{Hom}(CF^*(L_{n-1}, L_n; w_n H)[\mathbf{q}] \otimes \dots \otimes CF^*(L_0, L_1; w_1 H)[\mathbf{q}], CF^*(L_n, L_0; w_{n+1} H)[\mathbf{q}])$$

This involves discs S as in 3.11 for $d+1 = n+2$ except now mimicking 5.5:

- (1) We fix the parametrization of the domain by taking $S = \mathbb{D} \setminus \{\text{punctures}\}$, $\underline{z} = z_{n+1} = -1$, $z_0 = 0$ (compare 5.5 and 3.12).
- (2) $\underline{z} = -1$ is a *negative* boundary puncture, and it is distinguished (it receives the module output from $CF^*(L_n, L_0; w_{n+1} H)[\mathbf{q}]$). We denote it by a dash in Figure 4. The other boundary punctures, z_1, \dots, z_n , are positive.

- (3) The interior puncture z_0 is *positive* (it receives the 1-orbit input from $CF^*(w_0H)[\mathbf{q}] \subset SC^*(E)$), so we make a choice $\epsilon_0 : [0, \infty) \times S^1 \rightarrow S$ of a cylindrical end parametrization, with $s \rightarrow +\infty$ corresponding to z_0 . We fix $\epsilon_0(s, t) = \frac{1}{2}e^{2\pi(-s+it)}$ to be the parametrization by polar coordinates of $\{z \in \mathbb{C} : 0 < |z| \leq 1/2\} \subset \mathbb{D} \setminus \{\text{punctures}\}$. This fixes the S^1 parametrization at the asymptote (generators of $CF^*(w_0H)$ are *parametrized* 1-orbits). To z_0 we associate a closed form α_0 , and the subclosed forms β_f satisfying $p_f = 0$.
- (4) By the Rotational symmetry trick (4.6) it does not matter that the preferred points on the geodesic $(-1, 0)$ connecting \underline{z}, z_0 do not line up with the $t = 0$ direction (see 6.1).
- (5) Just as in Remark 5.2, if one uses the auxiliary data pulled back from the data for $\mu_{\mathcal{A}}$, one must deform this data when all punctures converge to the module output. Analogously to 6.1, one constructs consistent universal auxiliary data following 3.13, by extending the auxiliary data prescribed on bubbles counted by $\mu_{\mathcal{A}}, \mu_{\mathcal{M}}, \mu_{SC^*(E)}^1$.
- (6) For $p_f \neq 0$, so preferred points on the geodesics connecting \underline{z} to z_{p_f} , the corresponding map (compare 5.4) is $\phi_f : S \rightarrow \mathbb{D} \setminus \{-1, 1\}$, $\phi_f(\underline{z}) = -1$, $\phi_f(z_{p_f}) = 1$, and we identify the map with the associated preferred point $\phi_f(0)$. The $\phi_f \in \mathbb{D}$ belong to the open unit disc with the puncture $z_0 = 0$ filled in. That a collision between ϕ_f and z_0 is not a problem is implicitly concealed in the construction of the universal data, as β_f is required to vanish at z_0 (recall $\gamma = w_0 dt$ near z_0), so the support of $d\beta_f$ will not touch z_0 .

The analogous discussion as in 6.1 holds for CO. Namely it involves taking a signed sum over the choices of \mathbf{p}, \mathbf{w} , then extending linearly (since $SC^*(E) = \oplus CF^*(w_0H)[\mathbf{q}]$), and then taking the product of all Hom spaces as we let n and the L_j vary. This yields (6.2), and thus

$$\text{CO} : SH^*(E) \rightarrow \text{HH}^*(\mathcal{W}(E)). \quad (6.3)$$

Remark. *As in 6.1, the sum over \mathbf{p}, \mathbf{w} involves signs: $\text{CO}^{n, \mathbf{p}, \mathbf{w}}(\mathbf{q}^{i_0} x_0)(\mathbf{q}^{i_n} x_n \otimes \cdots \otimes \mathbf{q}^{i_1} x_1)$ is weighted $(-1)^\star$ where $\star = \sum_{j=1}^n j(\|x_j\| + 1) + \sum_{j \in F} \sigma(x)_{j+1}^n$.*

The 0-part of (6.2), for $L \in \text{Ob}(\mathcal{W}(E))$ is $\text{CO}^0 : SC^*(E) \rightarrow \text{Hom}(\mathbb{K}, CW^*(L, L))$. The CC^* -differential δ on those homomorphisms is just the differential $\mu_{\mathcal{M}}^{0|0} = -\mu_{\mathcal{A}}^1$ acting on $CW^*(L, L)$. Thus, evaluating at $1 \in \mathbb{K}$, the map $\text{CO}^0(\cdot)(1)$ defines

$$\text{CO}^0 : SH^*(E) \rightarrow HW^*(L, L). \quad (6.4)$$

Remark 6.2. *One can also define a map*

$$\text{CO} : QH^*(E) \rightarrow \text{HH}^*(\mathcal{A}) \quad (6.5)$$

by composing $c^* : QH^*(E) \rightarrow SH^*(E)$ with $\text{CO} : SH^*(E) \rightarrow \text{HH}^*(\mathcal{A})$. In particular, the 0-part $\text{CO}^0 = \text{CO}^0(\cdot)(1)$ of (6.5) yields

$$\text{CO}^0 : QH^*(E) \rightarrow HW^*(L, L). \quad (6.6)$$

Arguing as in the proof of Lemma 5.3, (6.5) can be defined at the chain level as before, except:

- (1) we replace the interior puncture at $z_0 = 0$ by an interior marked point, with lf-cycle intersection condition given by the $QC^*(E)$ input (compare 5.5),
- (2) we place a fixed preferred point ϕ_{fixed} on the geodesic $(-1, 0)$, say at $\phi_{\text{fixed}} = -1/2$ (the choice of location of ϕ_{fixed} and the choice of associated subclosed form β_{fixed} , will not matter up to chain homotopy, by a standard homotopy argument).

The second condition is necessary, otherwise (6.6) would have output weight zero, which is not allowed. The role of ϕ_{fixed} is essentially analogous to the Floer datum for the compact category $\mathcal{F}(E)$, which is necessary when defining $\text{CO}^0 : QH^*(E) \rightarrow HF^*(L, L)$ for $L \in \text{Ob}(\mathcal{F}(E))$.

7. THE ACCELERATION FUNCTOR

7.1. The A_∞ -categories $\mathcal{F}_\diamond(\mathbf{E}), \mathcal{W}_\diamond(\mathbf{E})$ and the cochain complex $\mathbf{SC}_\diamond^*(\mathbf{E})$. We will now construct an A_∞ -functor $\mathcal{AF} : \mathcal{F}(E) \rightarrow \mathcal{W}(E)$ from the compact category to the wrapped category. On objects, $\mathcal{AF}(L) = L$ is the inclusion of the compact Lagrangians $L \in \text{Ob}(\mathcal{F}(E)) \subset \text{Ob}(\mathcal{W}(E))$. On morphism spaces there is no obvious inclusion: we cannot in general allow weight $w = 0$ in Definition 3.4, because the Floer datum approach of 5.2 does not work for non-compact Lagrangians. The key idea will be to allow weight $w = 0$ only when L_0, L_1 are compact Lagrangians, which is all we need to define \mathcal{AF} .

The A_∞ -category $\mathcal{W}_\diamond(E)$ has the same objects as $\text{Ob}(\mathcal{W}(E))$ and the same morphism spaces $CW^*(L_0, L_1)$ as $\mathcal{W}(E)$ except when L_0, L_1 are both compact, in which case

$$\begin{aligned} \text{hom}_{\mathcal{W}_\diamond(E)}(L_0, L_1) &= CF^*(L_0, L_1)[\mathbf{q}] \oplus CW^*(L_0, L_1) \\ &= CF^*(L_0, L_1)[\mathbf{q}] \oplus \bigoplus_{w=1}^{\infty} CF^*(L_0, L_1; wH)[\mathbf{q}]. \end{aligned}$$

The A_∞ -operations are defined by the same disc counts as in the construction of $\mathcal{W}(E)$, with the novelty that if one of the boundary punctures carries an input from $CF^*(L_0, L_1)[\mathbf{q}]$ then no closed 1-form α is associated to that puncture (see 3.11), so $(du - X \otimes \gamma)^{0,1} = 0$ turns into the J -holomorphic equation $du^{0,1} = 0$ near that puncture when L_0, L_1 are transverse (since the solution has finite energy it extends continuously over the puncture, so that puncture can be replaced by a marked point receiving as input an intersection point from $L_0 \cap L_1$). When L_0, L_1 are not transverse then the equation becomes $(du - X_{H_{L_0, L_1}} \otimes c(s)dt)^{0,1} = 0$ on the strip-like end of that puncture, using the Floer datum machinery of 5.2.

To extend the differential μ^1 of Definition 3.4 to $CF^*(L_0, L_1)[\mathbf{q}]$, we need to define

$$\mathfrak{d} : CF^*(L_0, L_1) \rightarrow CF^{*+1}(L_0, L_1) \quad \text{and} \quad \mathfrak{R} : CF^*(L_0, L_1) \rightarrow CF^*(L_0, L_1; H).$$

The first (following 3.6) is precisely the Floer differential as defined for the compact Fukaya category (in particular, using a Floer datum when L_0, L_1 are not transverse, see 5.2). The second (following 3.9) counts finite-energy Floer continuation solutions on the strip $\mathbb{R} \times [0, 1]$, where the homotopy H_s interpolates $H_s = H$ for $s \ll 0$ with $H_s = 0$ for $s \gg 0$ (this becomes $H_s = H_{L_0, L_1}$ for $s \gg 0$ using the Floer datum, when L_0, L_1 are not transverse). In the transverse setup, the solutions for the \mathfrak{R} map extend continuously over the puncture at $s = +\infty$. We remark that on cohomology, the map

$$\mathfrak{R} : HF^*(L_0, L_1) \rightarrow HF^*(L_0, L_1; H)$$

is the analogue of the PSS map for open strings (compare 4.4).

Let $\mathcal{F}_\diamond(E)$ be the full subcategory of $\mathcal{W}_\diamond(E)$ whose objects are the compact Lagrangians. So $\mathcal{F}_\diamond(E)$ enlarges the morphism spaces of $\mathcal{F}(E)$ from $CF^*(L_0, L_1)$ to $\text{hom}_{\mathcal{W}_\diamond}(L_0, L_1)$.

Similarly, enlarge $SC^*(E)$ from 4.2 to

$$\begin{aligned} SC_\diamond^*(E) &= QC^*(E)[\mathbf{q}] \oplus SC^*(E) \\ &= QC^*(E)[\mathbf{q}] \oplus \bigoplus_{w=1}^{\infty} CF^*(wH)[\mathbf{q}]. \end{aligned}$$

We can extend the product μ^2 , by allowing an insertion from $QC^*(E)$ at a puncture: this means that no α form is associated to that puncture, so the Floer equation becomes holomorphic near that puncture, and as we only count finite-energy solutions this implies that the solution extends holomorphically over that puncture. In particular, if both of the two positive punctures defining μ^2 involve $QC^*(E)$ -inputs, and there are no preferred geodesics, then $\gamma = 0$ so μ^2 counts holomorphic spheres and coincides with the quantum product of the $QC^*(E)$ -inputs.

To extend $\mu^1 = \mu_{SC^*(E)}^1$ defined in (4.1) to $QC^*(E)[\mathbf{q}]$, we again need to define

$$\mathfrak{d} : QC^*(E) \rightarrow QC^{*+1}(E) \quad \text{and} \quad \mathfrak{R} : QC^*(E) \rightarrow CF^*(H).$$

Let \mathfrak{d} be the differential for the chain complex $QC^*(E)$: in our convention, we pick locally finite cycles representing a basis for $H_*^{lf}(E) \cong H^{2n-*}(E)$, so $\mathfrak{d} = 0$. Let \mathfrak{R} be the Floer continuation map for the cylinder $\mathbb{R} \times S^1$, with $H_s = H$ for $s \ll 0$ and $H_s = 0$ for $s \gg 0$, with the lf-cycle intersection condition at $s = +\infty$ (finite-energy implies the solution holomorphically extends over the puncture at $s = +\infty$). On cohomology, \mathfrak{R} is the PSS map from 4.4,

$$\mathfrak{R} : QH^*(E) \rightarrow HF^*(H).$$

The natural inclusion on subcomplexes $SC^*(E) \rightarrow SC_{\diamond}^*(E)$ yields an isomorphism

$$SH^*(E) \rightarrow SH_{\diamond}^*(E) \tag{7.1}$$

on cohomology. This is because the telescope construction corresponds to a direct limit on cohomology, so we can compare cocycles for large slopes (indeed slope $w \geq 1$ suffices), but in that case the two complexes coincide. There is a natural map

$$c^* : QH^*(E) \rightarrow SH_{\diamond}^*(E) \tag{7.2}$$

induced by the inclusion of the subcomplex $QC^*(E) \rightarrow SC_{\diamond}^*(E)$, and our previous comments about the μ^2 map imply that c^* is an algebra homomorphism. On cohomology c^* is unital since $c^*[E]$ is the unit for $SH_{\diamond}^*(E)$ (by the same proof as in 4.7). Also, (7.2) factorizes as $c^* : QH^*(E) \rightarrow SH^*(E) \rightarrow SH_{\diamond}^*(E)$ through the usual c^* map and the isomorphism (7.1) (at the chain level, (7.2) and that factorization are homotopic via the homotopy $x \mapsto x \mathbf{q}$).

Finally, one can extend the string maps compatibly with the above enlargements,

$$\text{OC} : \text{HH}_*(\mathcal{W}_{\diamond}(E)) \rightarrow SH_{\diamond}^*(E) \quad \text{and} \quad \text{CO} : QH^*(E) \rightarrow SH_{\diamond}^*(E) \rightarrow \text{HH}^*(\mathcal{W}_{\diamond}(E)),$$

by allowing zero-weight insertions for pairs of compact Lagrangians L_i, L_{i+1} , using the Floer datum 5.2 when L_i, L_{i+1} are not transverse. To define $\text{CO} : SH_{\diamond}^*(E) \rightarrow \text{HH}^*(\mathcal{W}_{\diamond}(E))$ we use Remark 6.2: we always place a fixed preferred point ϕ_{fixed} at $-1/2$, whenever $z_0 = 0$ receives an lf-cycle intersection condition from $QC^*(E)$. We still allow a (free) preferred point to lie on $(-1, 0)$; this does not cause symmetry cancellations because ϕ_{fixed} is fixed.

7.2. The acceleration functor \mathcal{AF} . Define the A_{∞} -functor

$$f : \mathcal{W}(E) \rightarrow \mathcal{W}_{\diamond}(E)$$

to be the identity on objects, the inclusion on morphisms, and zero on higher order tensors.

Lemma 7.1. *f is a quasi-isomorphism (analogously $\mathcal{F}(E) \rightarrow \mathcal{F}_{\diamond}(E)$ is a quasi-isomorphism).*

Proof. f is an isomorphism at cohomology because the wrapped cohomology is a direct limit over the weights $w \rightarrow \infty$ (see 3.10), so it does not matter if the weights start at $w = 0$ or $w = 1$. The analogous claim for the inclusion $\mathcal{F}(E) \rightarrow \mathcal{F}_{\diamond}(E)$ follows by Lemma 5.1. \square

Since A_{∞} -functors which are quasi-isomorphisms can be inverted up to homotopy [41, Sec.(1i)], we can pick an inverse $f^{-1} : \mathcal{W}_{\diamond}(E) \rightarrow \mathcal{W}(E)$. Composing the natural inclusion $\mathcal{F}(E) \rightarrow \mathcal{F}_{\diamond}(E) \rightarrow \mathcal{W}_{\diamond}(E)$ (inclusion on objects and morphisms, zero on higher order tensors) with f^{-1} yields an A_{∞} -functor, which we call the *acceleration functor*,

$$\mathcal{AF} : \mathcal{F}(E) \rightarrow \mathcal{F}_{\diamond}(E) \rightarrow \mathcal{W}_{\diamond}(E) \rightarrow \mathcal{W}(E).$$

By Lemma 2.10, this induces chain maps on Hochschild homology, so:

$$\text{HH}_*(\mathcal{AF}) : \text{HH}_*(\mathcal{F}(E)) \rightarrow \text{HH}_*(\mathcal{F}_{\diamond}(E)) \rightarrow \text{HH}_*(\mathcal{W}_{\diamond}(E)) \rightarrow \text{HH}_*(\mathcal{W}(E))$$

where the first and last maps are in fact isomorphisms by Theorem 2.12 and Lemma 7.1.

Hochschild cohomology is not in general covariantly functorial (see 2.18). Nevertheless, $\mathcal{F}_\diamond(E) \rightarrow \mathcal{W}_\diamond(E)$ is the inclusion of a full subcategory, so by Lemma 2.11 it induces a restriction map $\mathrm{HH}^*(\mathcal{W}_\diamond(E)) \rightarrow \mathrm{HH}^*(\mathcal{F}_\diamond(E))$. By Theorem 2.12 and Lemma 7.1, we obtain:

$$\mathrm{HH}^*(\mathcal{A}\mathcal{F}) : \mathrm{HH}^*(\mathcal{W}(E)) \cong \mathrm{HH}^*(\mathcal{W}_\diamond(E)) \rightarrow \mathrm{HH}^*(\mathcal{F}_\diamond(E)) \cong \mathrm{HH}^*(\mathcal{F}(E)).$$

7.3. The acceleration diagram.

Theorem 7.2. *There are commutative diagrams of linear maps*

$$\begin{array}{ccc} \mathrm{HH}_*(\mathcal{F}(E)) & \xrightarrow{\mathrm{HH}_*(\mathcal{A}\mathcal{F})} & \mathrm{HH}_*(\mathcal{W}(E)) \\ \mathrm{OC} \downarrow & & \downarrow \mathrm{OC} \\ \mathrm{QH}^*(E) & \xrightarrow{c^*} & \mathrm{SH}^*(E) \end{array} \qquad \begin{array}{ccc} \mathrm{HH}^*(\mathcal{F}_\diamond(E)) & \xleftarrow{\mathrm{HH}^*(\mathcal{A}\mathcal{F})} & \mathrm{HH}^*(\mathcal{W}_\diamond(E)) \\ \mathrm{CO} \uparrow & \nearrow \mathrm{CO} & \uparrow \mathrm{CO} \\ \mathrm{QH}^*(E) & \xrightarrow{c^*} & \mathrm{SH}_\diamond^*(E) \end{array}$$

In 8.13 we prove that all maps are $\mathrm{QH}^*(E)$ -module homomorphisms, and that in the diagram on the right all maps are unital algebra homomorphisms.

Proof. For the first diagram, we prove that the following diagram commutes,

$$\begin{array}{ccccc} \mathrm{HH}_*(\mathcal{F}(E)) & \xrightarrow{\mathrm{HH}_*(\mathrm{incl})} & \mathrm{HH}_*(\mathcal{W}_\diamond(E)) & \xleftarrow{\mathrm{HH}_*(f)} & \mathrm{HH}_*(\mathcal{W}(E)) \\ \mathrm{OC} \downarrow & & \mathrm{OC} \downarrow & & \mathrm{OC} \downarrow \\ \mathrm{QH}^*(E) & \xrightarrow{c^*} & \mathrm{SH}_\diamond^*(E) & \xleftarrow{\mathrm{incl}} & \mathrm{SH}^*(E) \end{array}$$

The square on the right tautologically commutes since f is the inclusion. The left square commutes at the chain level because the extension of the OC map to $\mathcal{W}_\diamond(E)$ has been precisely defined so as to coincide with the definition of OC for the compact category. Indeed $\mathrm{CC}_*(\mathcal{F}(E)) \subset \mathrm{CC}_*(\mathcal{W}_\diamond(E))$ is a subcomplex, and since no \mathfrak{q} -terms appear in this subcomplex, the domains involved in the OC map do not carry preferred points. Therefore there are no β_f forms. By construction, when inputs come from $\mathrm{CF}^*(L_i, L_{i+1})$ there are also no α_i forms. So the total form used is $\gamma = 0$ (or a perturbation datum, when non-transverse L_i, L_{i+1} arise, in which case $X \otimes \gamma$ vanishes outside of a compact subset of E). Thus the output weight is $w_0 = 0$, so the solutions land in the subcomplex $\mathrm{QC}^*(E) \subset \mathrm{SC}_\diamond^*(E)$.

The second diagram commutes by construction. Note that $\mathrm{HH}^*(\mathcal{A}\mathcal{F}) \circ \mathrm{CO} \circ c^*$ takes values in $\mathrm{CF}^*(L_0, L_n; H) \subset \mathrm{hom}_{\mathcal{F}_\diamond(E)}(L_0, L_n)$ rather than $\mathrm{CF}^*(L_0, L_n) \subset \mathrm{hom}_{\mathcal{F}(E)}(L_0, L_n)$, due to the fixed preferred point, as explained at the end of 7.1; these values agree with the values of $\mathrm{CO} : \mathrm{QC}^*(E) \rightarrow \mathrm{CC}^*(\mathcal{F}_\diamond(E))$ up to the quasi-isomorphism $x \mapsto \mathfrak{R}(x)$. \square

Remark 7.3. *The second diagram can be rephrased in terms of the ψ -string map of 8.13. The following is a commutative diagram of unital algebra homomorphisms.*

$$\begin{array}{ccc} \mathrm{End}(\mathcal{F}_\diamond(E)) & \xleftarrow{\text{restriction}} & \mathrm{End}(\mathcal{W}_\diamond(E)) \\ \psi \uparrow & \nearrow \psi & \uparrow \psi \\ \mathrm{QH}^*(E) & \xrightarrow{c^*} & \mathrm{SH}_\diamond^*(E) \end{array}$$

8. ALGEBRAIC STRUCTURES

8.1. The ψ -string map. We use the notation $\mathcal{A} = \mathcal{W}(E)$, $\mathcal{M} = \mathcal{W}(E)$ from (6.1). Our discussion applies equally well to the compact category, with obvious simplifications (8.12).

Theorem 8.1. *There is a unital algebra homomorphism*

$$SH^*(E) \rightarrow \text{End}(\mathcal{M}) \rightarrow \text{End}(\text{HH}_*(\mathcal{M})), \quad c \mapsto \text{HH}_*(\psi_c),$$

defining a module action of $SH^*(E)$ on $\text{HH}_*(\mathcal{M})$.

By Corollary 2.3, we just need to construct a linear map $SH^*(E) \rightarrow \text{End}(\mathcal{M})$, $c \mapsto \psi_c$, which is a unital algebra homomorphism. We also claim:

Theorem 8.2. *There is a commutative diagram of unital algebra homomorphisms*

$$\begin{array}{ccccccc} QH^*(E) & \xrightarrow{c^*} & SH^*(E) & \xrightarrow{\psi} & \text{End}(\mathcal{M}) & \xrightarrow{\text{HH}_*(\cdot)} & \text{End}(\text{HH}_*(\mathcal{M})) \\ & \searrow & \downarrow \text{CO} & & \uparrow \mu_{\mathcal{M}} \circ \text{left} & & \nearrow \text{cap product} \\ & & & & \text{HH}^*(\mathcal{M}) & & \end{array}$$

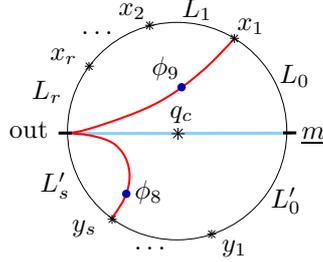
using composition product for $\text{End}(\mathcal{M})$ and $\text{End}(\text{HH}_*(\mathcal{M}))$, cup product for $\text{HH}^*(\mathcal{M})$ as in 2.14, the vertical map is the algebra homomorphism in 2.16, and see 2.17 for the cap product.

The A_∞ -category \mathcal{A} is cohomologically unital: the cohomological units are

$$e_L = \text{CO}^0(c^*[E]) \in HW^*(L, L).$$

By Corollary 2.8, it follows that $\text{CO}(c^*[E]) \in \text{HH}^*(\mathcal{A})$ is a unit for the cup product and the vertical map above is an isomorphism. In particular, the 0-part $\text{CO}^0 : SH^*(E) \rightarrow HW^*(L, L)$ (see 6.2) is a unital algebra homomorphism, using $\mu_{\mathcal{W}(E)}^2$ on $HW^*(L, L)$.

8.2. The construction of ψ . Consider the following domain S ,



which is the same domain and auxiliary data as for $\mu_{\mathcal{M}}^{**}$ (see 3.16), except we also have a positive interior puncture at $q_c = 0$ (compare 6.2). Associated to c there is a closed form α_c , and any β_f form with $p_f = q_c$ (i.e. if the preferred point ϕ_f lies on the new geodesic $(-1, 0)$).

Remark. *We fixed the $PSL(2, \mathbb{R})$ -parametrization of the disc by fixing punctures at $-1, 0, +1$. If one wanted to work with moduli of domains, one would allow the module input/output boundary punctures to move as well, and q_c would be allowed to move on the hyperbolic geodesic connecting those two boundary punctures.*

Given a 1-orbit \tilde{c} for $w_c H$, for some integer weight $w_c \geq 1$, consider

$$c = \tilde{c} \mathbf{q}^{i_c} \in CF^*(w_c H)[\mathbf{q}] \subset SC^*(E) = \bigoplus_{w=1}^{\infty} CF^*(wH)[\mathbf{q}].$$

We may abusively refer to “the asymptotic condition c at q_c ” when we mean \tilde{c} . We now define

$$\psi_c^{r|s} : \mathcal{A}(L_r, \dots, L_0) \otimes \mathcal{M}(L_0, L'_0) \otimes \mathcal{A}(L'_0, \dots, L'_s) \rightarrow \mathcal{M}(L_r, L'_s),$$

for $r, s \geq 0$. We follow the routine of 3.15: we first define the map on a vector space generator,

$$\psi_c^{r|s}(\mathbf{q}^{i_r} x_r, \dots, \mathbf{q}^{i_1} x_1, \mathbf{q}^{i_0} \underline{m}, \mathbf{q}^{j_1} y_1, \dots, \mathbf{q}^{j_s} y_s) \in CF^*(L'_s, L_r; wH)[\mathbf{q}]$$

where $\underline{m} \in CF^*(L'_0, L_0; w_0H)$, $x_i \in CF^*(L_{i-1}, L_i; w_iH)$, $y_i \in CF^*(L'_i, L'_{i-1}; w'_iH)$. We define the \mathbf{q}^0 -coefficient of the output to be the count (as usual, with weights $t^{E_{\text{top}}(u)}$ and with orientation signs) of the isolated solutions $u : S \rightarrow E$ of Floer's equation $(du - X \otimes \gamma)^{0,1} = 0$, on the above decorated domains (compare 6.2). In particular

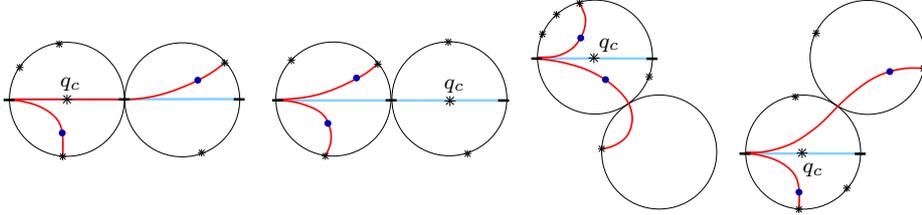
$$\gamma = w_c \alpha_c + w_0 \alpha_0 + \sum w_i \alpha_i + \sum w'_j \alpha'_j + \sum \beta_f \quad \text{and} \quad w = w_c + w_0 + \sum w_i + \sum w'_j + |F|$$

where $|F|$ is the number of preferred points. The exponents $i_r, \dots, i_1, i_0, j_1, \dots, j_s \in \{0, 1\}$ determine whether or not a preferred point arises on the geodesic connecting -1 to the relevant boundary puncture, and $i_c \in \{0, 1\}$ is associated to the geodesic $(-1, 0)$. As usual, for symmetry reasons [8, Lemma 3.7], the counts of Floer solutions which involve more than one preferred point on a geodesic cancel out. As always, we extend ψ_c so that it is $\partial_{\mathbf{q}}$ -equivariant and Λ -multilinear (compare 3.15).

Remark about signs. *As in 3.15 and 6.2, solutions are counted with a sign $(-1)^{k+\sum j \text{Deg}(z_j)}$ where: $(z_{r+s+1}, \dots, z_1) = (x_r, \dots, x_1, \underline{m}, y_1, \dots, y_s)$ in that order, $\text{Deg}(z_j) = \|z_j\| + 1$ for non-module inputs, $\text{Deg}(z_j) = \text{deg}(\underline{m}) + 1$ for the module input (when $j = s + 1$). The k is the Koszul sign obtained from viewing \mathbf{q} as an operator of degree -1 acting from the left (for example if $j_2 = 1$, then $\mathbf{q}^{j_2} y_2$ above contributes $\|x_r\| + \dots + \|x_1\| + \text{deg}(\underline{m}) + \|y_1\|$ to k).*

Remark 8.3. *Geometrically, ψ_c is similar to $\text{CO}(c)$, with the following differences. For ψ_c there is a fixed input boundary puncture at $+1$ receiving the module input, whereas for CO all input boundary punctures are free to move (CO does not receive a distinguished module input). By Remark 5.2, for CO one cannot simply pull back the auxiliary data used for $\mu_{\mathcal{A}}$, whereas for ψ we can pull back the data used for $\mu_{\mathcal{M}}$ (which is the same as for $\mu_{\mathcal{A}}$) because q_c can only bubble off at a module input/output puncture. In the bubbling analysis for CO , there can arise a disc bubble carrying q_c with just one output puncture, counted by $\text{CO}^0(c)$; for ψ_c a disc bubble carrying q_c will always have both an input and an output puncture.*

8.3. Bubbling analysis: the possible degenerations. We now consider the degenerations for a 1-family of Floer solutions on such domains.



As an illustration, in the case $R = 3$, $S = 2$, the above pictures (1)-(4) are four possible degenerations. In (1), a subset of the boundary punctures converges to $+1$. In (2) a subset of the boundary punctures converges to -1 . In (3) a subset of the boundary punctures on the lower arc connecting $-1, +1$ converges to a puncture on the lower arc different from ± 1 . In (4) the analogue of (3) happens for the upper arc. In all four cases, a subset of the preferred points may converge to the same puncture, so these will reappear on the new disc bubble.

The main component of the above broken solutions is the one carrying q_c , and it is counted by the ψ map. The new disc bubbles that appear, which do not carry q_c , are counted by $\mu_{\mathcal{M}}^{*|*}$ in pictures (1)-(2) and by $\mu_{\mathcal{A}}^*$ in pictures (3)-(4). We recall that the auxiliary data has been constructed universally and consistently as in 3.13, so that those new disc bubbles (which do not carry q_c) will be endowed with the universal data constructed for $\mu_{\mathcal{M}}$ and $\mu_{\mathcal{A}}$.

The only other remaining degeneration, occurs when a Floer cylinder breaks off at $q_c = 0$, possibly carrying a subset of the preferred points. These cylinders are counted by the differential $\mu_{SC^*(E)}^1(c)$ on $SC^*(E)$.

8.4. Proof that ψ_c is a morphism of bimodules for $[c] \in \mathbf{SH}^*(\mathbf{E})$. Extending ψ_c linearly in c , we may now replace c by a cycle in $SC^*(E)$ (i.e. a representative of a class $[c] \in SH^*(E)$). The Floer cylinders breaking off at q_c all cancel out, since $\mu_{SC^*(E)}^1(c) = 0$, so the above degenerations (1)-(4) prove that ψ_c is a morphism of \mathcal{A} -bimodules (as defined in 2.8)

$$\psi_c : \mathcal{M} \rightarrow \mathcal{M}.$$

As usual, it is the condition of $\partial_{\mathbf{q}}$ -equivariance that ensures that the algebra accurately reflects the bubbling analysis (compare the Remarks in 3.15). By Lemma 2.1, we obtain a chain map associated with the bimodule map ψ_c ,

$$CC_*(\psi_c) : CC_*(\mathcal{A}, \mathcal{M}) \rightarrow CC_*(\mathcal{A}, \mathcal{M}).$$

We may abusively write ψ_c when we mean $CC_*(\psi_c)$. Explicitly, the chain map $CC_*(\psi_c)$ on a vector space generator $\underline{m} \otimes x_n \otimes \cdots \otimes x_1 \in CC_n(\mathcal{A}, \mathcal{M})$ equals

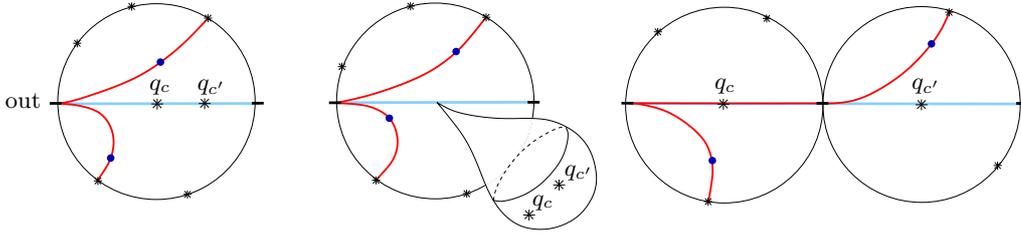
$$\sum (-1)^{\sigma_1(\deg(\underline{m}) + \sigma_{r+1})} \psi_c^{r|n-s+1}(x_r \otimes \cdots \otimes x_1 \otimes \underline{m} \otimes x_n \otimes \cdots \otimes x_s) \otimes x_{s-1} \otimes \cdots \otimes x_{r+1}$$

summing over all obvious choices of $r, s \geq 0$. Thus, on homology, we obtain:

$$HH_*(\psi_c) : HH_*(\mathcal{A}, \mathcal{M}) \rightarrow HH_*(\mathcal{A}, \mathcal{M}).$$

A standard argument shows that $HH_*(\psi_c)$ does not depend on the choice of representative c for $[c] \in SH^*(E)$ (one checks that changing c by a boundary chain corresponds, by the above degeneration analysis, to changing ψ_c by the boundary of a pre-morphism, as in 2.8).

8.5. Proof that ψ is an $\mathbf{SH}^*(\mathbf{E})$ -module action. We consider a ‘‘higher-order’’ ψ -string map: in addition to $q_c = 0$, we introduce another positive interior puncture $q_{c'} \in (0, 1)$ which is allowed to move along the geodesic joining $q_c = 0$ and $+1$.



The possible degenerations in a 1-family are:

- (1) If $q_{c'}$ converges to $q_c = 0$, then a 3-punctured sphere counted by the $\mu_{SC^*(E)}^2$ product appears as in picture 2. These breakings are counted by $\psi_{c*c'}$.
- (2) If $q_{c'}$ converges to $+1$ (possibly together with a subset of the boundary punctures and a subset of the preferred points), then a disc bubble carrying $q_{c'}$ appears which is counted by $\psi_{c'}$ as in picture 3. These degenerations are counted by $\psi_c \circ \psi_{c'}$.
- (3) If $q_{c'}$ converges to a point in the interval $(0, 1)$, then a subset of the boundary punctures (possibly together with a subset of the preferred points) may converge to $+1$, or -1 , or a boundary puncture different from ± 1 . These disc bubbles are counted respectively by the composition maps $\mu_{\mathcal{M}}, \mu_{\mathcal{M}}, \mu_{\mathcal{A}}$. The main component of the broken configuration still carries $q_c, q_{c'}$, and the count of these (now rigid) main components defines a pre-morphism $K : \mathcal{M} \rightarrow \mathcal{M}$ as in 2.8. The count of the broken configurations is therefore δK , where δ is the differential on pre-morphisms.

This proves that $SH^*(E) \rightarrow \text{End}(\mathcal{M})$ is an algebra homomorphism. By Lemma 2.2(1),

$$CC_*(\psi_{c*c'}) - CC_*(\psi_c) \circ CC_*(\psi_{c'}) = CC_*(\delta K) = CC_*(K) \circ b + b \circ CC_*(K),$$

so $CC_*(K)$ is a chain homotopy which shows $HH_*(\psi_{c*c'}) = HH_*(\psi_c) \circ HH_*(\psi_{c'})$, as required.

8.6. **The ψ -string map for $c \in \mathbf{QH}^*(E)$.** Analogously, for any locally finite cycle $c \in QH^*(E)$, we can construct an algebra homomorphism

$$\psi : QH^*(E) \rightarrow \text{End}(\mathcal{M}), \quad c \mapsto \psi_c.$$

In the construction in 8.2 we now require q_c to be an interior marked point, rather than an interior puncture, and we have no auxiliary forms associated to q_c . In 8.2, all auxiliary forms vanished near q_c except for the α form associated to q_c . So in this new setup, $\gamma = 0$ near q_c , therefore the Floer solutions $u : S \rightarrow E$ are J -holomorphic near the marked point q_c . So it remains to discuss the intersection condition at q_c .

As the solutions u are J -holomorphic near q_c , we may require that they intersect the given lf-cycle c at q_c , so $u(q_c) \in c$. More precisely, consider the moduli space \mathcal{S} of all solutions u but without imposing the condition $u(q_c) \in c$. We can ensure that the lf-cycles chosen for the definition of $QH^*(E)$ are transverse to the evaluation map $\text{ev} : \mathcal{S} \rightarrow E$ at $q_c = 0 \in \mathbb{D}$. Then $\text{ev}^{-1}(c) \subset \mathcal{S}$ is the space of all solutions u described above satisfying the condition $u(q_c) \in c$. Then ψ_c is the count of the isolated points (i.e. the zero-dimensional components) of $\text{ev}^{-1}(c) \subset \mathcal{S}$ with appropriate orientation signs and with weight $t^{E_{\text{top}}(u)}$ (as in 3.14).

Following the same arguments as in 8.4, 8.3 and 8.5, we obtain a $QH^*(E)$ -module structure on $\text{HH}_*(\mathcal{A}, \mathcal{M})$ (in particular in 8.5 the 3-punctured sphere bubbles which carry lf-cycle intersection conditions at $q_c, q_{c'}$ are precisely those which count the quantum product $c * c'$).

Lemma 8.4. *There is a commutative diagram*

$$\begin{array}{ccc} SH^*(E) & \xrightarrow{\psi} & \text{End}(\mathcal{M}) \longrightarrow \text{End}(\text{HH}_*(\mathcal{M})) \\ c^* \uparrow & \nearrow \psi & \\ QH^*(E) & & \end{array}$$

Proof. This can be proved just like Lemma 5.3. A cleaner proof, is to first define

$$\psi : SH_{\diamond}^*(E) \rightarrow \text{End}(\mathcal{W}_{\diamond}(E))$$

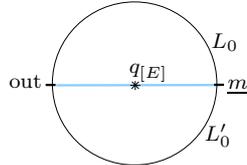
using notation from 7.1, by combining the two constructions from 8.5 and 8.6 (see 8.13). Now observe that by definition this is compatible with the inclusion $c^* : QC^*(E) \rightarrow SC_{\diamond}^*(E)$. \square

8.7. Proof that ψ is unital.

Lemma 8.5. *For $\psi : QH^*(E) \rightarrow \text{End}(\mathcal{M})$,*

$$\psi_{[E]} = (-1)^{1+\text{deg}(\cdot)} 1_{\mathcal{M}} \in \text{End}(\mathcal{M}).$$

Proof. $\psi_{[E]} : \mathcal{M}(L_0, L'_0) \rightarrow \mathcal{M}(L_0, L'_0)$ is a count of J -holomorphic discs as in the picture,



involving a tautological intersection condition $[E]$ at $q_{[E]} = 0$. We claim that $\psi_{[E]}$ is the identity bimodule map $\mathcal{M} \rightarrow \mathcal{M}$ from 2.6 (up to sign). In the case when there are $r + s \geq 1$ non-module inputs, if the moduli space of solutions for $\psi_{[E]}^{r|s}$ were non-empty, then it would contain a non-constant solution (as $r + s \geq 1$, there are at least three Lagrangians involved, and these do not have a triple intersection, for generic Hamiltonian perturbation data). But this non-constant solution u gives rise to a 1-family of solutions $u \circ g$ where $g \in SL(2, \mathbb{R})$

belongs to the \mathbb{R} -family of hyperbolic isometries which fix ± 1 (the module output/input). These isometries move the $r + s \geq 1$ other boundary markers and the marker $q_{[E]} = 0$, but we can simply reposition the marker $q_{[E]}$ to be at the origin of the disc since the intersection condition with $[E]$ is automatically satisfied also for $u \circ g$. These solutions are therefore not rigid, which is not allowed for dimension reasons, so the moduli space must in fact be empty.

Clarification. *The auxiliary data in the construction of $\psi_{[E]}$ is, by construction, $PSL(2, \mathbb{R})$ -equivariant so $u \circ g$ is a legitimate solution counted by $\psi_{[E]}$ (it automatically hits $[E]$ at 0). Although for $\mu_{\mathcal{M}}$ the solutions $u, u \circ g$ are considered to be the same (since we must quotient out by the residual \mathbb{R} -freedom of the $PSL(2, \mathbb{R})$ maps fixing ± 1), for $\psi_{[E]}$ these solutions are considered to be distinct because the $r + s \geq 1$ boundary markers have moved. This reparametrization trick is legitimate: transversality for solutions of $\psi_{[E]}$ is automatically implied by the transversality for solutions of $\mu_{\mathcal{M}}$.*

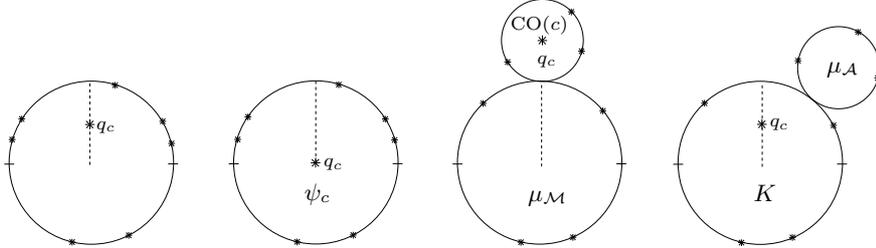
When $r + s = 0$, there are only the module output/input markers, and only two Lagrangians L_0, L_1 are involved. The previous argument still holds, so solutions are not rigid unless they are constant, except that now constant solutions are allowed: an intersection point in $L_0 \cap L_1$ when L_0, L_1 are transverse (if they are not transverse, then we identify $\mathbb{D} \setminus \{\pm 1\} \cong \mathbb{R} \times S^1$ and use the perturbation datum $H_{L_0, L_1}, J_{L_0, L_1}$ which is s -independent, see 5.2, so for any Hamiltonian chord $x \in CF^*(L_0, L_1) \equiv CF^*(L_0, L_1; H_{L_0, L_1})$ the analogue of the ‘‘constant solutions’’ are the s -independent solutions $u(s, t) = x(t)$). So $\psi_{[E]}^{0|0}$ is the identity.

Remark. *The sign $\psi_{[E]}^{0|0}(\underline{m}) = (-1)^{\deg(\underline{m})+1} \underline{m}$ is caused by the Remark about signs in 8.2, consistently with the sign for cohomological units $\mu_{\mathcal{M}}^2(e_X, \underline{m}) = (-1)^{\deg(\underline{m})+1} \underline{m}$ (see 2.10). \square*

Corollary 8.6. $\psi : SH^*(E) \rightarrow \text{End}(\mathcal{M}) \rightarrow \text{End}(\text{HH}_*(\mathcal{M}))$ is a unital algebra homomorphism.

Proof. We already proved that it is an algebra homomorphism in 8.5. Unitality follows from Lemmas 8.4 and 8.5 because c^* is a unital ring homomorphism. \square

8.8. Proof of Theorem 8.2. We already proved in Lemma 2.6 that the right half of the diagram commutes, we now prove the left half commutes.



Consider the domains from 8.2, except now we allow the positive interior puncture q_c to move along the imaginary line segment $[0, 1)i$, rather than fixing it at 0, as in the first picture (in the picture we have omitted preferred points for simplicity). Consider the possible degenerations of a 1-family of solutions for such data. The two interesting degenerations occur if $q_c \in [0, 1)i$ converges to one of the end-points: if $q_c \rightarrow 0$ we get the second picture above which is just the ψ_c map; if $q_c \rightarrow i$ then a bubble appears which carries q_c and these bubbles are counted by $\text{CO}(c)$, and the total broken configuration (in the third picture) is counted by $\mu_{\mathcal{M}} \circ_{\text{left}} \text{CO}(c)$.

Technical Remark. *We used the residual \mathbb{R} -symmetry for the domains used to define $\mu_{\mathcal{M}}$, to fix at $+i$ the input coming from $\text{CO}(c)$.*

There other degenerations involve a $\mu_{\mathcal{A}}$ -bubble or a $\mu_{\mathcal{M}}$ -bubble, with q_c converging to a point in $[0, 1)i$. The main component is a rigid solution as in the first picture above: call $K_c : \mathcal{M} \rightarrow \mathcal{M}$ the pre-morphism of bimodules which counts such rigid solutions. Thus, these

degenerations are counted by δK_c . The total count of the boundaries of those 1-families prove that $\psi_c = \mu_{\mathcal{M} \circ_{\text{left}}} \text{CO}(c)$ up to a δ -boundary term. Thus, on homology, $\text{HH}_*(\psi) = \mu_{\mathcal{M} \circ_{\text{left}}} \text{CO}$, as required. In 8.9 and 8.10 we prove that CO is a unital algebra homomorphism.

8.9. The cohomological unit.

Lemma 8.7. *The maps in the diagram in Theorem 8.2 are unital, namely*

- (1) \mathcal{A} is cohomologically unital: $e_L = \text{CO}^0(1) \in \text{HW}^*(L, L)$ are cohomological units.
- (2) $\text{CO}(1) \in \text{HH}^*(\mathcal{A})$ is a unit for the cup product defined in 2.14.

Proof. Let $m \in \mathcal{A}(L_1, L_0) = \mathcal{M}(L_1, L_0)$ be a $\mu_{\mathcal{A}}^1$ -cycle, and consider $\text{CO}^0(1) \in \mathcal{A}(L_1, L_1)$. By the commutativity of the left half of the diagram in Theorem 8.2,

$$\begin{aligned} \mu_{\mathcal{A}}^2(\text{CO}^0(1), m) &= -\mu_{\mathcal{M}}^2(\text{CO}^0(1), \underline{m}) \\ &= -(\mu_{\mathcal{M} \circ_{\text{left}}} \text{CO}(1))(\underline{m}) \\ &= -(\psi_E(\underline{m}) + \delta K_c(\underline{m})) \\ &= -(-1)^{\deg(\underline{m})+1} m - \mu_{\mathcal{M}}^{0|0}(K_c(\underline{m})) \\ &= (-1)^{\|\underline{m}\|+1} m + \mu_{\mathcal{A}}^1(K_c(\underline{m})), \end{aligned}$$

where K_c was defined in 8.8. Thus $\mu_{\mathcal{A}}^2(e_L, m) = (-1)^{\|\underline{m}\|+1} m$ on cohomology, as required by 2.10. A similar argument as in 8.8, using the obvious right-composition $\mu_{\mathcal{M} \circ_{\text{right}}}$ making $q_c \rightarrow -i$ instead of $+i$ (we are not using compatibility with product structures in this argument), shows that $\mu_{\mathcal{A}}^2(m, e_L) = m$ up to a $\mu_{\mathcal{A}}^1$ -boundary. So \mathcal{A} is cohomologically unital. The second claim follows by combining the first claim with Corollary 2.8. \square

Remark 8.8. *In general, constructing the cohomological unit for \mathcal{A} is not straightforward, even in the exact setup (see the element Φ_H in [41, Sec.(8c) and (8k)]). For L, E exact, one can construct a relative PSS-map $c^* : H^*(L) \rightarrow \text{HF}^*(L, L)$ and show that it is a unital algebra isomorphism using TQFT arguments (see $c^* : H^*(L) \rightarrow \text{HW}^*(L, L)$ in [34]). Another approach (see the final comment of [41, Sec.(8l)]), is to use a Morse-Bott chain level model for $\text{HF}^*(L, L)$: in the monotone setup, Auroux [12] takes $\text{CF}^*(L, L) = C_{n-*}(L)$ and shows that $[L]$ is a cohomological unit. However, it is not straightforward to make the Morse-Bott model compatible with the A_{∞} -structure (indeed, we used the Floer datum approach, see 5.2).*

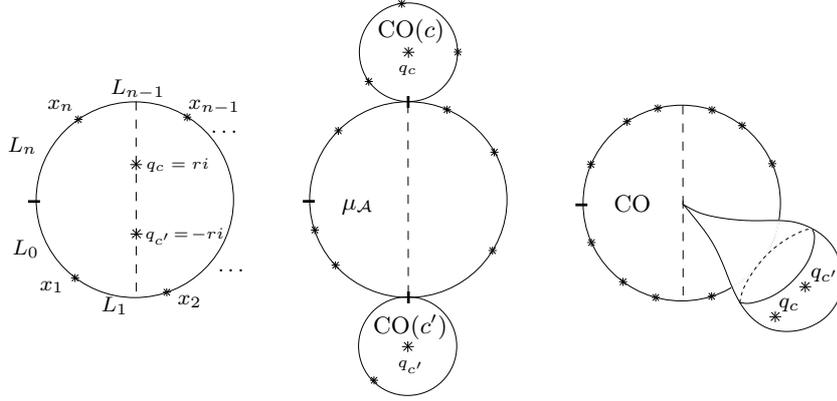
8.10. The CO map is a unital algebra homomorphism.

Corollary 8.9. *The vertical map in Theorem 8.2 is an isomorphism, and the left half of the diagram consists of unital algebra homomorphisms.*

Proof. By Lemma 8.7(1) and Corollary 2.8, the first claim follows. Since ψ is a unital algebra homomorphism, the second claim follows by the left half of the commutative diagram. \square

One can also show directly that $\text{CO} : SH^*(E) \rightarrow \text{HH}^*(\mathcal{M})$ is an algebra homomorphism by considering the degenerations of a 1-family of Floer solutions shown in the first picture below: we place two positive interior punctures $q_c = ri$, $q_{c'} = -ri$ on the imaginary axis, where $r \in (0, 1)$. Apart from the obvious bubblings of type $\mu_{\mathcal{A}}, \mu_{\mathcal{M}}$ which give δ -boundary terms, there are degenerations when r converges to 0 or 1. If $r \rightarrow 0$, the bubbling at 0 is counted by the quantum product; see picture 3. If $r \rightarrow 1$, we obtain the broken configuration in picture 2, which is counted by the cup product $\text{CO}(c) * \text{CO}(c')$; see 2.14.

Technical Remark. *We used the residual $\text{PSL}(2, \mathbb{R})/\mathbb{R}$ -symmetry for the domains that define $\mu_{\mathcal{A}}$, to fix at $\pm i$ the two inputs coming from $\text{CO}(c), \text{CO}(c')$.*



The above Corollary also shows that the 0-part $\text{CO}^0 : SH^*(E) \rightarrow HW^*(L, L)$ (see 6.2), for any $L \in \text{Ob}(\mathcal{A})$, is a unital algebra homomorphism, because the cup product on CC^* restricted to the 0-parts is the usual $\mu_{\mathcal{A}}^2$ -product on $HW^*(L, L)$. Alternatively, the same direct proof as above applies, ignoring all boundary punctures except $\underline{z} = -1$.

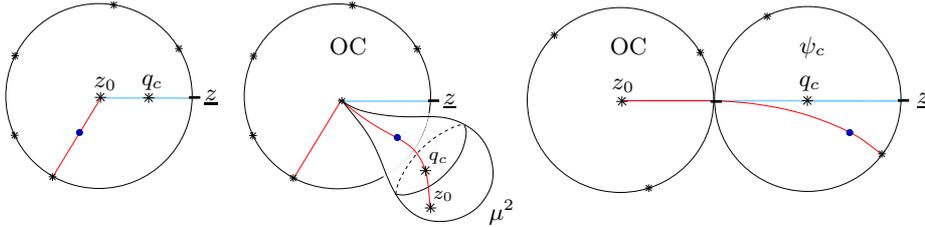
8.11. The CO and OC string maps respect the module structure. That the CO map respects the module structure is immediate from the fact that it is a ring homomorphism (the module action by c is defined by the cup product action of $\text{CO}(c)$ on $\text{HH}^*(\mathcal{M})$).

Theorem 8.10. $\text{OC} : \text{HH}_*(\mathcal{W}(E)) \rightarrow SH^*(E)$ respects the $SH^*(E)$ -module structure:

$$c \bullet \text{OC}(x) = \text{OC}(\psi_c(x))$$

where \bullet is the $\mu_{SH^*(E)}^2$ -product, $c \in SH^*(E)$, $x \in \text{HH}_*(\mathcal{W}(E))$. Similarly, $\text{OC} : \text{HH}_*(\mathcal{F}(E)) \rightarrow QH^*(E)$ and $\text{OC} : \text{HH}_*(\mathcal{W}(E)) \rightarrow SH^*(E)$ both respect the $QH^*(E)$ -module structure.

Proof. We will just prove the first claim, since the other claims are analogous (indeed simpler). We consider a “higher-order” OC-map: in addition to the negative interior puncture $z_0 = 0$, we introduce a positive interior puncture $q_c \in (0, 1)$ which is allowed to move along the geodesic joining $z_0 = 0$ and the module input puncture $+1$, as in the first picture below.



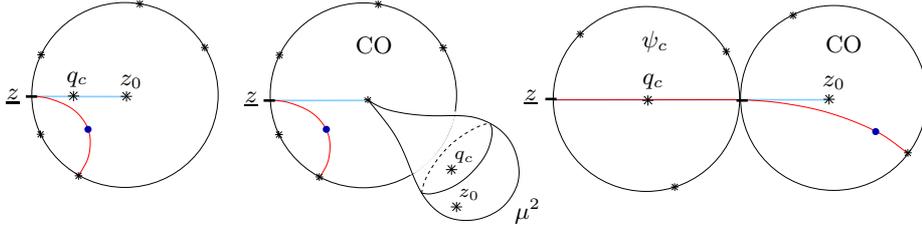
The domains are the same as those used to define OC except we introduce a new positive interior puncture $q_c \in (0, 1)$ which is allowed to move on the geodesic connecting $z_0 = 0$, $\underline{z} = +1$ (this is an interior marked point when we work with the $QH^*(E)$ -action). In particular, there may be preferred points on the geodesics connecting any puncture to $z_0 = 0$ (when working with the compact category these do not arise), and there may also be preferred points on the geodesic connecting $z_0 = 0$ to q_c (these do not arise when working with the compact category, or when working with the $QH^*(E)$ -action for the wrapped category). As usual, all $\beta_f = 0$ near the puncture q_c , and the construction of universal auxiliary data follows the routine of 3.13 (compare 4.6, 6.1, 6.2). In particular, the auxiliary data is chosen to extend

the data already chosen for $\mu_{\mathcal{A}}^*$, $\mu_{\mathcal{M}}^{*|*}$, $\mu_{SC^*(E)}^1$, $\mu_{SC^*(E)}^2$, OC, ψ_c (bubbles counted by these maps arise in the compactification of the above moduli spaces).

Let us consider the possible degenerations for a 1-family of Floer solutions for these decorated domains. The uninteresting degenerations are: disc bubbling at the module input \underline{z} counted by $\mu_{\mathcal{M}}$ or disc bubbling at another boundary puncture counted by $\mu_{\mathcal{A}}$ (these two contribute to the bar differential acting on $CC_*(\mathcal{A}, \mathcal{M})$); or a Floer cylinder breaks off at z_0 possibly carrying preferred points (these contribute to $\mu_{SC^*(E)}^1$); or a Floer cylinder breaks off at q_c (possibly carrying preferred points). In all three cases, the main component is a rigid solution whose domain is as in the first picture above. The count K_c of these rigid solutions is the chain homotopy. These degenerations are therefore counted by $K_{\mu_{SC^*(E)}^1(c)} + K_c \circ b_{CC_*(\mathcal{W}(E))} + \mu_{SC^*(E)}^1 \circ K_c$, and therefore disappear on cohomology.

The interesting degenerations occur when $q_c \rightarrow z_0 = 0$ or $q_c \rightarrow \underline{z} = +1$, in the second and third pictures above. When $q_c \rightarrow 0$, a bubble appears carrying two punctures z_0, q_c , and these contribute to $\mu_{SC^*(E)}^2$. When $q_c \rightarrow +1$, a disc bubbles off carrying q_c (which due to the residual \mathbb{R} -translation freedom can be fixed to be located at $q_c = 0$ of the new disc domain) which are counted by ψ_c , whilst the main component is counted by OC. Thus, these degenerations are counted by $OC \circ \psi_c - \mu_{SC^*(E)}^2(c, \cdot) \circ OC$, as required. \square

The analogous ‘‘higher order’’ CO-map, with an additional positive interior puncture $q_c \in (-1, 0)$, and the analogous two degenerations as in the previous picture are:



This implies that on cohomology,

$$\boxed{CO(c \bullet x) = \psi_c \circ CO(x)}$$

where \bullet is the $\mu_{SH^*(E)}^2$ -product, $c, x \in SH^*(E)$, and \circ is the composition in the sense of 2.3.

8.12. The ψ -structure for the compact Fukaya category. The construction of ψ_c for the category $\mathcal{F}(E)$, and a given lf-cycle $c \in QH^*(E)$, is a simplification of 8.2. There are no auxiliary $\alpha_k, \beta_f, \gamma$ forms, no preferred points / geodesics, so the Floer equation is replaced by the J -holomorphic equation $du^{0,1} = 0$ when each pair of Lagrangians is transverse, whereas in the non-transverse case we apply the Floer datum machinery of 5.2. By the same algebra formalism as in 8.2, ψ is determined by the counts

$$\psi_c^{r|s}(x_r, \dots, x_1, \underline{m}, y_1, \dots, y_s) \in CF^*(L'_s, L_r)$$

where $\underline{m} \in CF^*(L'_0, L_0)$, $x_i \in CF^*(L_{i-1}, L_i)$, $y_i \in CF^*(L'_i, L'_{i-1})$. Notice no Hamiltonians are present (except those coming from the Floer datum machinery, which are suppressed from the notation). These maps count J -holomorphic curves $u : S \rightarrow E$ (as usual, with weight $t^{E_{\text{top}}(u)}$ and orientation signs) on the domains S as in the picture in 8.2 except:

- the geodesic/preferred point decorations are omitted,
- $q_c = 0$ is now just an interior marked point, rather than a puncture.

Therefore, the auxiliary data in the construction of $\psi_c^{r|s}$ is the same as the universal data which defines $\mu_{\mathcal{M}}^{r|s}$ for the compact category in 5.3 with the exception that we have additionally an

interior marked point $q_c = 0$. Recall that the data for $\mu_{\mathcal{M}}^{\tau|s}$ is, by construction, $PSL(2, \mathbb{R})$ -equivariant (recall 3.12), but unlike 5.3 for the domains S above we no longer have any residual \mathbb{R} -translation freedom because we fixed a new marker at $q_c = 0$.

The discussion of the intersection condition at q_c is the same as in 8.6. In the solution counts, the weights $t^{E_{\text{top}}(u)}$ are still defined as in 3.14. In the case of transverse Lagrangians, when no Floer datum is needed, $E_{\text{top}}(u) = \int_{\mathbb{D}} u^* \omega$ is the usual energy of the J -holomorphic curve, but in general the Floer datum Hamiltonians will appear in the formula for E_{top} .

8.13. The acceleration functor respects the algebraic structures. Recall the definition of $\mathcal{W}_{\diamond}(E)$ in 7.1. Just as we extended the OC, CO maps to $\mathcal{W}_{\diamond}(E)$, we can extend the module action. This reduces to extending the definition of the bimodule endomorphism ψ_c to $\psi_c \in \text{end}(\mathcal{W}_{\diamond}(E))$ (see the proof of Lemma 8.4), so geometrically we need to explain what we do when we feed $CF^*(L_i, L_{i+1})$ inputs. As usual, this means we do not associate an α -form to the input (but there may be a preferred geodesic and preferred points, as usual, determining β_f forms). In the case when L_i, L_{i+1} are transverse, the Floer solution is holomorphic near the boundary puncture, so as usual finite-energy implies that the Floer solution extends continuously over the puncture. When L_i, L_{i+1} are not transverse, we use the Floer datum machinery from 5.2. The properties of the module structure for $\text{HH}_*(\mathcal{W}_{\diamond}(E))$ and $\text{HH}^*(\mathcal{W}_{\diamond}(E))$ now carry over just as for $\mathcal{W}(E)$. To conclude the proof of the final claim in Theorem 7.2, in light of 8.11, it remains to show the following result.

Theorem 8.11. *For the acceleration functor \mathcal{AF} from 7.2, $\text{HH}_*(\mathcal{AF})$ and $\text{HH}^*(\mathcal{AF})$ are $QH^*(E)$ -module maps, and $\text{HH}^*(\mathcal{AF})$ is a unital algebra homomorphism.*

Proof. Let $c \in QC^*(E)$. Consider what happens when $\psi_c \in \text{end}(\mathcal{W}_{\diamond}(E))$ receives all its inputs from $CF^*(L_i, L_{i+1})$ summands. By the construction described above, in this case the Floer solutions counted by $\psi_c \in \text{end}(\mathcal{W}_{\diamond}(E))$ are precisely those counted by the map $\psi_c \in \text{end}(\mathcal{F}(E))$ (in particular, the total output weight is zero, since no preferred points and no α forms are present, so ψ_c lands again in a summand of the type $CF^*(L, L')$). Therefore

$$\text{incl} \circ \psi_c = \psi_c \circ \text{incl}$$

at the level of bimodule maps, and hence also on Hochschild homology, where $\text{incl} : \mathcal{F}(E) \rightarrow \mathcal{W}_{\diamond}(E)$ is the natural inclusion. Recall the natural inclusion $f : \mathcal{W}(E) \rightarrow \mathcal{W}_{\diamond}(E)$ from 7.2. Since the action on $\mathcal{W}_{\diamond}(E)$ is defined by extending the action defined on $\mathcal{W}(E)$,

$$f \circ \psi_c = \psi_c \circ f$$

at the level of bimodule maps, and hence on Hochschild homology. Finally $\text{HH}_*(\mathcal{AF}) = \text{HH}_*(\text{incl}) \circ \text{HH}_*(f)^{-1}$ is a composite of maps which preserve the module structure, so the composite also preserves the module structure.

Similarly, the cup product action of $\text{CO}(c)$ on $\text{CC}^*(\mathcal{W}_{\diamond}(E))$ with all inputs coming from $CF^*(L_i, L_{i+1})$ agrees with $\text{CO}(c)$ acting on $\text{CC}^*(\mathcal{F}_{\diamond}(E))$ (up to a chain homotopy, due to the fixed preferred point; compare the proof of Theorem 7.2). Thus

$$\text{HH}^*(\text{incl}) \circ \text{CO}(c) = \text{CO}(c) \circ \text{HH}^*(\text{incl})$$

where $\text{CC}^*(\text{incl}) : \text{CC}^*(\mathcal{W}_{\diamond}(E)) \rightarrow \text{CC}^*(\mathcal{F}_{\diamond}(E))$ is the restriction map from Lemma 2.11. So $\text{HH}^*(\mathcal{AF})$ preserves the $QH^*(E)$ -module structure. That the inclusion $\mathcal{F}_{\diamond}(E) \rightarrow \mathcal{W}_{\diamond}(E)$ respects the cup product action on CC^* follows by homological algebra (since the $\mu_{\mathcal{A}}, \mu_{\mathcal{M}}$ maps agree on the subcategory). Thus $\text{HH}^*(\mathcal{AF})$ preserves the cup product, so it is an algebra homomorphism. The fact that it is unital follows from the second diagram in Theorem 7.2:

$$\text{HH}^*(\mathcal{AF})(\text{CO}(c^*[E])) = (\text{HH}^*(\mathcal{AF}) \circ \text{CO} \circ c^*)[E] = \text{CO}([E]) \in \text{HH}^*(\mathcal{F}(E)). \quad \square$$

9. THE STRING MAPS RESPECT EIGENSUMMANDS

9.1. Eigensummands. Let $a = c_1(TE) - \lambda 1 \in QH^*(E)$, and let $a^N = a * a * \dots * a$ be the quantum product N times. We call λ -eigensummand $QH^*(E)_\lambda$ the subalgebra of $QH^*(E)$ given by the generalized eigenspace for $c_1(TE)$:

$$QH^*(E)_\lambda = \{q \in QH^*(E) : a^N * q = 0 \text{ for some } N \in \mathbb{N}\}.$$

These are of course only non-zero when λ is an eigenvalue of $c_1(TE)* \in \text{End}(QH^*(E))$.

Similarly, since $SH^*(E)$ is a $QH^*(E)$ -module via c^* (see 4.8) we get a subalgebra $SH^*(E)_\lambda$ of $SH^*(E)$ given by the generalized λ -eigensummand for multiplication by $c^*(c_1(TE))$.

Remark. $SH^*(E)$ is typically infinite-dimensional, so it is not clear whether it fully decomposes into eigensummands or whether there is a well-defined projection $SH^*(E) \rightarrow SH^*(E)_\lambda$.

9.2. Relating m_0 and c_1 -eigenvalues. Recall $m_0(L)$ is the count of J -holomorphic discs, of Maslov index 2, with a boundary marked point constraint $[\text{pt}] \in C_*(L)$, modulo $PSL(2, \mathbb{R})$ -reparametrization. Following [12, Lemma 6.7], each disc hits $\text{PD}[c_1(TE)]$ once algebraically [12, Lemma 3.1] so we fix the parametrization of the domain so that the intersection with $\text{PD}[c_1(TE)]$ occurs at 0 and the intersection with $[\text{pt}]$ occurs at -1 . This defines $\text{CO}^0(c_1(TE))$ in the Morse-Bott model $CF_{\text{Bott}}^*(L, L) = C_{n-*}(L)$ (compare Remark 8.8). This argument implies the following result, due to Kontsevich, Seidel and Auroux [12],

$$\boxed{\text{CO}^0(c_1(TE)) = m_0(L) \text{CO}^0([E]) = m_0(L) [L] \in CF_{\text{Bott}}^*(L, L)} \quad (9.1)$$

where $\text{CO}^0([E]) = [L]$ since only constant discs contribute, as other solutions are not isolated (the intersection condition $[E]$ is automatic). Recall [12, Lemma 6.7] requires that the monotone Lagrangian L is disjoint from a representative lf-cycle D for $c_1(TE)$, to rule out a contribution of constant discs in the calculation of $\text{CO}(D)$, implying (9.1). We use:

Lemma 9.1 ([36, Lemma 6.4]). *For any orientable Lagrangian L in a monotone symplectic manifold E , there is a representative lf-cycle D of $\text{PD}[c_1(TE)]$ that does not intersect L .*

The idea of the proof, is that a representative for $\text{PD}[c_1(TE)]$ can be obtained from the vanishing set of a generic smooth section a complex line bundle on E with Chern class $c_1(TE)$. Then, using that L is orientable, we can deform this section near L so that it does not vanish at points of L , because $c_1(TE)|_L = \lambda_E \omega|_L = 0$ (using monotonicity of E , and $\omega|_L = 0$).

Remark 1. *For non-orientable L this may fail: $\mathbb{R}P^2 \subset \mathbb{C}P^2$ always intersects $\text{PD}[c_1(T\mathbb{C}P^2)]$.*

Remark 2. *If one imposes the condition $\text{char } \mathbb{K} \neq 2$, then one can drop the assumptions that L is orientable and that E is monotone. Following [46, Prop 3.1], $H^2(E, L) \cong H_{\dim E - 2}(E \setminus L)$ so the Maslov class³ μ_L can be chosen to be disjoint from L . Finally $\text{CO}(c_1(TE)) = \frac{1}{2} \text{CO}(\mu_L)$.*

Lemma 9.2. *Let $L \in \text{Ob}(\mathcal{W}_\lambda(E))$, so $m_0(L) = \lambda$.*

The map $\text{CO}^0 : QC^(E) \rightarrow CW^*(L, L)$, defined in Remark 6.2, satisfies*

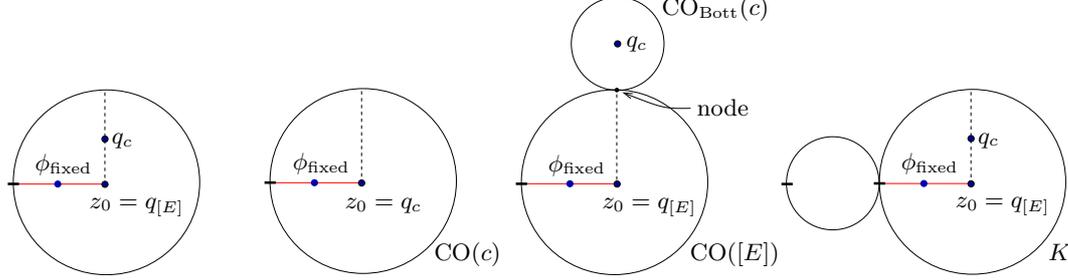
$$\boxed{\text{CO}^0(c_1(TE)) = \lambda e_L + \mu^1(K) \in CW^*(L, L)} \quad \text{for some } K \in CW^*(L, L),$$

where $e_L = \text{CO}^0([E]) \in CW^*(L, L)$ is the cohomological unit (Lemma 8.7).

Proof. Let $c = D$ be the lf-cycle from Lemma 9.1. We run an argument similar to that in 8.8. Consider the first picture below: the auxiliary data is obtained via the forgetful map from the data defining $\text{CO}^0([E])$, as described in Remark 6.2, forgetting the new interior marked point $q_c \in [0, 1]i$. Explicitly, the negative puncture at -1 will define the output in

³The Maslov index $\mu_L : H_2(E, L) \rightarrow \mathbb{Z}$ defines a class $\mu_L \in H^2(E, L)$, and $\mu_L \mapsto 2c_1(TE)$ via $H^2(E, L) \rightarrow H^2(E)$. The Maslov index of a disc bounding L equals the homological intersection number with $\text{PD}[\mu_L]$.

$CF^*(L, L; H) \subset CW^*(L, L)$; the interior marked point $z_0 = 0$ receives the automatic lf-cycle intersection condition $[E]$; the new interior marked point $q_c \in [0, 1)i$ receives the lf-cycle intersection condition c ; and the fixed preferred point is at $-1/2$.



Call K the count of rigid solutions as in picture 1. Now consider a 1-family of solutions. If $q_c \rightarrow 0$ a bubble appears, counting the quantum product $[E] * c = c$, and as it only involves constant solutions we get picture 2, counted by $\text{CO}(c)$. If $q_c \rightarrow i$, in picture 3, a disc bubble appears which is connected to the main component along a node, not a puncture (compare Remark 5.2), in particular the bubble will carry a $PSL(2, \mathbb{R})$ -invariant almost complex structure. This bubble imposes the automatic cycle intersection condition $[L]$ at the node of the main component, so the total count is $m_0(L)$ times the output of the main component, which is $\text{CO}^0([E])$. If q_c converges to a point in $[0, 1)i$, in picture 4, bubbling can occur at -1 , counted by $\mu_{CW^*(L, L)}^1$ of the output of the main component, which is K . \square

Remark. Lemma 9.2 also holds for $\mathcal{F}_\lambda(E)$ taking $\text{CO}^0 : QC^*(E) \rightarrow CF^*(L, L)$ (using the Floer datum model of 5.2). The Floer datum then plays the role of ϕ_{fixed} .

9.3. Producing a nilpotent action of $c_1(TE) - \lambda I$ at the chain level. Let \mathcal{A} be the A_∞ -category $\mathcal{F}_\lambda(E)$ or $\mathcal{W}_\lambda(E)$, and let $\mathcal{M} = \mathcal{A}$ be the diagonal bimodule.

Theorem 9.3. $c = c_1(TE) - \lambda I \in QH^*(E)$ acts nilpotently on $\text{HH}_*(\mathcal{M})$ (for any cycle $w \in \text{HH}_*(\mathcal{M})$, there is an integer N possibly depending on w , such that $\text{HH}_*(\psi_{c^N})(w) = 0$).

Proof. We will construct a pre-morphism $\mathcal{K} \in \text{end}(\mathcal{M})$ such that $(\psi_c + \delta\mathcal{K})^{0|0} = 0$ (at the chain level). Since on homology

$$\text{HH}_*(\psi_c) = \text{HH}_*(\psi_c + \delta\mathcal{K}) \in \text{End}(\text{HH}_*(\mathcal{M})),$$

the Theorem then follows by Lemma 2.4 applied to $\psi_c + \delta\mathcal{K} \in \text{end}(\mathcal{M})$.

Following the left half of the diagram in Theorem 8.2 (which commutes on homology) we obtain a bimodule pre-morphism $\mathcal{K}_1 \in \text{end}(\mathcal{M})$ such that

$$\mu_{\mathcal{M}} \circ_{\text{left}} \text{CO}(c) = \psi_c + \delta\mathcal{K}_1.$$

Define $\mathcal{K}_2 \in \text{CC}^*(\mathcal{M})$ taking $\mathcal{K}_2^0 : \mathbb{K} \rightarrow \mathcal{M}(L, L)$ linear with $\mathcal{K}_2^0(1) = K \in CW^*(L, L)$ as in Lemma 9.2, and $\mathcal{K}_2^r = 0$ for $r \geq 1$. By construction, $(\delta\mathcal{K}_2)^{0|0}(1) = \mu_{\mathcal{M}}^{0|0}(\mathcal{K}_2^0(1)) = -\mu_{\mathcal{A}}^1(K)$, so

$$(\text{CO}(c) + \delta\mathcal{K}_2)^{0|0}(1) = 0 \tag{9.2}$$

by Lemma 9.2, at the chain level. Finally, define $\mathcal{K} = \mathcal{K}_1 + \mu_{\mathcal{M}} \circ_{\text{left}} \mathcal{K}_2 \in \text{end}(\mathcal{M})$. Since $\mu_{\mathcal{M}} \circ_{\text{left}}$ is a chain map, we have

$$\psi_c + \delta\mathcal{K} = \mu_{\mathcal{M}} \circ_{\text{left}} \text{CO}(c) + \delta(\mu_{\mathcal{M}} \circ_{\text{left}} \mathcal{K}_2) = \mu_{\mathcal{M}} \circ_{\text{left}} (\text{CO}(c) + \delta\mathcal{K}_2).$$

The $0|0$ part evaluated on any $\underline{m} \in \mathcal{M}(L, L)$, using Equation (9.2), is zero as required:

$$(-1)^{\deg(\underline{m})} \mu_{\mathcal{M}}^{0|0}((\text{CO}(c) + \delta\mathcal{K}_2)^{0|0}(1), \underline{m}) = 0. \quad \square$$

Remark 9.4. One can equivalently define the module action by the lower half of Theorem 8.2, so $\text{CO}(c)$ acts by cap product on $\text{HH}_*(\mathcal{M})$. To see that Equation 9.2 ensures that the action of $\text{CO}(c) + \delta\mathcal{K}_2$ is nilpotent at the chain level it then suffices to notice that the analogue of Lemma 2.4 holds: each time we apply cap product by $\text{CO}(c) + \delta\mathcal{K}_2$, the cap product from 2.17 must feed at least one of the non-module inputs to $\text{CO}(c) + \delta\mathcal{K}_2$ (otherwise we get zero as the 0 part of the map vanishes), thus shortening the word in $\text{CC}_*(\mathcal{M})$.

9.4. The OC map respects eigensummand decompositions.

Theorem 9.5.

(1) $\text{OC} : \text{HH}_*(\mathcal{F}_\lambda(E)) \rightarrow QH^*(E)$ lands in the eigensummand $QH^*(E)_\lambda$,

$$\text{OC} : \text{HH}_*(\mathcal{F}_\lambda(E)) \rightarrow QH^*(E)_\lambda.$$

(This also holds if we replace E by a closed monotone symplectic manifold B .)

(2) $\text{OC} : \text{HH}_*(\mathcal{W}(E)) \rightarrow SH^*(E)$ lands in the eigensummand $SH^*(E)_\lambda$,

$$\text{OC} : \text{HH}_*(\mathcal{W}_\lambda(E)) \rightarrow SH^*(E)_\lambda.$$

(3) The first acceleration diagram respects these structures:

$$\begin{array}{ccc} \text{HH}_*(\mathcal{F}_\lambda(E)) & \xrightarrow{\text{HH}_*(\mathcal{A}\mathcal{F})} & \text{HH}_*(\mathcal{W}_\lambda(E)) \\ \text{OC} \downarrow & & \downarrow \text{OC} \\ QH^*(E)_\lambda & \xrightarrow{c^*} & SH^*(E)_\lambda \end{array}$$

Proof. This follows immediately from Theorem 9.3 (and the fact that OC and the acceleration diagram respect the module structures, see 8.11 and 8.13). \square

9.5. The CO map respects eigensummand decompositions. The 0-parts

$$\text{CO}^0 : QH^*(E) \rightarrow HF^*(L, L) \quad \text{CO}^0 : SH^*(E) \rightarrow HW^*(L, L)$$

restricted to an eigensummand $QH^*(E)_\lambda$, respectively $SH^*(E)_\lambda$, vanish unless $m_0(L) = \lambda$. This follows by Lemma 9.2 via Theorem 8.2 as $c_1(TE) - \mu \text{Id}$ acts by multiplication by the scalar $\lambda - \mu \neq 0 \in \Lambda$ on $HF^*(L, L)$ and $HW^*(L, L)$, for $\mu \neq \lambda$. The acceleration diagram respects the above decomposition: for any $L \in \text{Ob}(\mathcal{F}_\lambda(E))$,

$$\begin{array}{ccc} HF^*(L, L) & \xleftarrow{\text{HH}^0(\mathcal{A}\mathcal{F})} & HW^*(L, L) \\ \text{CO}^0 \uparrow & & \uparrow \text{CO}^0 \\ QH^*(E)_\lambda & \xrightarrow{c^*} & SH^*(E)_\lambda \end{array}$$

We claim more generally that $\text{CO} : QH^*(E) \rightarrow \text{HH}^*(\mathcal{F}_\lambda(E))$ factorizes:

$$\text{CO} = \bigoplus_{\lambda} \text{CO}_\lambda : QH^*(E) = \bigoplus_{\lambda} QH^*(E)_\lambda \rightarrow \bigoplus_{\lambda} \text{HH}^*(\mathcal{F}_\lambda(E)). \quad (9.3)$$

If we replace E by a closed monotone symplectic manifold then (9.3) follows directly from Theorem 9.5 and a duality argument relating OC and CO maps (see Sheridan [45, Prop 2.6, Cor 2.10]). In the non-compact setup, this duality no longer holds; OC is defined using ordinary cycles but CO is defined using locally-finite cycles. Nick Sheridan and the anonymous referee pointed out to us a way to adapt Theorem 9.5 directly to CO bypassing this dualization.

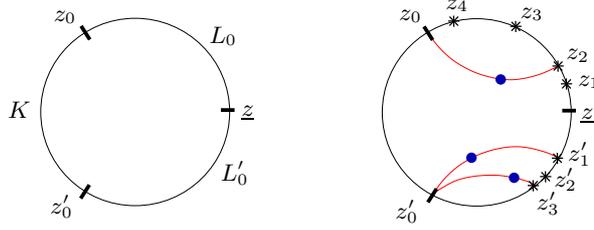


FIGURE 5. $\Delta^{0|0} : CF^*(L'_0, L_0; \underline{w}H) \rightarrow CF^*(K, L_0; w_0H) \otimes CF^*(L'_0, K; w'_0H)$
and $\Delta^{4|3}(x_4, x_3, \mathbf{q}x_2, x_1, \underline{x}, \mathbf{q}x'_1, x'_2, \mathbf{q}x'_3)$.

Theorem 9.6. *If $\mu \neq \lambda$, then the following maps vanish,*

$$\begin{aligned} \text{CO}_{\lambda, \mu} &: QH^*(E)_\mu \rightarrow \text{HH}^*(\mathcal{F}_\lambda(E)), \\ \text{CO}_{\lambda, \mu} &: SH^*(E)_\mu \rightarrow \text{HH}^*(\mathcal{W}_\lambda(E)). \end{aligned}$$

Proof. Let $c = c_1(TE)$. Given $q \in QH^*(E)_\mu$ (respectively $SH^*(E)_\mu$), we know $(c - \mu)^N * q = 0$ for some $N \in \mathbb{N}$. Since CO is an algebra homomorphism (see 8.10),

$$\text{CO}(c - \mu)^N * \text{CO}(q) = 0 \in \text{HH}^*(\mathcal{F}_\lambda(E)) \quad (\text{respectively } \text{HH}^*(\mathcal{W}_\lambda(E))). \quad (9.4)$$

By Lemma 9.2, $\text{CO}^0(c - \mu)^N = (\lambda - \mu)^N \cdot e_L \in CF^*(L, L)$ (respectively $CW^*(L, L)$) where e_L is the cohomological unit and $m_0(L) = \lambda$. If $\mu \neq \lambda$ then Lemma 2.5 applied to $\varphi = (\lambda - \mu)^{-N} \text{CO}(c - \mu)^N$ forces $\text{CO}(c - \mu)^N * \cdot$ to be an isomorphism, so $\text{CO}(q) = 0$ by (9.4). \square

Corollary 9.7. *Equation (9.3) holds. Also the restriction of CO to $\bigoplus SH^*(E)_\lambda$ decomposes:*

$$\text{CO} = \bigoplus_{\lambda} \text{CO}_{\lambda} : \bigoplus_{\lambda} SH^*(E)_\lambda \rightarrow \bigoplus_{\lambda} \text{HH}^*(\mathcal{W}_\lambda(E))$$

Moreover, the acceleration diagram respects the eigensummand decomposition: the following is a commutative diagram of unital algebra homomorphisms (via Theorems 8.2 and 8.11).

$$\begin{array}{ccc} \text{HH}^*(\mathcal{F}_\lambda(X)) & \xleftarrow{\text{HH}^*(\mathcal{A}\mathcal{F})} & \text{HH}^*(\mathcal{W}_\lambda(X)) \\ \text{CO} \uparrow & & \uparrow \text{CO} \\ QH^*(E)_\lambda & \xrightarrow{c^*} & SH^*(E)_\lambda \end{array}$$

In particular, $\text{CO}^0 : SH^(E)_\lambda \rightarrow \text{HW}^*(L, L)$ is unital, where $L \in \text{Ob}(\mathcal{W}_\lambda(E))$.*

10. THE COPRODUCT

10.1. The coproduct as a bimodule map. Continue with the notation $\mathcal{A} = \mathcal{W}(E)$ and $\mathcal{M} = \mathcal{A} = \mathcal{W}(E)$ from 8.1. Fix an object K in $\text{Ob}(\mathcal{A})$. Let \mathcal{L}, \mathcal{R} denote the Yoneda \mathcal{A} -modules

$$\begin{aligned} \mathcal{L}(X) &= \text{hom}_{\mathcal{A}}(K, X) = CW^*(K, X), \\ \mathcal{R}(X) &= \text{hom}_{\mathcal{A}}(X, K) = CW^*(X, K), \end{aligned}$$

where $X \in \text{Ob}(\mathcal{A})$. Recall \mathcal{R} represents the object K (see [41, Sec.(2g)]). By 2.20,

$$\text{HH}_*(\mathcal{A}, \mathcal{L} \otimes \mathcal{R}) \cong H^*(\mathcal{R} \otimes_{\mathcal{A}} \mathcal{L}).$$

We want to construct a *coproduct* Δ as an \mathcal{A} -bimodule map

$$\Delta : \mathcal{M} \rightarrow \mathcal{L} \otimes \mathcal{R},$$

in particular $H^*(\Delta^{0|0}) : \text{HW}^*(L'_0, L_0) \rightarrow \text{HW}^*(K, L_0) \otimes \text{HW}^*(L'_0, K)$ is the usual coproduct.

10.2. **The auxiliary data for the coproduct.** Consider domains S as in the second picture in Figure 5. We consider a universal family $\mathcal{R}^{r|s, \mathbf{p}, \mathbf{p}'} \rightarrow \mathcal{R}^{r+s+3}$ in which there are:

- two distinguished negative boundary punctures z_0, z'_0 on S for the outputs of Δ ,
 - one distinguished positive puncture \underline{z} for the input of Δ from the module $CF^*(L'_0, L_0; \underline{w}H)$.
- We call z_0 the *left output* and z'_0 the *right output*. The distinguished punctures \underline{z}, z_0, z'_0 are fixed to lie at $1, e^{2\pi i/3}, e^{4\pi i/3}$, after identifying \overline{S} with the disc \mathbb{D} . This kills the reparametrization group of \mathcal{R}^{r+s+3} and the local coordinates for \mathcal{R}^{r+s+3} are given by the angular coordinates of the positive punctures z_1, \dots, z_r , called *left punctures*, lying on the arc \underline{z}, z_0 and of the positive punctures z'_s, \dots, z'_1 , called *right punctures*, lying on the arc z'_0, \underline{z} . The punctures are ordered as shown in Figure 5. The puncture \underline{z} will be considered both a left and a right puncture. The \mathbf{p} -labelling map now breaks into two maps

$$\mathbf{p} : F \rightarrow \{r, \dots, 1, \infty\} \quad \text{and} \quad \mathbf{p}' : F' \rightarrow \{\infty, 1, \dots, s\} \quad (10.1)$$

where F, F' are finite sets of indices f, f' and where we use ∞ to denote $\underline{z} = z_\infty$. These give rise to labels $p_f, p'_{f'}$ (possibly not distinct). The map \mathbf{p} encodes the presence of preferred points ϕ_f on geodesics connecting the left-output puncture z_0 with a left-input puncture $z_r, \dots, z_1, z_\infty$, whereas the \mathbf{p}' map encodes the presence of preferred points $\phi_{f'}$ on geodesics connecting the right-output z'_0 and the right-inputs $z_\infty, z'_1, \dots, z'_s$. The module input puncture $\underline{z} = z_\infty$ can play both left/right roles: there can be geodesics connecting it to either z_0 or z'_0 .

For the non-module inputs, the α -forms are separated into two classes: α_k forms for the pairs (z_0, z_k) and α'_k forms for the pairs (z'_0, z'_k) . For the module input puncture \underline{z} , we have a form α_L associated to the left-pair (z_0, \underline{z}) and a form α_R associated to the right-pair (z'_0, \underline{z}) . The weight data is as follows, where all entries are integers ≥ 1 :

$$\mathbf{w} = (w_r, \dots, w_1) \quad \mathbf{w}' = (w'_1, \dots, w'_s) \quad \underline{w} = \underline{w}_L + \underline{w}_R.$$

The total one-form γ on S is:

$$\gamma = \sum_{k=1}^r w_k \alpha_k + \underline{w}_L \alpha_L + \underline{w}_R \alpha_R + \sum_{k=1}^s w'_k \alpha'_k + \sum_{f \in F} \beta_f + \sum_{f' \in F'} \beta_{f'},$$

which satisfies $d\gamma \leq 0$; $d\gamma = 0$ near ∂S ; the pull-back of γ to ∂S is 0. The output weights are:

$$w_0 = \sum w_k + \underline{w}_L + |F| \quad \text{and} \quad w'_0 = \sum w'_k + \underline{w}_R + |F'|.$$

Near the ends, in the relevant parametrization, γ has the form: $w_k dt, \underline{w} dt, w'_k dt, w_0 dt, w'_0 dt$ respectively near $z_k, \underline{z}, z'_k, z_0, z'_0$. This is consistent with Stokes' Theorem: $0 \leq -\int d\gamma = w_0 + w'_0 - \sum w_k - \sum w'_k - \underline{w} = |F| + |F'|$.

By mimicking 3.14, we obtain a moduli space $\mathcal{R}^{r|s, \mathbf{p}, \mathbf{p}'}(\mathbf{x}, \underline{x}, \mathbf{x}')$ of solutions $u : S \rightarrow E$ of $(du - X \otimes \gamma)^{0,1} = 0$ satisfying the relevant asymptotic and boundary conditions. Arguing as in [8, Lemma 3.7], we can always assume that there is never more than one preferred point on a geodesic, since for symmetry reasons the contributions of isolated solutions will cancel when \mathbf{p} or \mathbf{p}' is not injective. When defining Δ as a sum of contributions $\Delta^{r|s, \mathbf{p}, \mathbf{w}, \underline{w}_L, \underline{w}_R, \mathbf{w}', \mathbf{p}'}$, we only sum over inclusions of subsets

$$\mathbf{p} : F \subset \{r, \dots, 1, \infty\} \quad \text{and} \quad \mathbf{p}' : F' \subset \{\infty, 1, \dots, s\}. \quad (10.2)$$

Because \underline{z} can arise for both \mathbf{p}, \mathbf{p}' (i.e. there are two geodesics connecting to \underline{z}), we first need a chain model for \mathcal{M} having two separate \mathbf{q} variables.

10.3. **A chain level direct limit using a double telescope construction.** Let

$$\overline{CW}(X, X') = \bigoplus_{n \geq 2} \bigoplus_{a+b=n, a \geq 1, b \geq 1} CF^*(X, X'; (a+b)H)[\mathbf{q}_L, \mathbf{q}_R]$$

for $X, X' \in \text{Ob}(\mathcal{A})$. This vector space comes with two gradings $a, b \geq 1$ in \mathbb{N} , and it has two variables $\mathbf{q}_L, \mathbf{q}_R$ in degree -1 satisfying $\mathbf{q}_L^2 = 0, \mathbf{q}_R^2 = 0, \mathbf{q}_L \mathbf{q}_R = -\mathbf{q}_R \mathbf{q}_L$.

The differential is defined linearly and equivariantly with respect to $\partial_{\mathbf{q}_L}, \partial_{\mathbf{q}_R}$. Summing along the row, the differential of $x + \mathbf{q}_L y + \mathbf{q}_R z + \mathbf{q}_L \mathbf{q}_R w$ has the following contributions:

Input	Contributions to $\mu^1(\text{Input})$			
x	$(-1)^{ x } \partial x$			
$\mathbf{q}_L y$	$(-1)^{ y } (\mathfrak{K}_L - \text{id}) y$	$(-1)^{ y } \mathbf{q}_L \partial y$		
$\mathbf{q}_R z$	$(-1)^{ z } (\mathfrak{K}_R - \text{id}) z$		$(-1)^{ z } \mathbf{q}_R \partial z$	
$\mathbf{q}_L \mathbf{q}_R w$		$-(-1)^{ w } \mathbf{q}_R (\mathfrak{K}_L - \text{id}) w$	$(-1)^{ w } \mathbf{q}_L (\mathfrak{K}_R - \text{id}) w$	$(-1)^{ w } \mathbf{q}_L \mathbf{q}_R \partial w$

where we mimic Definition 3.4, in particular the \mathfrak{K} -maps are continuation maps:

$$\begin{aligned} \mathfrak{K}_L : CF^*(X, X'; (a+b)H)[\mathbf{q}_L, \mathbf{q}_R] &\rightarrow CF^*(X, X'; (a+1+b)H)[\mathbf{q}_L, \mathbf{q}_R] \\ \mathfrak{K}_R : CF^*(X, X'; (a+b)H)[\mathbf{q}_L, \mathbf{q}_R] &\rightarrow CF^*(X, X'; (a+b+1)H)[\mathbf{q}_L, \mathbf{q}_R] \end{aligned} \quad (10.3)$$

where we stipulate that on $\overline{CW}^*(X, X')$ the map \mathfrak{K}_L increases a by 1, and \mathfrak{K}_R increases b by 1. At the analysis level (i.e. choices of auxiliary data) we construct the two maps in (10.3) to be precisely the same if one ignores the grading by (a, b) :

$$\mathfrak{K}_L = \mathfrak{K}_R \quad (\text{in particular } \mathfrak{K}_L \circ \mathfrak{K}_R = \mathfrak{K}_R \circ \mathfrak{K}_L).$$

An easy verification shows that $\mu^1 \circ \mu^1 = 0$, so $\overline{CW}^*(X, X')$ is a chain complex.

Remark. A double telescope at the topological level would usually result in an asymmetry in the chain complex. This would be present above if we had chosen more generally to consider the sum over $CF^*(X, X'; aH_L + bH_R)[\mathbf{q}_L, \mathbf{q}_R]$ for two different Hamiltonians H_L, H_R , with $\mathfrak{K}_L, \mathfrak{K}_R$ increasing respectively a, b by one. In this case $\mathfrak{K}_L, \mathfrak{K}_R$ are genuinely different, and to ensure $\mu^1 \circ \mu^1 = 0$ an additional contribution $(-1)^{|w|} \mathbf{q}_L \mathbf{q}_R T(w)$ would be required in μ^1 for a chain homotopy T satisfying $\mathfrak{K}_L \circ \mathfrak{K}_R - \mathfrak{K}_R \circ \mathfrak{K}_L = \partial \circ T + T \circ \partial$.

Define an ‘‘inclusion’’ $i_L : CW^*(X, X') \rightarrow \overline{CW}^*(X, X')$ by sending $\mathbf{q} \mapsto \mathbf{q}_L$ and

$$CF^*(X, X'; nH) \rightarrow CF^*(X, X'; (n+1)H), \quad x \mapsto \mathfrak{K}_R(x).$$

Define the projection $\text{pr} : \overline{CW}^*(X, X') \rightarrow CW^*(X, X')$ by sending $\mathbf{q}_L \mapsto \mathbf{q}, \mathbf{q}_R \mapsto \mathbf{q}$, and

$$CF^*(X, X'; (a+b)H) \rightarrow CF^*(X, X'; nH), \quad x \mapsto x \quad \text{where } n = a + b.$$

Observe that $\text{pr} \circ i_L = \mathfrak{K}_R|_{CW^*(X, X')}$.

Lemma 10.1. $\overline{CW}^*(X, X')$ is quasi-isomorphic to $CW^*(X, X')$, intertwining the action of $\partial_{\mathbf{q}}^L + \partial_{\mathbf{q}}^R$ on $\overline{CW}^*(X, X')$ with the action of $\partial_{\mathbf{q}}$ on $CW^*(X, X')$. Explicitly, i_L and pr are quasi-isomorphisms, and $\text{pr} \circ i_L$ and $i_L \circ \text{pr}$ are homotopic to identity maps.

Proof. Observe that, taking x, y, z to be ∂ -cycles, implies the identifications $[\mathfrak{K}_L y] = [y]$ and $[\mathfrak{K}_R z] = [z]$ on cohomology, and $[\mathfrak{K}_L \mathfrak{K}_R x] = [\mathfrak{K}_R \mathfrak{K}_L x]$ already at the chain level. In particular, any ∂ -cycle x defines a class in the cohomology $\overline{HW}^*(X, X')$. Although many copies of this appear for given n , since $x \in CF^*(X, X'; (a+b)H)$ for all $a, b \geq 1$ with $a+b = n$, these are all identified on cohomology since we can compare any two copies by applying $\mathfrak{K}_R, \mathfrak{K}_L$ appropriately many times until they both lie in the same summand of $CW^*(X, X')$. So there is a copy of each cohomology representative $[x] \in HF^*(L, L'; nH)$ in grading $(a, b) = (n-1, 1)$ (for $n \geq 2$), thanks to the cohomological identifications provided by the \mathfrak{K} -maps.

We now run an argument analogous to [8, Sec.3.7].

Abbreviate $F_{x,y} = CF^*(X, X'; (x+y)H)$ for integers $x, y \geq 1$. Abbreviate

$$C_{k,\ell} = \bigoplus_{1 \leq x < k} \bigoplus_{1 \leq y < \ell} F_{x,y}[\mathbf{q}_L, \mathbf{q}_R] \oplus \bigoplus_{1 \leq x < k} F_{x,\ell}[\mathbf{q}_L] \oplus \bigoplus_{1 \leq y < \ell} F_{k,y}[\mathbf{q}_R] \oplus F_{k,\ell}.$$

Imagine these summands arranged as a grid of squares in the xy -plane, so \mathfrak{K}_R maps upwards by one square, \mathfrak{K}_L maps to the right by one square, and all non- \mathfrak{K} terms in the differential map from a square to itself. Thus, $C_{k,\ell}$ are subcomplexes which exhaust $CW^*(X, X')$. The key observation is that if we take any one of the squares (i.e. a summand above), and we consider the differential obtained by dropping \mathfrak{K} -terms, then this square becomes acyclic.

Consider the smaller subcomplex C' obtained by removing the summand $F_{1,1}[\mathbf{q}_L, \mathbf{q}_R]$. Then C' is quasi-isomorphic to $C_{k,\ell}$, because the quotient $C_{k,\ell}/C'$ is the acyclic square in position (1,1). More generally, up to quasi-isomorphism we can inductively remove one square at a time (always removing a square with minimal occurring x value or minimal y or both minimal). Thus $C_{k,\ell}$ is quasi-isomorphic to $F_{k,\ell}$. Just as in the original argument [8, Sec.3.7], one checks that $\mathfrak{K}_R : F_{k,\ell} \rightarrow F_{k,\ell+1}$ (respectively $\mathfrak{K}_L : F_{k,\ell} \rightarrow F_{k+1,\ell}$) commutes up to chain homotopy with the inclusions $C_{k,\ell} \subset C_{k,\ell+1}$ (respectively $C_{k,\ell} \subset C_{k+1,\ell}$); the chain homotopy is $x \mapsto \mathbf{q}_R x$ (respectively $x \mapsto \mathbf{q}_L x$). On cohomology, the direct limit of the cohomologies $H^*(F_{k,\ell})$ over the maps $\mathfrak{K}_R, \mathfrak{K}_L$ is $HW^*(X, X')$ essentially by definition. Thus $H^*(CW^*(X, X'); \mu^1) \cong HW^*(X, X')$ are canonically isomorphic. \square

10.4. A new construction of the diagonal bimodule. We will define a new \mathcal{A} -bimodule $\overline{\mathcal{M}}$, which is quasi-isomorphic to \mathcal{M} , satisfying

$$\overline{\mathcal{M}}(X, X') = \overline{CW}^*(X', X).$$

The composition maps $\mu_{\overline{\mathcal{M}}}^{r|s} : \mathcal{A}(L_r, \dots, L_0) \otimes \overline{\mathcal{M}}(L_0, L'_0) \otimes \mathcal{A}(L'_0, \dots, L'_s) \rightarrow \overline{\mathcal{M}}(L_r, L'_s)$ are defined by mimicking 3.16 using the same auxiliary data, so that we count the same Floer solutions as for $\mu_{\mathcal{M}}^{r|s}$, but we now “separate” the total output by distributing it according to weights. Explicitly, $\mu_{\overline{\mathcal{M}}}^{r|s}$ on a generator

$$\mathbf{q}^{i_r} x_r \otimes \dots \otimes \mathbf{q}^{i_1} x_1 \otimes \mathbf{q}_L^{i_\infty} \mathbf{q}_R^{i'_\infty} \underline{m} \otimes \mathbf{q}^{i'_1} y_1 \otimes \dots \otimes \mathbf{q}^{i'_s} y_s \quad (10.4)$$

with input weights $w_r, \dots, w_1, \underline{w}_L + \underline{w}_R, w'_1, \dots, w'_s$ (where the module input has grading $(a, b) = (\underline{w}_L, \underline{w}_R)$) will contribute to the output summand

$$CF^*(L'_s, L_r; (w_{0,L} + w_{0,R})H) \subset \overline{\mathcal{M}}(L_r, L'_s)$$

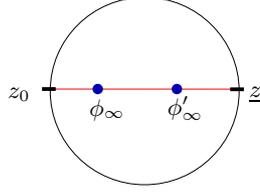
in grading $w_{0,L} = \sum w_k + \underline{w}_L + |F|$ and $w_{0,R} = \sum w'_k + \underline{w}_R + |F'|$ where F, F' are defined as in (10.2). Note we use two \mathbf{p} -maps as in (10.1), one for the left punctures and one for the right punctures. So we place a preferred point on the geodesic connecting (z_0, z_k) , (respectively (z_0, z'_k)) if $i_k = 1$ (respectively $i'_k = 1$). To define the $\mathbf{q}_L, \mathbf{q}_R, \mathbf{q}_L \mathbf{q}_R$ outputs of $\mu_{\overline{\mathcal{M}}}^{r|s}$, we impose $\partial_{\mathbf{q}}^L$ and $\partial_{\mathbf{q}}^R$ equivariance. Here, we define the left and right operators $\partial_{\mathbf{q}}^L, \partial_{\mathbf{q}}^R$ to act respectively on the left and right inputs, associated to \mathbf{p} and \mathbf{p}' . These should be viewed in grading $+1$ and as acting from the left, so:

$$\begin{aligned} \partial_{\mathbf{q}}^L(x_r \otimes \dots \otimes x_1 \otimes \underline{x} \otimes x'_1 \otimes \dots \otimes x'_s) &= \partial_{\mathbf{q}}(x_r \otimes \dots \otimes x_1 \otimes \underline{x}) \otimes x'_1 \otimes \dots \otimes x'_s \\ \partial_{\mathbf{q}}^R(x_r \otimes \dots \otimes x_1 \otimes \underline{x} \otimes x'_1 \otimes \dots \otimes x'_s) &= (-1)^{\sigma(x)_r} x_r \otimes \dots \otimes x_1 \otimes \partial_{\mathbf{q}}(\underline{x} \otimes x'_1 \otimes \dots \otimes x'_s) \end{aligned}$$

and for the outputs we define $\partial_{\mathbf{q}}^L(x, y) = \partial_{\mathbf{q}}(x) \otimes y$ and $\partial_{\mathbf{q}}^R(x, y) = (-1)^{\deg(x)x} \otimes \partial_{\mathbf{q}}(y)$.

The A_∞ -relations required for $\overline{\mathcal{M}}$ to be an \mathcal{A} -bimodule hold because we know that these relations hold for the diagonal \mathcal{A} -bimodule \mathcal{M} and because the gluing/breaking involved in compositions respects the grading on weights $(a, b) = (\underline{w}_L, \underline{w}_R)$ described above.

Example. Given $x \in CF^*(X', X; (a+b)H)$, the following picture contributes to the \mathbf{q}^0 -output of $\mu_{\overline{\mathcal{M}}}^{0|0}(\mathbf{q}_L \mathbf{q}_R x)$ lying in $CF^*(X', X; (a+1+b+1)H)$. Here $F = \{\infty\}$, $F' = \{\infty\}$.



Given such a Floer solution, there is another Floer solution contributing to the same output (in the same grading) with $\phi_\infty, \phi_\infty'$ interchanged (analytically it is the same PDE solution, since our auxiliary data is constructed independently of the labelling of the preferred points). The usual symmetry argument (see the Remark about symmetry in 3.15) implies that these two Floer solutions are counted with opposite orientation signs, so they cancel.

Lemma 10.2. *The \mathcal{A} -bimodule $\overline{\mathcal{M}}$ is quasi-isomorphic to the diagonal \mathcal{A} -bimodule \mathcal{M} , intertwining the action of $\partial_{\mathbf{q}}^L + \partial_{\mathbf{q}}^R$ on $\overline{\mathcal{M}}$ with the action of $\partial_{\mathbf{q}}$ on \mathcal{M} .*

An explicit quasi-isomorphism $\mathbf{pr} : \overline{\mathcal{M}} \rightarrow \mathcal{M}$ of bimodules is given by:

$$\begin{aligned} \mathbf{pr}^{r|s} &: \mathcal{A}(X_r, \dots, X_0) \otimes \overline{\mathcal{M}}(X_0, X'_0) \otimes \mathcal{A}(X'_0, \dots, X'_s) \rightarrow \mathcal{M}(X_r, X'_s) \\ \mathbf{pr}^{0|0} &= \mathbf{pr} : \overline{\mathcal{M}}(X_0, X'_0) \rightarrow \mathcal{M}(X_0, X'_0). \\ \mathbf{pr}^{r|s} &= 0 \text{ if } r + s \geq 1. \end{aligned}$$

where \mathbf{pr} is the projection defined in Lemma 10.1.

Proof. It remains to check that the map \mathbf{pr} satisfies the A_∞ -relations described in 2.8. This reduces to checking the identity

$$\mathbf{pr}^{0|0}(\mu_{\overline{\mathcal{M}}}^{r|s}(x_r, \dots, x_1, \underline{m}, y_1, \dots, y_s)) = \mu_{\mathcal{M}}^{r|s}(x_r, \dots, x_1, \underline{\mathbf{pr}^{0|0}(\underline{m})}, y_1, \dots, y_s).$$

This equation simply says that the solutions that $\mu_{\overline{\mathcal{M}}}^{r|s}$ counts are the same as those counted by $\mu_{\mathcal{M}}^{r|s}$ if we ignore the (a, b) -grading of the outputs. This holds by construction. \square

Remark. *A priori this A_∞ -bimodule quasi-isomorphism has a quasi-inverse $\mathcal{M} \rightarrow \overline{\mathcal{M}}$, but it is probably difficult to describe it explicitly.*

10.5. Definition of the coproduct. The \mathcal{A} -bimodule map $\Delta : \overline{\mathcal{M}} \rightarrow \mathcal{L} \otimes \mathcal{R}$ is determined by the $\partial_{\mathbf{q}}^L$ - and $\partial_{\mathbf{q}}^R$ -equivariant maps

$$\begin{aligned} &\mathcal{W}(L_r, \dots, L_0; \mathbf{w}) \otimes CF^*(L'_0, L_0; (\underline{w}_L + \underline{w}_R)H)[\mathbf{q}_L, \mathbf{q}_R] \otimes \mathcal{W}(L'_0, \dots, L'_s; \mathbf{w}') \\ &\longrightarrow CF^*(K, L_r; w_0 H)[\mathbf{q}] \otimes CF^*(L'_s, K; w'_0 H)[\mathbf{q}], \end{aligned}$$

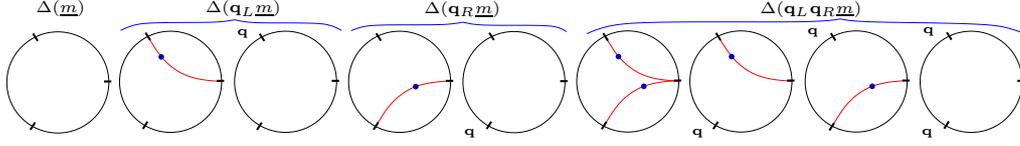
counting Floer solutions for the auxiliary data described in 10.2. Thus, given the input as in (10.4), we require $k \in \mathbf{p}(F)$ iff $i_k = 1$ (respectively $k \in \mathbf{p}'(F')$ iff $i'_k = 1$). Equivalently: there is a \mathbf{q} iff there is a preferred point on the geodesic connecting the relevant puncture with z_0 (respectively z'_0). For the module input we have a left/right \mathbf{q} -variable which keeps track of the presence of a preferred point on the left/right geodesic connecting to $\underline{z} = z_\infty$. The $x_0 \otimes x'_0$ coefficient of the output is then the count with weight $\pm t^{E_{\text{top}}(u)}$ of the isolated solutions $u \in \mathcal{R}^{r|s, \mathbf{p}, \mathbf{p}'}(x_r, \dots, x_1, \underline{x}, x'_1, \dots, x'_s)$.

Remark. *Solutions are counted with a sign $(-1)^\ddagger$ [2, (4.22)] where*

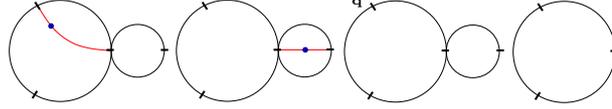
$$\ddagger = k + \sum_{j=1}^s (s - j + 1) \deg(x'_j) + s \deg(\underline{x}) + \sum_{j=1}^r (j + s) \deg(x_j)$$

where k is the Koszul sign obtained from viewing the \mathbf{q} variables as operators of degree -1 acting from the left (compare 3.15 and 8.2).

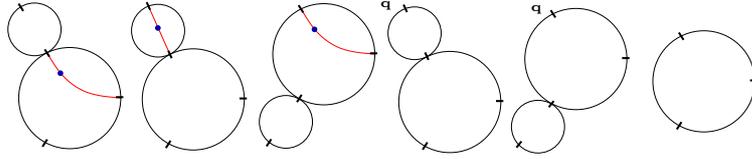
As an illustration, we show the decorated domains counted by $\Delta^{0|0}$, where we label an output with \mathbf{q} if it contributes to the \mathbf{q} -part of that output:



To illustrate the A_∞ -equations, we show the pictures involved in the equation $\Delta^{0|0}(\mu_{\overline{\mathcal{M}}}^{0|0}(\mathbf{q}_L \underline{m})) = \mu_{\mathcal{L} \otimes \mathcal{R}}^1(\Delta^{0|0}(\mathbf{q}_L \underline{m}))$. Ignoring signs, the left-hand side is $\Delta^{0|0}(\mathbf{q}_L \partial \underline{m} + (\mathfrak{K}_L - \text{id})(\underline{m}))$ counting:



and the right-hand side is $(\mu_{\mathcal{A}}^1 \otimes \text{id} + \text{id} \otimes \mu_{\mathcal{A}}^1) \circ \Delta^{0|0}(\mathbf{q}_L \underline{m})$ counting:



With the obvious labelling: pictures 1.1, 1.2, 2.1, 2.2, 2.3 cancel as they are the boundary of a 1-dimensional moduli space; similarly for 1.3, 2.4, 2.5; pictures 1.4, 2.6 cancel due to signs.

That Δ satisfies the A_∞ -equations required from an \mathcal{A} -bimodule map $\overline{\mathcal{M}} \rightarrow \mathcal{L} \otimes \mathcal{R}$ is now an exercise in bookkeeping, compare Remark 4 in 3.15. The key observation is that in codimension 1, either a subset of the punctures (and possibly some preferred points) converges to a common point on the boundary. If the common point is not z_0, z'_0, \underline{z} , it gives rise to a $\mu_{\mathcal{A}}^d$ -term (part of the bar differential for the \mathcal{A} -bimodule $\overline{\mathcal{M}}$). If the common point is z_0 it gives rise to a $\mu_{\mathcal{L} \otimes \mathcal{R}}^{R|0}$ -term; if z'_0 , it gives rise to a $\mu_{\mathcal{L} \otimes \mathcal{R}}^{0|S}$ -term; if \underline{z} , it gives rise to a $\mu_{\overline{\mathcal{M}}}^{R|S}$ -term.

10.6. The OC map using the new diagonal bimodule. We define

$$\text{OC} : \text{CC}_*(\mathcal{A}, \overline{\mathcal{M}}) \rightarrow \text{CC}_*(\mathcal{A}, \mathcal{M}) \rightarrow \text{SC}^*(E)$$

by composing $\text{CC}_*(\text{pr})$ and the usual OC map, where pr was defined in Lemma 10.2. The new OC map can also be defined directly, mimicking $\mu_{\overline{\mathcal{M}}}^{r|s}$, and all weights contribute to the output weight for $\text{SC}^*(E)$ (which only uses a \mathbf{q} -variable; we do not form a double complex).

Remark. We will not need a map CO which lands in $\text{CC}^*(\mathcal{A}, \overline{\mathcal{M}})$. For such a map, one would need to use a double telescope chain model for $\text{SC}^*(E)$.

11. THE GENERATION CRITERION IN THE MONOTONE SETTING

11.1. The generation criterion in the monotone setting.

Definition 11.1 (Split-generating [41, Sec.4c,3j]). *A full subcategory \mathcal{S} split-generates an A_∞ -category \mathcal{A} if for each $A \in \text{Ob}(\mathcal{A})$, there is a twisted complex built from $\text{Ob}(\mathcal{S})$ admitting A as a summand. We sometimes say $\text{Ob}(\mathcal{S})$ split-generates \mathcal{A} .*

Generation Criterion (Abouzaid [2]). *Let $(M, d\theta)$ be an exact symplectic manifold conical at infinity. Suppose \mathcal{S} is a full subcategory of $\mathcal{W}(M)$ such that 1 lies in the image of*

$$\text{HH}_*(\mathcal{S}) \xrightarrow{\text{HH}_*(\text{incl})} \text{HH}_*(\mathcal{W}(M)) \xrightarrow{\text{OC}_M} \text{SH}^*(M).$$

Then \mathcal{S} split-generates $\mathcal{W}(M)$.

By Theorem 9.5, it in fact suffices that $\text{im}(\text{OC}_M)$ contains an $\text{SH}^*(M)$ -invertible:

Lemma 11.2. *If $s \in \text{im}(\text{OC}_M)$ is invertible in $SH^*(M)$ then $1 \in \text{im}(\text{OC}_M)$.*

Proof. $\text{OC}_M(x) = s$ for some $x \in \text{HH}_*(\mathcal{W}(M))$, so since OC_M is an $SH^*(M)$ -module map by Theorem 9.5 (which of course holds also in the exact setup), $\text{OC}_M(s^{-1} \cdot x) = s^{-1} \cdot s = 1$. \square

Theorem 11.3. *Let E be a monotone symplectic manifold conical at infinity. Let \mathcal{S} be a full subcategory of $\mathcal{W}_\lambda(E)$. If*

$$\text{HH}_*(\mathcal{S}) \xrightarrow{\text{HH}_*(\text{incl})} \text{HH}_*(\mathcal{W}_\lambda(E)) \xrightarrow{\text{OC}} SH^*(E)_\lambda,$$

hits an invertible element (in the notation of 9.1), then \mathcal{S} split-generates $\mathcal{W}_\lambda(E)$.

11.2. Outline of the proof of the generation criterion. Fix an object K in $\mathcal{W}(E)$ (as usual, in the monotone setting, we always mean $\mathcal{W}_\lambda(E)$ for a fixed m_0 -value λ). Recall \mathcal{S} is a full subcategory of $\mathcal{W}(E)$. Let \mathcal{L}, \mathcal{R} denote the Yoneda \mathcal{S} -modules

$$\begin{aligned} \mathcal{L}(X) &= \text{hom}_{\mathcal{S}}(K, X) = CW^*(K, X), \\ \mathcal{R}(X) &= \text{hom}_{\mathcal{S}}(X, K) = CW^*(X, K), \end{aligned}$$

where $X \in \text{Ob}(\mathcal{S})$. Recall \mathcal{R} represents the object K (see [41, Sec.(2g)]). By 2.20,

$$\text{HH}_*(\mathcal{S}, \mathcal{L} \otimes \mathcal{R}) \cong H^*(\mathcal{R} \otimes_{\mathcal{S}} \mathcal{L}). \quad (11.1)$$

Recall from Section 10 that the coproduct Δ is an \mathcal{S} -bimodule map

$$\Delta : \mathcal{S} \rightarrow \mathcal{L} \otimes \mathcal{R},$$

where \mathcal{S} is the diagonal \mathcal{S} -bimodule as in 10.4. Consider the composites

$$\begin{aligned} C_1 : \text{HH}_*(\mathcal{S}, \mathcal{S}) &\xrightarrow{\text{OC}} SH^*(E) \xrightarrow{\text{CO}^0} HW^*(K, K) \\ C_2 : \text{HH}_*(\mathcal{S}, \mathcal{S}) &\xrightarrow{\text{HH}_*(\Delta)} \text{HH}_*(\mathcal{S}, \mathcal{L} \otimes \mathcal{R}) \cong H^*(\mathcal{R} \otimes_{\mathcal{S}} \mathcal{L}) \xrightarrow{\mu_{\mathcal{S}}} HW^*(K, K) \end{aligned}$$

using the $\mu_{\mathcal{S}}^n$ to compose all morphisms of the chain complex $\mathcal{R} \otimes_{\mathcal{S}} \mathcal{L}$.

The key ingredient of the proof is to show that these composites are equal (up to a global sign $(-1)^{n(n+1)/2}$, where $\dim E = 2n$), which we will prove in 11.3. The hypothesis that OC hits $1 \in SH^*(E)$ implies C_1 hits the identity $e_K \in HW^*(K, K)$ since CO^0 is a unital algebra homomorphism (see Theorem 8.2). So C_2 hits e_K . For purely algebraic reasons [2, Appendix A] this implies that \mathcal{R} (which represents K) is a summand of a twisted complex built from objects in \mathcal{S} . Thus any object K in $\mathcal{W}(E)$ is split-generated by objects in \mathcal{S} , as required.

Remark 11.4. *Suppose \mathcal{S} is a full subcategory of $\mathcal{W}_\lambda(E)$, such that the composite*

$$\text{CO}^0 \circ \text{OC} : \text{HH}_*(\mathcal{S}) \rightarrow SH^*(E) \rightarrow HW^*(K, K)$$

hits the identity e_K . Suppose K is Floer cohomologically essential: $HW^(K, K) \neq 0$. We claim that $m_0(K) = \lambda$, i.e. $K \in \text{Ob}(\mathcal{W}_\lambda(E))$. Indeed, by Theorem 9.5, $\text{OC}(\text{HH}_*(\mathcal{S}))$ must land in the eigensummand $SH^*(E)_\lambda$, and the restricted map $\text{CO}^0 : SH^*(E)_\lambda \rightarrow HW^*(K, K)$ will vanish unless $m_0(K) = \lambda$ (see 9.5) and it is unital when $m_0(K) = \lambda$, by Corollary 9.7. So it suffices to consider Lagrangians with a single m_0 -value λ in the proof of the generation criterion. Also, in the Generation Criterion and in Lemma 11.2 we only care that $\text{OC}|_{\mathcal{S}}$ hits the idempotent summand of 1 in the generalized eigensummand $SH^*(E)_\lambda$.*

This outline is precisely the argument of Abouzaid [2]. To prove Theorem 1.1 we will therefore only highlight the novelties of the proof in our monotone setup.

11.3. The moduli space of marked annuli. We now prove that the composites C_1, C_2 from 11.2 agree. This involves studying a moduli space of annuli, which in Abouzaid’s case is described in [2, Appendix C4,C5], and in our case will be endowed with additional decorations coming from geodesics and preferred points. Keeping track of this additional data may seem cumbersome at first, but it has the advantage that all gluings are natural: so in fact we will not need the first homotopy of Abouzaid [2, Sec.6.2] which changes the weights so that gluing becomes possible, and we will only need to construct the analogue of his second homotopy [2, Sec.6.3]. Moduli spaces of annuli with n boundary marked points are given Deligne–Mumford compactifications by being viewed as moduli spaces of elliptic curves with real structures via doubling, and an étale forgetful map to $\mathcal{M}_{1,n}$, see [21, Ch.2].

Lemma 11.5. *Inside the Deligne–Mumford moduli space of annuli with two marked points, there is a real one-dimensional submanifold (not transverse to the boundary) whose boundary points are depicted in the first and last pictures of Figure 6 (ignoring decorations).*

Sketch Proof. For $r \in (1, \infty)$, consider the subfamily defined by

$$\varepsilon_r = \{z \in \mathbb{C} : 1 \leq |z| \leq r\} \setminus \{-1, +r\}. \tag{11.2}$$

(Recall the space of conformal structures on annuli in \mathbb{C} centred at 0 is parametrized by the ratio of the outer and inner radius.) By doubling up the annulus, (11.2) determines a real 1-family in $\mathcal{M}_{1,2}$ parametrised by e^r . Abouzaid shows in [2, Appendix C.4] that the boundary strata carry two nodes (even though this would typically be a complex codimension two phenomenon). When $r \rightarrow \infty$, one of the two nodes forms in the interior of the annulus, thus in the limit the annulus breaks into two discs joined at an interior point as in picture 1. When $r \rightarrow 1$, the two nodes arise as a limit of two pairs of boundary points of the annulus connected by two shrinking geodesic segments (the dotted lines in picture 3). In the limit this yields two discs joined at two nodes as in picture 5. \square

Analogous families of annuli can be constructed when there are several marked points on the outer circle $|z| = r$, which we view as families lying over the family (11.2) via the forgetful map. Thus the punctures at $-1, +r$ are distinguished (marked by dashes in Figure 6). The dotted lines in pictures 2 and 3 of Figure 6 are examples of hyperbolic geodesics which shrink in the limit to give respectively pictures 1 and 5. In picture 2, we work in the (s, t) plane, where $z = e^{s+it}$ parametrizes the annulus, so we’ve passed to the cover $\mathbb{R} \rightarrow \mathbb{R}/\mathbb{Z}$ of the angular variable. The limit $r \rightarrow \infty$ corresponds to stretching the strip in picture 2 near the dotted line. Picture 4 is biholomorphic to 3, and shrinking the dotted lines in 3 corresponds to stretching the two handles in 4 (corresponding to the limit $r \rightarrow 1$). The dotted lines in pictures 2 and 3 play the same role as the “cuts” in Figure 1 of 3.11 and Figure 2 of 3.15.

Observe our moduli space of annuli in picture 2 or 3 of Figure 6. The distinguished boundary punctures are $\underline{z} = -1$ (which will be the output) and $\underline{z}' = \underline{z}_\infty = +r$ (which will be the bimodule input), and the other boundary punctures z_1, \dots, z_n are ordered counterclockwise on the outer boundary. Mimicking the construction of the coproduct, the auxiliary data involves two \mathbf{p} -maps as in (10.1) which keep track of the presence of preferred points on geodesics joining \underline{z} with the other punctures, with the important feature that there are two types of (deformed) geodesics joining $\underline{z}, \underline{z}'$ (in Figure 6, the *left geodesic* is the one carrying ϕ_∞ , the *right geodesic* carries ϕ'_∞). The geodesics are deformed so that they come in orthogonally to the boundary. We have weights w_1, \dots, w_n associated to the non-module inputs (coming from \mathcal{A}), and a graded weight $(a, b) = (\underline{w}'_L, \underline{w}'_R)$ for the module input (coming from the diagonal bimodule $\mathcal{S} \subset \overline{\mathcal{M}}$). There are closed forms $\underline{\alpha}_L, \underline{\alpha}_R, \alpha_1, \dots, \alpha_n$; sub-closed forms $\beta_f, \beta_{f'}$. These define a total 1-form $\gamma = \sum w_k \alpha_k + \underline{w}'_L \underline{\alpha}_L + \underline{w}'_R \underline{\alpha}_R + \sum \beta_f + \sum \beta_{f'}$, which pulls back to $w_k dt, (\underline{w}'_L + \underline{w}'_R) dt, \underline{w} dt$ near $z_k, \underline{z}', \underline{z}$, where $\underline{w} = \sum w_k + \underline{w}'_L + \underline{w}'_R + |F| + |F'|$.

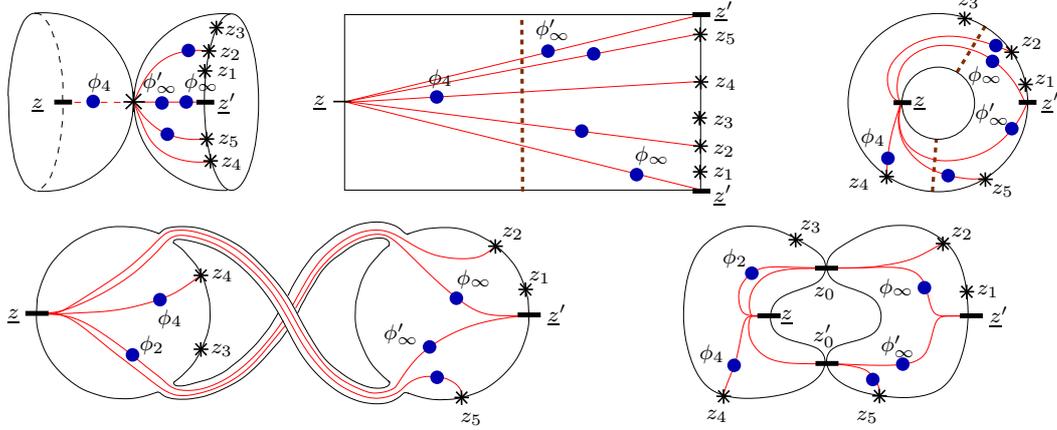


FIGURE 6. Breaking analysis for the moduli space of annuli.

As usual, the auxiliary data is constructed consistently via 3.13, so that on the broken configurations we are using the auxiliary data previously constructed for other operations. So in picture 1, the left disc-bubble carries the data of the $\text{CO}^0 : SC^*(E) \rightarrow CW^*(K, K)$ map, the right disc-bubble carries the data of the $\text{OC} : \text{CC}_*(\mathcal{A}, \overline{\mathcal{M}}) \rightarrow SC^*(E)$ map constructed in 10.6. In picture 5 the left bubble carries the data for $\mu_{\mathcal{A}}^*$ constructed in 3.15, and the right bubble carries the data for $\Delta : \overline{\mathcal{M}} \rightarrow \mathcal{L} \otimes \mathcal{R}$ constructed in 10.5.

Consider a 1-dimensional moduli space of solutions $u : S \rightarrow E$ of $(du - X \otimes \gamma)^{0,1} = 0$ defined on a fixed annulus S , and satisfying the relevant boundary and asymptotic conditions. Pictures 1 and 5 of Figure 6 are two examples of boundary points of this moduli space. We conclude that picture 1 is a contribution to the coefficient of \underline{x} in⁴

$$\text{CO}^0 \circ \text{OC}(\mathbf{q}_L \mathbf{q}_R \underline{x}', \mathbf{q}x_5, \mathbf{q}x_4, x_3, \mathbf{q}x_2, x_1),$$

whereas picture 5 is a contribution to the coefficient of \underline{x} in⁵

$$\mu_S^3 \circ \mathcal{T} \left(\underline{\Delta}^{2|1}(\mathbf{q}x_2, x_1, \mathbf{q}_L \mathbf{q}_R \underline{x}', \mathbf{q}x_5) \otimes \mathbf{q}x_4 \otimes x_3 \right),$$

where \mathcal{T} is the reordering isomorphism from 2.20 which induces (11.1) (notice we are inputting into \mathcal{T} a typical summand of $\text{CC}_*(\Delta)(\mathbf{q}x_2, x_1, \mathbf{q}_L \mathbf{q}_R \underline{x}', \mathbf{q}x_5, \mathbf{q}x_4, x_3)$; see Lemma 2.1).

This example now easily generalizes to show that the boundary points of such 1-dimensional moduli spaces include the counts involved in the maps $C_1 = \text{CO}^0 \circ \text{OC}$ and $C_2 = \mu_S^{\max} \circ \mathcal{T} \circ \text{CC}_*(\Delta)$. There are however also two other contributions which we now describe.

Let $\varphi : \text{CC}_*(\mathcal{A}, \mathcal{S}) \rightarrow CW^*(K, K)$ denote the weighted count of isolated solutions u defined on an annulus as above. In a 1-family, when a subset of the outer boundary punctures converges to a common point, we obtain bubbling configurations counted by $\varphi \circ b$ where b is the bar differential for $\text{CC}_*(\mathcal{A}, \mathcal{S})$. Secondly, when bubbling occurs on the inner boundary circle of the annulus, so at $\underline{z} = -1$, the broken configurations are counted by $\mu_{CW^*(K, K)}^1 \circ \varphi$.

We conclude that (up to signs) C_1 and C_2 differ at the chain level by $\varphi \circ b - \mu_{CW^*(K, K)}^1 \circ \varphi$, therefore they agree on cohomology. This concludes the proof of Theorem 11.3.

Technical Remark. *A more careful treatment of orientation signs [2, Lemma 6.8] shows*

⁴By $\partial_{\mathbf{q}_R}$ -equivariance, the $(\mathbf{q}^0, \mathbf{q}^0)$ -output of $\text{OC}(\mathbf{q}_L \mathbf{q}_R \underline{x}', \mathbf{q}x_5, x_4, x_3, \mathbf{q}x_2, x_1)$ contributes to the \mathbf{q}^1 -coefficient of $\text{OC}(\mathbf{q}_L \mathbf{q}_R \underline{x}', \mathbf{q}x_5, \mathbf{q}x_4, x_3, \mathbf{q}x_2, x_1)$. Compare the Remarks in 3.11 and 3.15.

⁵By $\partial_{\mathbf{q}_L}$ -equivariance, the $(\mathbf{q}^0, \mathbf{q}^0)$ -output of $\underline{\Delta}^{2|1}(x_2, x_1, \mathbf{q}_L \mathbf{q}_R \underline{x}', \mathbf{q}x_5)$ contributes to the $(\mathbf{q}^1, \mathbf{q}^0)$ -coefficient of $\underline{\Delta}^{2|1}(\mathbf{q}x_2, x_1, \mathbf{q}_L \mathbf{q}_R \underline{x}', \mathbf{q}x_5)$. Compare the Remarks in 3.11 and 3.15.

that $C_1 - (-1)^{n(n+1)/2}C_2 = \varphi \circ b + (-1)^n \mu_{CW^*(K,K)}^1 \circ \varphi$, so the maps C_1 and C_2 differ on homology by $(-1)^{n(n+1)/2}$, where $\dim E = 2n$.

11.4. Generation for the compact category. When working with $\mathcal{F}(E)$ or with $\mathcal{F}(B)$ for a closed monotone symplectic manifold B , the analogue of Theorem 11.3 holds, and the proof is analogous (indeed simpler, since one no longer needs a telescope chain model and no auxiliary forms α_k, β_f are in play, beyond using the Floer datum machinery of 5.2).

Theorem 11.6. *Let M be a closed monotone symplectic manifold, or a monotone symplectic manifold conical at infinity. Let \mathcal{S} be a full subcategory of $\mathcal{F}_\lambda(M)$. If*

$$\mathrm{HH}_*(\mathcal{S}) \xrightarrow{\mathrm{HH}_*(\mathrm{incl})} \mathrm{HH}_*(\mathcal{F}_\lambda(M)) \xrightarrow{\mathrm{OC}} QH^*(M)_\lambda,$$

hits an invertible element (in the notation of 9.1), then \mathcal{S} split-generates $\mathcal{F}_\lambda(M)$.

12. APPLICATIONS

12.1. Toric varieties and the Landau-Ginzburg superpotential. We assume the reader has some familiarity with the relevant literature (Cho-Oh [17], Auroux [12], Fukaya-Oh-Ohta-Ono [22]), and we refer to [36, Appendix] for details. Let (X, ω_X) be a closed real $2n$ -dimensional symplectic manifold together with an effective Hamiltonian action of the n -torus $U(1)^n$. This action determines a *moment map* $\mu_X : X \rightarrow \mathbb{R}^n$ whose image $\Delta = \mu_X(X) \subset \mathbb{R}^n$ is a convex polytope, called the *moment polytope*. By Delzant's theorem Δ determines, up to isomorphism, (X, ω_X) together with the action. The moment polytope has the form

$$\Delta = \{y \in \mathbb{R}^n : \langle y, e_i \rangle \geq \lambda_i \text{ for } i = 1, \dots, r\}, \quad (12.1)$$

where $\lambda_i \in \mathbb{R}$ are parameters and $e_i \in \mathbb{Z}^n$ are the primitive inward-pointing normal vectors to the codimension 1 faces of Δ . The *Landau-Ginzburg superpotential* is

$$W : (\Lambda^*)^n \rightarrow \Lambda, \quad W(z_1, \dots, z_n) = \sum_{i=1}^r t^{-\lambda_i} z^{e_i} \quad (12.2)$$

where Λ is the Novikov field defined over $\mathbb{K} = \mathbb{C}$ (see 3.3), and $\Lambda^* = \Lambda \setminus \{0\}$. The fibres $L_y = \mu_X^{-1}(y)$ of the moment map, for $y \in \mathrm{Int}(\Delta) \subset \mathbb{R}^n$, are Lagrangian tori. Let val_t denote the valuation for the t -filtration on Λ , whose value is the lowest exponent of t arising in the series. The conditions $\mathrm{val}_t(t^{-\lambda_i} z^{e_i}) > 0$ are equivalent to the equations in (12.1) since

$$\mathrm{val}_t(t^{-\lambda_i} z^{e_i}) = \langle y, e_i \rangle - \lambda_i. \quad (12.3)$$

Thus z determines a toric Lagrangian L_y , $y = \mathrm{val}_t(z) \in \mathrm{interior}(\Delta) \subset \mathbb{R}^n$, together with a choice of holonomy around each generating circle of the torus L_y given by $t^{-\mathrm{val}_t(z)} z \in H^1(L, \Lambda_0^\times)$, using the ring Λ_0^\times from 3.17. Let L_z be the Lagrangian L_y with that holonomy.

Lemma 12.1 ([22, Theorem 7.11] using [17, Theorem 8.1]).

For a closed monotone toric manifold X , the $z \in \mathrm{Crit}(W)$ always have

$$y = \mathrm{val}_t(z) = \text{barycentre} \in \mathrm{Int}(\Delta). \quad (12.4)$$

and L_y is a monotone Lagrangian. Thus L_z for $z \in \mathrm{Crit}(W)$ is always the same Lagrangian torus L_{crit} (the fibre over the barycentre) but the choice of holonomy data depends on z .

In general, the barycentre $y_{\mathrm{bar}} \in \mathrm{Int}(\Delta)$ of Δ is determined by the equations

$$\langle y_{\mathrm{bar}}, e_i \rangle - \lambda_i = \frac{1}{\lambda_X}. \quad (12.5)$$

In Lemma 12.1, given a critical point z , using (12.3) and (12.5) we deduce

$$\boxed{z = (\zeta_1 t^{a_1}, \dots, \zeta_n t^{a_n}) \quad \text{and} \quad W(z) = c t^{1/\lambda_X}} \quad (12.6)$$

where $\zeta_j \in \mathbb{C}^*$, $c \in \mathbb{C}$, $y_{\text{bar}} = (a_1, \dots, a_n)$ is the barycentre, and recall t^{1/λ_X} has grading 2. Also L_z , being monotone and orientable, defines an object of the (wrapped) Fukaya category in the sense of 3.17. We will only consider cases where X is monotone (so for a Fano toric variety X we always pick a monotone toric symplectic form), and in the non-compact case we require that X is convex. We remark that the proof of the above Lemma still holds in the non-compact setting, after defining *barycentre* as in the Appendix of [36] to mean (12.5). For $z \in \text{Crit}(W)$, it follows [12, Definition 3.3] that

$$m_0(L_z) = W(z).$$

Theorem 12.2 (Cho-Oh [17], see also Auroux [12]).

HF(L_z, L_z) ≠ 0 if and only if z ∈ Crit(W), indeed this holds if and only if the point class [pt] ∈ C*(L_z) is a Floer cycle in HF*(L_z, L_z) in the Morse-Bott model. In particular, if this holds, L_z is not displaceable by a Hamiltonian isotopy.*

Corollary 12.3. $z \in \text{Crit}(W) \Rightarrow \lambda = m_0(L_z) = W(z)$ is an eigenvalue of $c_1(TX) \in QH^*(X)$.

Proof. (Mimicking Auroux [12, Sec.6]) Let $c = c_1(TX) - \lambda \text{Id}$. If c were invertible in $QH^*(X)$, then since $\text{CO}^0 : QH^*(X) \rightarrow HF^*(L_z, L_z)$ is a unital algebra homomorphism (Theorem 8.2),

$$1_{L_z} \equiv \text{CO}^0(1) = \text{CO}^0(c^{-1} * c) = \mu^2(\text{CO}^0(c^{-1}), \text{CO}^0(c)) = 0,$$

as $\text{CO}^0(c) = 0$ by Lemma 9.2. But by Theorem 12.2, the unit $1_{L_z} \neq 0$ as $HF^*(L_z, L_z) \neq 0$. \square

For closed monotone toric manifolds B , a classical result due to Batyrev and Givental is

$$QH^*(B) \cong \text{Jac}(W) = \Lambda[z_1^{\pm 1}, \dots, z_n^{\pm 1}] / (\partial_{z_1} W, \dots, \partial_{z_n} W), \quad c_1(TB) \mapsto W. \quad (12.7)$$

Since the eigenvalues of multiplication by W acting on the Jacobian ring are precisely the critical values of W , one can strengthen Corollary 12.3:

Corollary 12.4. *For closed monotone toric B , the eigenvalues of $c_1(TB) \in QH^*(B)$ are precisely the critical values of W .*

Theorem 12.5 (Galkin [24]). *For B any closed monotone toric manifold, the superpotential $W : (\mathbb{C}^*)^n \rightarrow \mathbb{C}$ (working over \mathbb{C} instead of Λ) always has a non-degenerate critical point $p \in (\mathbb{C}^*)^n$ with strictly positive real coordinates.*

Corollary 12.6. *For B closed monotone and toric, $c_1(TB) \in QH^*(B)$ is not nilpotent.*

Proof. Observe that (putting $t = 1$) $W = \sum z^{e_i}$ will take a strictly positive real value on Galkin's critical point p , so $c_1(TB) \in QH^*(B; \mathbb{C})$ has a strictly positive real eigenvalue λ_p . The same holds true for the superpotential defined over Λ , since we can insert powers of t as dictated by the \mathbb{Z} -grading of $QH^*(B)$ (the eigenvalue will be $\lambda_p t^{1/\lambda_B}$, in grading 2). \square

12.2. Toric negative line bundles. A complex line bundle $\pi : E \rightarrow B$ over a closed symplectic manifold (B, ω_B) is a *negative line bundle* if

$$c_1(E) = -k[\omega_B] \quad \text{for some real } k > 0.$$

We refer to [35, 36] for a detailed exposition. The negative line bundle E is a symplectic manifold conical at infinity whose Reeb flow is the S^1 -rotation in the fibres and whose symplectic form ω can be chosen [36, Section 4] to satisfy

$$\begin{aligned} \omega|_{\text{fibre}=\mathbb{C}} &= \omega_{\mathbb{C}} \\ \omega|_{\text{base}=B} &= \omega_B, \end{aligned}$$

so $[\omega] \equiv [\omega_B]$ via $H^2(E) \cong H^2(B)$. We require that E is monotone, so $c_1(TE) = \lambda_E[\omega]$ for some $\lambda_E > 0$. Splitting TE using a Hermitian connection, $c_1(TE) = \pi^*c_1(TB) + \pi^*c_1(E)$, where $c_1(E)$ is the first Chern class of the line bundle. Therefore, monotonicity of E is equivalent to B being monotone, so $c_1(TB) = \lambda_B[\omega_B]$, and that $0 < k < \lambda_B$. In particular,

$$\lambda_E = \lambda_B - k$$

Our motivation for studying this particular class of non-compact manifolds, is because it is the only family of examples outside of exact cotangent bundles where SH^* was computed:

Theorem 12.7 (Ritter [35]). $c^* : QH^*(E) \rightarrow SH^*(E)$ induces a unital algebra isomorphism

$$SH^*(E) \cong QH^*(E)/QH^*(E)_0$$

(in the notation of 9.1), where we restrict⁶ $SH^*(E)$ to contractible loops (see 4.3). Moreover, $QH^*(E)_0$ is always non-trivial, and $c^*(c_1(TE))$ acts invertibly on $SH^*(E)$ (so $SH^*(E)_0 = 0$).

The Appendix in [36] constructs the moment polytope Δ_E for the complex line bundle $E = \mathcal{O}(\sum n_i D_i) \rightarrow B$, where $D_i \subset B$ are the toric divisors, that is: $c_1(E) = \sum n_i \text{PD}[D_i]$. If the moment polytope for B is $\Delta_B = \{y \in \mathbb{R}^n : \langle y, b_i \rangle \geq \lambda_i^B \text{ for } i = 1, \dots, r\} \subset \mathbb{R}^n$, then take $e_1 = (b_1, -n_1), \dots, e_r = (b_r, -n_r), e_{r+1} = (0, \dots, 0, 1) \in \mathbb{Z}^{n+1}$, $\lambda_i^E = \lambda_i^B$, and $\lambda_{r+1}^E = 0$:

$$\Delta_E = \{y \in \mathbb{R}^{n+1} : \langle y, e_i \rangle \geq \lambda_i^B \text{ for } i = 1, \dots, r \text{ and } y_{n+1} \geq 0\} \subset \mathbb{R}^{n+1}.$$

Although a closed Fano toric variety B has a canonical polytope Δ , with all $\lambda_i^\Delta = -1$, called the reflexive polytope, this will not be useful for our purposes. The reason is that $[\omega_\Delta] = c_1(TB)$, but we want a polytope Δ_B with $[\omega_B] \in H^2(B, \mathbb{Z})$ integral and primitive so that $c_1(TB) = \lambda_B[\omega_B]$, $c_1(E) = -k[\omega_B]$ for integers $\lambda_B > 0$, $k > 0$, because then we can take

$$E = \mathcal{O}(\sum k\lambda_i^B D_i)$$

since the toric symplectic form can be written in terms of the divisors as $[\omega_B] = \sum -\lambda_i^B \text{PD}[D_i]$. If we used Δ instead, then k would be fractional, and it would not be clear how to obtain E as an integral combination $\sum n_i D_i$. Such a polytope Δ_B was constructed in [36] (namely $\Delta_B = (\Delta - v)/\lambda_B$, for a vertex v of Δ). The construction ensures that $\lambda_i^B \leq 0$ are integers and the primitive inward normals b_i are the same as for Δ .

Lemma 12.8. *The barycentres $a \in \mathbb{R}^n$ of Δ_B , and $a' \in \mathbb{R}^{n+1}$ of Δ_E , are related by*

$$a'_j = \frac{\lambda_B}{\lambda_E} a_j \quad \text{and} \quad a'_{n+1} = \frac{1}{\lambda_E} \tag{12.8}$$

Proof. This follows from (12.5) using that $e_j = (b_j, -k\lambda_j^B)$ and $e_{n+1} = (0, \dots, 0, 1)$. \square

Although one could a posteriori translate the polytope Δ_E , so that at least Δ_B has barycentre at 0 (this would make $\Delta_B = \Delta/\lambda_B$ and $\lambda_i^B = -1/\lambda_B$), this would not enable us to make the following substitution trick which relates the superpotentials W_B, W_E of Δ_B, Δ_E , via the change of variables $y = tw^k$, abbreviating by $w = z_{n+1}$ the new variable associated to e_{n+1} .

$$\begin{aligned} W_E &= W_E(z_1, \dots, z_n, w) \\ &= \sum t^{-\lambda_i^B} z^{b_i} w^{-k\lambda_i^B} + w \\ &= \sum y^{-\lambda_i^B} z^{b_i} + w \\ &= W_B(y; z) + w \end{aligned}$$

⁶This is the full $SH^*(E)$ for toric negative line bundles, since toric varieties are simply connected.

where we view the superpotential W_B of B as a function not just of $z = (z_1, \dots, z_n) \in (\Lambda_B^*)^n$ but also of the Novikov variable $T \in \Lambda_B$ (which above undergoes the substitution $T = y$), so

$$W_B = W_B(T; z) = \sum T^{-\lambda_i^B} z^{b_i}.$$

Notice there are two Novikov variables in play: T, t respectively for B, E so respectively in the Novikov rings Λ_B, Λ , and lying in degrees $2\lambda_B, 2\lambda_E$. Our goal is to relate the critical points of W_E, W_B . This relationship is not straightforward because the resulting holonomies of the relevant line bundles on the barycentric torus fibres for B and E are quite different.

12.3. Phases and roots of Novikov series. Given a solution of $dW_B = 0$, to construct a solution of $dW_E = 0$ we may need fractional powers of y (e.g. the critical points for $B = \mathbb{P}^2$ involve $T^{1/3}$). For the purpose of proving the existence of a solution, it will be enough to prove that, given $y \in \Lambda$ involving only strictly positive powers of t , there exists *some* \mathbb{K} -algebra homomorphism $\Lambda_B \rightarrow \Lambda$ sending $T \mapsto y$. Our choice of substitution is governed by the following Definition/Claim.

Definition. $\theta = \theta(y) \in \mathbb{R}$ is a *phase* for $y \in \Lambda$ if $y = y_0 e^{i\theta} t^r + \mathcal{O}$, where we abusively write \mathcal{O} to mean “higher order t terms”, with $y_0 \in \mathbb{R}^*$. The shifts in phase, $\theta \mapsto \theta + 2\pi n$ for $n \in \mathbb{Z}$, preserve y but they will affect our choice of y^p . We call $r = \text{val}_t(y) \in \mathbb{R}$ the *valuation* of y , and $y_0 e^{i\theta} = \text{lead}_t(y) \in \mathbb{C}^*$ the *leading term* of y .

Claim. Given $y \in \Lambda$ with $\text{val}_t(y) > 0$, the group of \mathbb{K} -algebra homomorphisms $\Lambda_B \rightarrow \Lambda$, which send $T \mapsto y$, which are continuous in the T -adic and t -adic topologies, and for which $\text{lead}_t(y^p) \in \mathbb{C}^*$ is continuous in $p \in \mathbb{R}$, is isomorphic to the group \mathbb{Z} (more explicitly, we will construct such a homomorphism for each choice of the phase of y).

Proof. Observe that $y_0^{-1} e^{-i\theta} t^{-r} y = 1 + \mathcal{O}$. For any $p \in \mathbb{R}$, $1 + \mathcal{O}$ has a unique p -th power $(1 + \mathcal{O})^p$ of the form $1 + \mathcal{O}$, obtained by applying $\exp(p \log(\cdot))$ using the explicit series expansion (Newton’s generalised binomial theorem). We now define y^p by the Novikov series

$$y^p = y_0^p e^{ip\theta} t^{pr} (1 + \mathcal{O})^p$$

where $y_0^p \in \mathbb{R}^*$. The relations $y^p y^q = y^{p+q}$ follow because $\exp(p \log(\cdot))$ is a group homomorphism on formal series of type $1 + \mathcal{O}$. Any other choice of p -th power y^p would be $y^p = f(p) \cdot y_0^p e^{ip\theta} t^{pr} (1 + \mathcal{O})^p$ where $f : \mathbb{R} \rightarrow S^1$ is a group homomorphism (in particular $f(0) = 1$) satisfying the additional condition $f(1) = 1$. To ensure that $\text{lead}_t(y^p)$ is continuous in $p \in \mathbb{R}$, f must be continuous. Such f are classified: $f(p) = e^{2\pi i n p}$ for $n \in \mathbb{Z}$, so they are encoded by the shifts in phase $\theta \mapsto \theta + 2\pi n$. \square

Following [36, Lemma 7.15], we now prove that critical points arise in families. Using the above phase formalism, changing the Novikov variable t to $e^{2\pi i} t$ (viewed as a formal change in symbols) induces the following \mathbb{K} -algebra automorphism, which we call *phase shift*,

$$\varphi : \Lambda \rightarrow \Lambda, \sum k_j t^{r_j} \mapsto \sum k_j e^{2\pi i r_j} t^{r_j}.$$

Observe that φ is also t -adically continuous.

Remark. We will show that a non-zero object in the (wrapped) Fukaya category with holonomy systems gives rise to a family of non-zero objects by taking the orbit under the subgroup of the Galois group of Λ over the field $\mathbb{C}[[t]][[t^{-1}]]$ of Laurent series, generated by the phase shift (the full Galois group is a larger, profinite group, as the function f above need not in general be continuous). In retrospect, our approach is similar to that of Fukaya [19, Equation (3)].

Lemma 12.9. Let X be a monotone toric variety, whose polytope (12.1) involves integer parameters λ_i . Then the superpotential (12.2) is φ -equivariant, meaning

$$W \circ \varphi^{\oplus n} = \varphi \circ W,$$

and φ determines a free cyclic action $z \mapsto \varphi(z)$ on $\text{Crit}(W)$ of period λ_X , given on (12.6) by

$$(\zeta_1 t^{a_1}, \dots, \zeta_n t^{a_n}) \mapsto (\zeta_1 e^{2\pi i a_1} t^{a_1}, \dots, \zeta_n e^{2\pi i a_n} t^{a_n}) \quad \text{and} \quad W = c t^{1/\lambda_X} \mapsto c e^{2\pi i/\lambda_X} t^{1/\lambda_X}.$$

Proof. $W(z) = \sum t^{-\lambda_j} z^{e_j}$ so $\varphi(W(z)) = \sum e^{-2\pi i \lambda_j} t^{-\lambda_j} \varphi(z)^{e_j} = W(\varphi(z_1), \dots, \varphi(z_n))$ proves the first claim, using that $e^{-2\pi i \lambda_j} = 1$ since $\lambda_j \in \mathbb{Z}$. Now differentiate, so $(dW)(\varphi^{\oplus n}(z)) = (\varphi \circ dW)(z)$, proving the second claim. \square

We now clarify some notation we use later. Suppose we are given a critical point $z = z(T)$ of W_B , where we emphasize that this is a function of the Novikov variable $T \in \Lambda_B$. The critical value is $W_B(T; z(T))$. Notice that there is a difference between taking the derivative $\partial_T W_B(T; z)$ before evaluating at the critical point $z = z(T)$, and the derivative $\partial_T(W_B(T; z(T)))$ of the critical value. The chain rule implies

$$\frac{\partial}{\partial T} \Big|_{T=y} (W_B(T; z(T))) = \frac{\partial}{\partial T} \Big|_{T=y, z=z(y)} W_B(T; z) + \langle (d_z W_B)(y; z(y)), \frac{\partial}{\partial T} \Big|_{T=y} z(T) \rangle \quad (12.9)$$

where the second term is the sum of $(\partial_{z_j} W_B)(y; z(y)) \cdot \partial_T|_{T=y} z_j(T)$ over $j = 1, \dots, n$.

Example. Suppose $W_B(T; z) = T^3 z^2$ and $z(T) = T^2$, in the $n = 1$ case (so $z \in \Lambda_B^*$). Then $W_B(T; z(T)) = T^7$ so (12.9) reads: $(7T^6)|_{T=y} = (3T^2 z^2)|_{T=y, z=y^2} + (2T^3 z)|_{T=y, z=y^2} (2T)|_{T=y}$.

12.4. Generation results for toric negative line bundles. By “phase-shift invariant family” we mean an orbit of a critical point under the action of φ as in Lemma 12.9.

Lemma 12.10. *Each phase-shift invariant family in $\text{Crit}(W_B)$, with non-zero critical values, gives rise via (12.11)-(12.14) to a unique phase-shift invariant family in $\text{Crit}(W_E)$.*

Proof. Recall $W_E = W_B(y; z) + w$ and $y = tw^k$. By the chain rule, using $\frac{\partial y}{\partial w} = tkw^{k-1}$,

$$dW_E = (d_z W_B)(y; z) + (tkw^{k-1}(\partial_T W_B)(y; z) + 1) dw.$$

It follows that the $(z, w) \in \text{Crit}(W_E)$ are the solutions of the following three equations:

$$\begin{aligned} (i) \quad & tw^k = y \\ (ii) \quad & w = -ky(\partial_T W_B)(y; z) \\ (iii) \quad & (d_z W_B)(y; z) = 0. \end{aligned}$$

Remark. $\text{val}_t(w) > 0$ as $(z, w) \in \text{Crit}(W_E)$ lies in the interior of the polytope, indeed at the barycentre (12.3)-(12.5). So $\text{val}_t(y) > 0$, so we may substitute $T \mapsto y$ as explained in Sec.12.3.

Clarification. We seek a solution $z = z(y) \in (\Lambda^*)^n$, $w = w(y) \in \Lambda^*$ in terms of y , involving the Novikov variable $t \in \Lambda$, not $T \in \Lambda_B$. Observe that

$$z = z(T) \in \text{Crit}(W_B) \subset (\Lambda_B^*)^n \Rightarrow (d_z W_B)(T; z(T)) = 0 \Rightarrow (d_z W_B)(y; z(y)) = 0 \Rightarrow (iii).$$

(However, it is not clear whether the converse $(iii) \Rightarrow z \in \text{Crit}(W_B)$ holds.)

A consequence of (12.9) is that for any Novikov series $z(T)$ with $(z(y), w) \in \text{Crit}(W_E)$ we have $W_B(T; z(T)) \neq 0$, otherwise $(\partial_T W_B)(y; z(y)) = 0$ by (12.9) so $dW_E = dw \neq 0$.

Denote by $z = z(T) \in \text{Crit}(W_B)$ a generator of the given λ_B -family, so by assumption $W_B(z(T)) \neq 0$. We abbreviate

$$\boxed{\ell = \frac{1}{\lambda_B} \quad (\text{in particular } 1 - k\ell = \frac{\lambda_E}{\lambda_B})}$$

By (12.6), $z = z(T) = (\zeta_1 T^{a_1}, \dots, \zeta_n T^{a_n})$ and $W_B(T; z(T)) = cT^\ell$ where $\zeta_j \in \mathbb{C}^*$, $c \neq 0 \in \mathbb{C}$, $a = (a_1, \dots, a_n)$ is the barycentre of Δ_B . Thus

$$W_B(y; z(y)) = cy^\ell. \quad (12.10)$$

By (12.9) and equation (iii), namely $(d_z W_B)(y; z(y)) = 0$, we obtain:

$$(\partial_T W_B)(y; z(y)) = \partial_T|_{T=y}(W_B(T; z(T))) = \partial_T|_{T=y}(cT^\ell) = c\ell y^{\ell-1}.$$

Solving equation (ii), given y , we obtain $(z(y), w(y)) \in \text{Crit}(W_E)$ defined by⁷

$$\boxed{z(y) = (\zeta_1 y^{a_1}, \dots, \zeta_n y^{a_n}) \quad \text{and} \quad w(y) = -kc\ell y^\ell} \quad (12.11)$$

and the critical value is

$$\boxed{W_E = W_E(z(y), w(y)) = W_B(y; z(y)) + w(y) = \frac{\lambda_E}{\lambda_B} cy^\ell \neq 0} \quad (12.12)$$

where $y \neq 0$ since $w \in \Lambda^*$ by definition. It remains to solve for y . Combining (i) and (12.11) gives $y^{\lambda_E/\lambda_B} = (-kc\ell)^k t$, so:

$$\boxed{y = (-kc\ell)^{k\lambda_B/\lambda_E} t^{\lambda_B/\lambda_E}} \quad (12.13)$$

In particular, phases also need to be consistent:

$$\boxed{\theta(y) = \frac{k\lambda_E}{\lambda_E} \theta(c) \quad (\text{in particular } \theta(cy^\ell) = \frac{\lambda_E}{\lambda_E} \theta(c))} \quad (12.14)$$

By Lemma 12.9, the phase shift automorphism for B acts by $\zeta_j \mapsto \zeta_j e^{2\pi i a_j}$ and $c \mapsto ce^{2\pi i/\lambda_B}$. This induces the change $y \mapsto e^{2\pi i k/\lambda_E} y$ and hence $cy^\ell \mapsto e^{2\pi i/\lambda_E} cy^\ell$. Thus (12.11) changes by $\zeta_j y^{a_j} \mapsto \zeta_j e^{2\pi i a_j'} y^{a_j}$ (using $a_j' = a_j \lambda_B/\lambda_E$ from (12.8)) and $w(y) \mapsto e^{2\pi i/\lambda_E} w(y)$, and (12.12) changes by $W_E \mapsto e^{2\pi i/\lambda_E} W_E$. This agrees with the action of the phase shift automorphism for E , by Lemma 12.9. The claim follows. \square

Remark. Applying the phase shift λ_B times preserves $z(T) \in \text{Crit}(W_B)$, but it typically induces a change in $(z(y), w(y)) \in \text{Crit}(W_E)$ as the phase of y undergoes a shift by (12.14).

Theorem 12.11. *Let $E \rightarrow B$ be a monotone negative line bundle over a closed monotone toric symplectic manifold B . Then the monotone Lagrangian torus $\mathcal{L} \subset E$, lying over $L_{\text{crit}} \subset B$ in a sphere bundle of E , is non-displaceable and admits holonomy data such that $HF^*(\mathcal{L}, \mathcal{L}) \cong HW^*(\mathcal{L}, \mathcal{L}) \neq 0$. Moreover, $c_1(E) \in QH^*(E)$ is non-nilpotent and $SH^*(E) \neq 0$.*

Proof. By Corollaries 12.4 and 12.6, there exists a $z = z(T) \in \text{Crit}(W_B)$ with non-zero critical value. Lemma 12.10 implies that $\text{Crit}(W_E) \neq \emptyset$, so the first claim follows by Theorem 12.2 and Lemma 5.1. In particular, it is proved in the Appendix of [36] (by the same argument used for Lemma 12.1, using [17, Theorem 8.1] and (12.5)), that the toric fibre \mathcal{L} of the moment map for E over the barycentre of Δ_E is a monotone Lagrangian torus in a sphere bundle of E lying over the monotone Lagrangian torus $L_{\text{crit}} \subset B$ (corresponding to the barycentre of Δ_B). The sphere bundle is defined using an appropriate Hermitian metric [35, Section 7], and the radius is determined by the monotonicity condition. Concretely, \mathcal{L} has an additional fibre circle factor, arising from the natural $U(1)$ -toric rotation action on the fibres, and the natural filling disc u of that circle in the \mathbb{C} -fibre must satisfy the monotonicity condition: $\text{MaslovIndex}([u]) = 2\lambda_E \int_{\mathbb{D}} u^* \omega_E$.

That $c_1(TE) \in QH^*(E)$ is not nilpotent (so $SH^*(E) \neq 0$ by Theorem 12.7) follows because (12.12) is an eigenvalue of $c_1(TE)$ by Corollary 12.3. \square

Lemma 12.12. *All the critical points of W_E arise via (12.11)-(12.14) from the critical points of W_B which have non-zero critical value.*

⁷**Remark.** We already know a priori that $(z(y), w(y))$ lies at the barycentre of Δ_E by (12.5) (for $X = E$), but one can also check by hand that $\text{val}_t(z(y)) = \frac{\lambda_E}{\lambda_E} a$ and $\text{val}_t(w(y)) = \frac{1}{\lambda_E}$ (confirming (12.8) and (12.5)) and in particular $\text{val}_t(z, w) > 0$ in each entry, so $(z, w) \in \text{Int}(\Delta_E)$.

Proof. We reverse the previous proof: if $(z', w') \in \text{Crit}(W_E)$, we want to find $z = z(T) \in \text{Crit}(W_B)$ such that $z' = z(y)$, $w' = w(y)$ for some solution y of (12.13). By (12.5) for $X = E$,

$$z' = (\zeta'_1 t^{a'_1}, \dots, \zeta'_n t^{a'_n}) \quad w' = \xi' t^{\ell'} \quad W_E(z', w') = c' t^{\ell'}$$

where $\ell' = \frac{1}{\lambda_E} = \frac{1}{\lambda_B - k}$, $a' \in \mathbb{R}^n$ is the barycentre of Δ_E so (12.8) holds, $\zeta' \in (\mathbb{C}^*)^n$, $\xi' \in \mathbb{C}^*$, $c' \in \mathbb{C}$. By (12.6) we want $z = (\zeta_1 t^{a_1}, \dots, \zeta_n t^{a_n})$, with $\zeta \in (\mathbb{C}^*)^n$ and $a \in \mathbb{R}^n$ the barycentre of Δ_B , such that $z' = z(y)$, $w' = w(y)$ where

$$y = t w'^k = \xi'^k t^{\lambda_B / \lambda_E}$$

(using $1 + k\ell' = \lambda_B / \lambda_E$). So $z' = z(y)$ forces the equality: $\zeta'_j t^{a'_j} = \xi'^k t^{\lambda_B / \lambda_E} \zeta_j t^{a_j \lambda_B / \lambda_E}$. By (12.8), this equality holds for $z(T)$ by taking

$$\boxed{\zeta_j = \zeta'_j \xi'^{-k a_j}, \text{ or equivalently: } z_j(T) = (z'_j w'^{-k a_j})|_{t=T}} \quad (12.15)$$

By equation (iii) of the previous proof, $(d_z W_B)(y; z(y)) = 0$. We claim that $(d_z W_B)(T; z) = 0$, so $z = z(T) \in \text{Crit}(W_B)$ (compare the Clarification in the proof of Lemma 12.10). Indeed, because $y = \xi'^k t^{\lambda_B / \lambda_E}$ is homogeneous in t , y plays the same role as a formal variable in the equation (iii): $(d_z W_B)(y, z(y)) = 0$, so we can replace it with the formal variable T . \square

Corollary 12.13. *There are precisely $(\lambda_B - k)N / \lambda_B$ distinct choices of holonomy for $\mathcal{L} \subset E$ in Theorem 12.11 for which $HF^*(\mathcal{L}, \mathcal{L}) \neq 0$, where $N > 0$ is the number of critical points $z \in \text{Crit}(W_B)$ with non-zero critical value $W_B(T; z) \neq 0$. Moreover, the critical values of W_E are $\mu t^{1/\lambda_E}$ where $\mu \in \mathbb{C}^*$ are the roots of*

$$\left(\frac{\mu}{\lambda_E}\right)^{\lambda_E} = \left(\frac{c}{\lambda_B}\right)^{\lambda_B} (-k)^k \quad (12.16)$$

as cT^{1/λ_B} varies over possible non-zero critical values of W_B , and recall $\lambda_E = \lambda_B - k$.

Proof. The first claim follows by the one-to-one correspondence between families of critical points established by Lemmas 12.10 and 12.12. That critical values of W_E correspond to eigenvalues of $c_1(TE)$ follows by Corollary 12.3. By (12.12), the W_E -critical value is $\mu = c \frac{\lambda_E}{\lambda_B} y^\ell$, so by (12.13), $\mu^{\lambda_E} = (c \frac{\lambda_E}{\lambda_B})^{\lambda_E} (-k c \ell)^k t = \left(\frac{c}{\lambda_B}\right)^{\lambda_B} \lambda_E^{\lambda_E} (-k)^k t$. \square

Conjecture. $\mathcal{W}(E)$ is split-generated by \mathcal{L} taken with the holonomies from Corollary 12.13. This holds when W_B is Morse by [36], but it should also hold in general.

Remark 1. In fact [36], $\pi^*[\omega_B] = \frac{c_1(TE)}{\lambda_E}$, $[\omega_B] = \frac{c_1(TB)}{\lambda_B}$ and $c_1(E)^k = (-k)^k [\omega_B]^k$ are the cause of the constant factors in (12.16).

Remark 2. By Theorem 12.7, zero-eigenvalues for $c_1(TE)$ must exist, but in the proof of Lemma 12.10 we showed these cannot arise from critical values of W_E (which are non-zero). It turns out [36] that all the non-zero eigenvalues of $c_1(TE)$ arise as critical values of W_E .

12.5. Generation for $\mathcal{O}(-k) \rightarrow \mathbb{P}^m$. We will consider the negative line bundles $E = \mathcal{O}(-k)$ over $B = \mathbb{P}^m$, so $c_1(E) = -k[\omega_{\mathbb{P}^m}]$. These are monotone for $1 \leq k \leq m$. Moreover,

$$\lambda_B = 1 + m \quad \lambda_E = 1 + m - k.$$

Theorem 12.14 (Ritter [35, 36]). *Let $E = \text{Tot}(\mathcal{O}(-k) \rightarrow \mathbb{P}^m)$ with $1 \leq k \leq m$. Then $QH^*(E)$ and $SH^*(E)$ are generated over Λ by $w = \pi^*[\omega_B] \in QH^*(E)$ (viewing $SH^*(E)$ as a quotient of $QH^*(E)$ by Theorem 12.7) subject to the relations*

$$\begin{aligned} w^k (w^{1+m-k} - (-k)^k t) &= 0 & \text{in } QH^*(E), \\ w^{1+m-k} - (-k)^k t &= 0 & \text{in } SH^*(E). \end{aligned}$$

Recall that $QH^*(\mathbb{P}^m) \cong \Lambda[w]/(w^{1+m} - T)$, $\omega_{\mathbb{P}^m} \mapsto w$, where T is the Novikov variable for $B = \mathbb{P}^m$. So the eigenvalues of $c_1(TB) \cdot \ast$ on $QH^*(B)$ involve $(1+m)$ -th roots of unity (up to a fixed constant factor). On the other hand, for $E = \mathcal{O}_{\mathbb{P}^m}(-k)$ the above calculation shows that the non-zero eigenvalues of $c_1(TE) \cdot \ast$ on $QH^*(E)$ involve $(1+m-k)$ -th roots of unity (up to a constant factor). The moment polytope for $B = \mathbb{P}^m$ is determined by $b_1 = (1, 0, \dots, 0)$, \dots , $b_m = (0, \dots, 0, 1)$, $b_{m+1} = (-1, \dots, -1)$, and $\lambda_i^B = 0$ for $i \leq m$ and $\lambda_{m+1}^B = -1$, so:

$$\begin{aligned} \Delta_B &= \{y \in \mathbb{R}^m : y_1 \geq 0, \dots, y_m \geq 0, \sum y_j \leq 1\} \\ W_B(T; z) &= z_1 + \dots + z_m + Tz_1^{-1} \dots z_m^{-1}. \end{aligned}$$

The critical points of W_B are: for each λ_B -th root of unity ζ , there is one critical point

$$z_{\text{crit}} = (1, \dots, 1) \cdot \zeta T^{1/\lambda_B} \quad W_B(z_{\text{crit}}) = \lambda_B \zeta T^{1/\lambda_B}.$$

The toric divisors $D_i \subset B$ are given by setting the i -th homogeneous coordinate of \mathbb{P}^m to zero (using the 0-th coordinate for D_{m+1}). So each D_i is Poincaré dual to ω_B . The moment polytope for $E = \mathcal{O}(-k) = \mathcal{O}(-kD_{m+1}) \rightarrow \mathbb{P}^m$ is determined by $e_1 = (b_1, 0), \dots, e_m = (b_m, 0), e_{m+1} = (b_{m+1}, k)$, and $\lambda_i^E = 0$ for all $i = 1, \dots, m+2$ except $\lambda_{m+1}^E = \lambda_{m+1}^B = -1$, so:

$$\begin{aligned} \Delta_E &= \{y \in \mathbb{R}^{m+1} : y_1 \geq 0, \dots, y_m \geq 0, \sum_{j=1}^m y_j - ky_{m+1} \leq 1, y_{m+1} \geq 0\} \\ W_E &= z_1 + \dots + z_m + tz_1^{-1} \dots z_m^{-1} w^k + w = W_B(tw^k; z) + w. \end{aligned}$$

The critical points of W_E are: for each λ_E -th root of unity ζ' , the solution $x = \zeta'(-k)^{k/\lambda_E} t^{1/\lambda_E}$ of $x^{1+m-k} = (-k)^k t$ corresponds to a critical point $(x, \dots, x, -kx)$ with $W_E = \lambda_E x$:

$$z_{\text{crit}} = (1, \dots, 1, -k) \cdot \zeta' (-k)^{k/\lambda_E} t^{1/\lambda_E} \quad W_E(z_{\text{crit}}) = \lambda_E \zeta' (-k)^{k/\lambda_E} t^{1/\lambda_E}.$$

In particular, we obtain

$$y = tw^k = t(-k)^k x^k = (-k)^{k\lambda_B/\lambda_E} (\zeta')^k t^{\lambda_B/\lambda_E}.$$

Since (12.13) becomes $y = (-k\zeta)^{k\lambda_B/\lambda_E} t^{\lambda_B/\lambda_E}$, the one-to-one correspondence between families of critical points for B and for E boils down to the relation $(\zeta')^{\lambda_E} = \zeta^{\lambda_B} = 1$, and one can confirm by hand that (12.11), (12.12), (12.15), (12.16) indeed hold.

The Jacobian ring $\text{Jac}(W_E) \cong \Lambda[z_1^{\pm 1}, \dots, z_m^{\pm 1}, w^{\pm 1}]/(\partial_{z_1} W_E, \dots, \partial_{z_m} W_E, \partial_w W_E)$ is

$$\text{Jac}(W_E) \cong \Lambda[x]/(x^{1+m-k} - (-k)^k t).$$

Corollary 12.15. $SH^*(E) \cong \text{Jac}(W_E)$ for $1 \leq k \leq m$, and $W_E = (1+m-k)x \mapsto c_1(TE)$.

This Corollary is noteworthy since it is a proof of closed-string mirror symmetry in this example. Notice that $SH^*(E)$ now plays the role that $QH^*(E)$ did in the closed case (12.7).

Theorem 12.16. Let $E = \mathcal{O}_{\mathbb{P}^m}(-k)$, for $1 \leq k \leq m$, and let \mathcal{L} be the lift of the Clifford torus in the base \mathbb{P}^m to the sphere subbundle. Then

- (1) The wrapped category $\mathcal{W}(E)$ is compactly generated. Indeed, for each of the $1+m-k$ choices of non-zero $\lambda \in \text{Spec}(c_1(TE))$, the Fukaya categories $\mathcal{W}_\lambda(E)$ and $\mathcal{F}_\lambda(E)$ are both split-generated by \mathcal{L} with a choice of holonomy determined by λ .
- (2) $\mathcal{W}(E)$ is proper (cohomologically finite), so $\dim HW^*(L_1, L_2) < \infty$ for any L_1, L_2 .
- (3) $HF^*(\mathcal{L}, \mathcal{L}) \cong HW^*(\mathcal{L}, \mathcal{L}) \neq 0$ (Theorem 12.11), so \mathcal{L} is non-displaceable.
- (4) $\text{OC}^0 : HW^*(\mathcal{L}, \mathcal{L}) \rightarrow SH^*(E)$ is non-zero.

Proof. By Theorem 12.14, $QH^*(E)$ splits as a direct sum of $QH^*(E)_0$ and 1-dimensional field summands corresponding to the non-zero eigenvalues, and by Theorem 12.7, the map $c^* : QH^*(E) \rightarrow SH^*(E)$ is the quotient by $QH^*(E)_0$.

Let $\mathcal{L} = \mathcal{L}_z$ be a Lagrangian torus in E with appropriate holonomy arising from $z \in \text{Crit}(W_E)$. Call $\lambda_{\mathcal{L}} = m_0(\mathcal{L}) = W_E(z)$ the non-zero eigenvalue of $c_1(TE)$ corresponding to \mathcal{L} .

By Theorem 12.14 and the above calculation of the critical points of W_E , all except for the 0-eigenvalue of $c_1(TE)$ arise as $\lambda_{\mathcal{L}}$ for some critical point.

Using the Morse-Bott model [12], and using that [pt] is a cycle by Theorem 12.2, observe that $\text{OC} : HF^*(\mathcal{L}, \mathcal{L}) \rightarrow QH^*(E)$ satisfies

$$\text{OC}([\text{pt}]) = -kx \cdot \pi^* \omega_{\mathbb{P}^m}^m + (\text{higher order } t \text{ terms}) \neq 0 \in QH^*(E),$$

where the first term arises from the regular Maslov 2 disc in the fibre (the standard disc filling $S^1 \subset \mathbb{C}$) and $-kx \neq 0 \in \Lambda$ is the holonomy for the S^1 -fibre direction of \mathcal{L} (here $-kx$ also includes the factor t^{1/λ_E} lying in degree 2). Observe that there is no order t^0 term because the constant disc would sweep the lf-cycle $[\text{pt}] \in H_0^{lf}(E) \cong H^{2m+2}(E) = 0$. The $\pi^* \omega_{\mathbb{P}^m}^m$ factor arises because the fibre disc intersects once transversely the cycle intersection condition $[\mathbb{P}^m] \in C_*(E)$, and the intersection dual of this is the lf-cycle $[\text{fibre}] \in H_2^{lf}(E)$ which corresponds to $\pi^* \omega_{\mathbb{P}^m}^m \in H^{2m}(E)$ (the fibre is the preimage of a point in the base, and taking preimages is Poincaré dual to pull-back on cohomology). It is a general feature of the toric setting that the Maslov index of a disc bounding a toric Lagrangian is twice the number of intersections of the disc with the anticanonical divisor (see Auroux [12, Lemma 3.1]). In our case, the anticanonical divisor viewed as an lf-cycle, consists of the base \mathbb{P}^m together with the preimage under projection of the anticanonical divisor in the base \mathbb{P}^m . So the fibre disc is the only Maslov two disc that can intersect the cycle $[\mathbb{P}^m]$. By monotonicity, all other discs with Maslov index greater than 2 will contribute higher order t terms. Thus:

- (i) By Theorem 9.5, $\text{OC}([\text{pt}])$ can only be non-zero in the eigensummand $QH^*(E)_{\lambda_{\mathcal{L}}}$.
- (ii) That eigensummand is a field, $\text{OC}([\text{pt}])$ is non-zero, hence $\text{OC}([\text{pt}])$ is invertible in the eigensummand.
- (iii) Thus, by Theorem 11.6, \mathcal{L} split-generates $\mathcal{F}_{\lambda_{\mathcal{L}}}(E)$.
- (iv) Moreover, by the acceleration diagram (Theorem 7.2), also the map $\text{OC}^0 : HW^*(\mathcal{L}, \mathcal{L}) \rightarrow SH^*(E)$ hits an invertible element, since by Theorem 12.14 all eigensummands except the 0-eigensummand survive via $c^* : QH^*(E) \rightarrow SH^*(E)$.
- (v) By Theorem 11.3, \mathcal{L} split-generates $\mathcal{W}_{\lambda_{\mathcal{L}}}(E)$.

The finite-dimensionality claim follows because the split-generators are closed Lagrangians. In particular, $HF^*(\mathcal{L}, \mathcal{L}) \cong HW^*(\mathcal{L}, \mathcal{L})$ (Lemma 5.1) is generated at the chain level by the finitely many intersections of \mathcal{L} with a Hamiltonianly deformed copy of \mathcal{L} . \square

Remark 1. *It is not clear what the split-generators for $\mathcal{F}_0(E)$ might be. By Theorem 12.2, they cannot be Lagrangian torus fibres of the moment map, with holonomy data, as those would be detected by W_E . But W_E does not detect the zero eigenvalues of $c_1(TE)$.*

Remark 2. *We emphasise that although the wrapped category is proper (cohomologically finite), it does contain objects which are non-compact Lagrangians – such as real line bundles over suitable Lagrangians in the base or Lefschetz thimbles for appropriate Lefschetz fibrations – whose wrapped Floer cochain complexes are infinitely generated.*

12.6. Generation results for closed monotone manifolds. In Lemma 12.17, we use the Morse-Bott model $CF^*(L, L) = C_{n-*}(L)$ of [12], so there is a well-defined element $[\text{pt}] \in CF^n(L, L)$. Note this need not a priori be a Floer cocycle (if it is not, the Lemma gives no information, as $[\text{pt}]$ will not survive to cohomology).

Lemma 12.17. *For any closed monotone manifold B , and any monotone Lagrangian $L \in \text{Ob}(\mathcal{F}_{\lambda}(B))$, the map $\text{OC}^0 : CF^*(L, L) \rightarrow QC^*(B)$ satisfies*

$$\text{OC}([\text{pt}]) = \text{vol}_B + (\text{higher order } t \text{ terms}) \neq 0 \in QH^*(B).$$

Proof. Since $L \subset B$ is monotone, Maslov index zero discs are constant. So the first term above corresponds to the constant disc, which is the lowest Maslov index disc, using that the volume form vol_B is the Poincaré dual of $[\text{pt}] \in H_0(B)$. All other terms will involve strictly positive powers of t since they involve non-constant discs. \square

Theorem 12.18. *Let B be a closed monotone toric manifold, such that the eigenvalues of $c_1(TB)$ acting on $QH^*(B)$ are all distinct and have multiplicity one. Then $\mathcal{F}(B)$ is split-generated by L_{crit} with the holonomies described in Lemma 12.1.*

Proof. The assumption on the eigenvalues of $c_1(TB)$ ensures that $QH^*(B)$ splits into 1-dimensional eigensummands. By (12.7), each 1-dimensional eigensummand corresponds to a critical point $z \in \text{Crit}(W_B)$ which determines a toric Lagrangian L_z with suitable holonomy (see 12.1). By Lemma 12.17, $\text{OC}^0 : CF^*(L_z, L_z) \rightarrow QH^*(B)$ is non-zero on $[\text{pt}]$. By Theorem 12.2, $[\text{pt}] \in CF^*(L_z, L_z)$ is a cycle, so $\text{OC}^0 : HF^*(L_z, L_z) \rightarrow QH^*(B)$ is non-zero.

By the $QH^*(B)$ -module structure of OC , the element $\text{OC}([\text{pt}])$ can only be non-zero in the eigensummand $QH^*(B)_{\lambda_{L_z}}$ corresponding to the eigenvalue $\lambda_{L_z} = W_B(z) = m_0(L_z)$ of $c_1(TB)$. Being non-zero in a field implies invertibility, so OC hits an invertible element in the relevant eigensummand. By Theorem 11.6, the Lagrangians L_z split-generate $\mathcal{F}(B)$. \square

12.7. Generation results for convex symplectic manifolds. Let (E, ω) be a monotone symplectic manifold conical at infinity. Let $L \in \text{Ob}(\mathcal{F}(E)_\lambda)$, so $L \subset E$ is a monotone closed orientable Lagrangian with $m_0(L) = \lambda$.

Theorem 12.19. *If $\lambda \neq 0$ and $[\text{pt}]$ is a Floer cycle in $HF^*(L, L)$, then*

- (1) $\text{OC}([\text{pt}]) \neq 0 \in QH^*(E)$;
- (2) *if the generalized eigensummand $QH^*(E)_\lambda$ is 1-dimensional, then L split-generates $\mathcal{F}(E)_\lambda$ and $\mathcal{W}(E)_\lambda$, in particular $\mathcal{W}(E)_\lambda$ is proper (cohomologically finite).*

Proof. By Lemma 9.2, the image under $\text{CO}^0 : QH^*(E) \rightarrow HF^*(L, L)$ of $c_1(TE)$ is $m_0(L)e_L$. Now $c_1(TE)$ is a multiple of $[\omega]$, and ω is exact at infinity. Therefore $c_1(TE)$ can be represented by a closed de Rham form which is compactly supported, which in turn is Poincaré dual to a cycle $C \in H_2(E)$. Thus we can represent $c_1(TE)$ by a *compact* lf-cycle C . Since the disc count above does not depend on the choice of representative, we can use C .

Denote by $\langle \cdot, \cdot \rangle_E$ the intersection pairing in E between lf-cycles and cycles, and by $\langle \cdot, \cdot \rangle_L$ the intersection pairing in L between cycles. Then observe that, by definition, the following two disc counts are equal:

$$\langle \text{OC}([\text{pt}]), C \rangle_E = \langle [\text{pt}], \text{CO}^0(C) \rangle_L = m_0(L). \quad (12.17)$$

This proves the first claim since $m_0(L) = \lambda \neq 0$ by assumption, so $\text{OC}([\text{pt}]) \neq 0$. The second claim follows by the same argument as in the proof of Theorem 12.16 and by the acceleration diagram (Theorem 7.2). \square

Corollary 12.20. *Let E be any (non-compact) toric Fano variety convex at infinity. Let $z \in \text{Crit}(W)$ be a critical point of the superpotential W with non-zero critical value $\lambda = W(z) \neq 0$. If the generalized eigensummand $QH^*(E)_{W(z)}$ is 1-dimensional then the Lagrangian torus L_z split-generates $\mathcal{F}(E)_\lambda$ and $\mathcal{W}(E)_\lambda$, in particular $\mathcal{W}(E)_\lambda$ is proper (cohomologically finite).*

Proof. The fact that $[\text{pt}]$ is a Floer cycle in $HF^*(L_z, L_z)$ follows by Theorem 12.2. The Lagrangian L_z (with holonomy) associated to z has eigenvalue $\lambda = m_0(L_z) = W(z) \neq 0$. So the claim follows by the previous Theorem. \square

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A. F. RITTER, MATHEMATICAL INSTITUTE, UNIVERSITY OF OXFORD, ENGLAND.
E-mail address: ritter@maths.ox.ac.uk

I. SMITH, CENTRE FOR MATHEMATICAL SCIENCES, UNIVERSITY OF CAMBRIDGE, ENGLAND.
E-mail address: I.Smith@dpms.cam.ac.uk